

Two-Photon Ionization of He through a Superposition of Higher Harmonics

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Outline

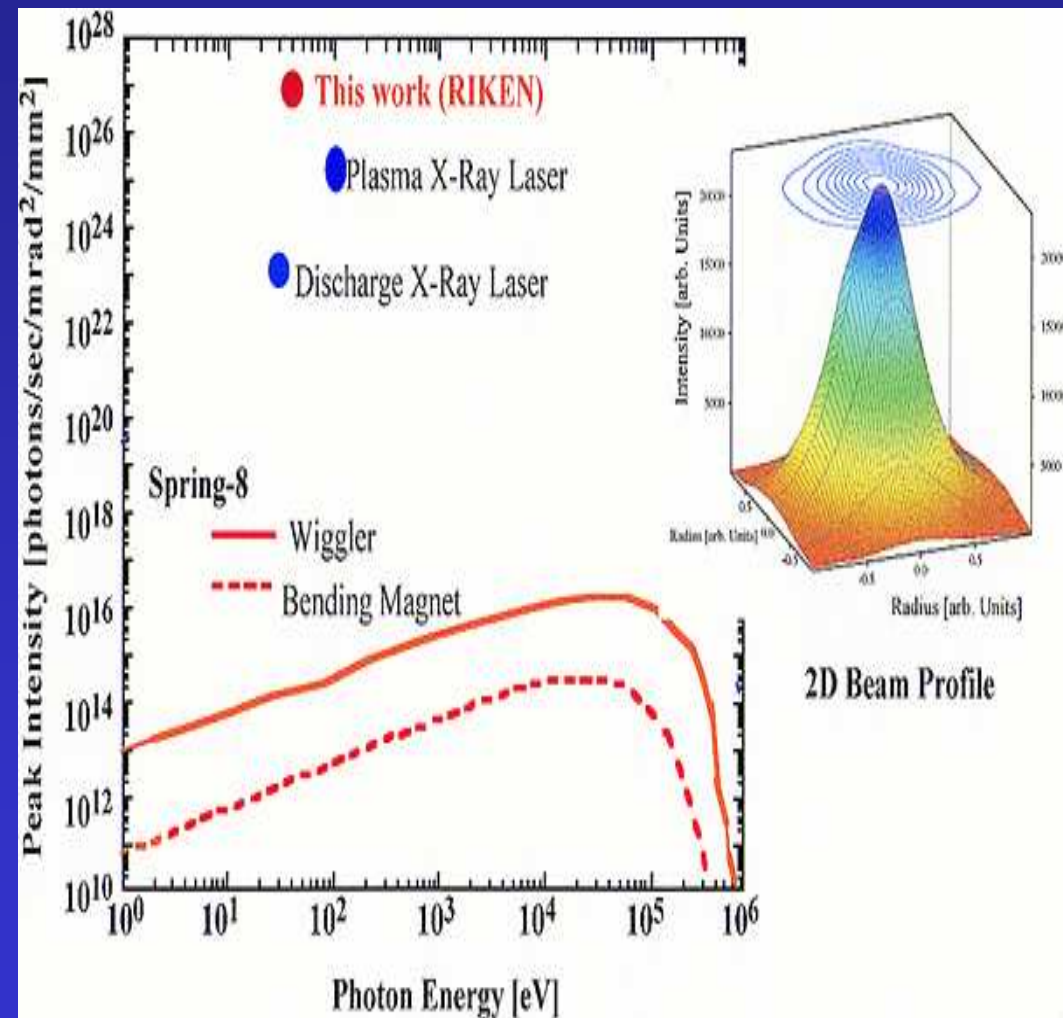
- Motivation
- The applied Coupled-Channel Method
- Results for Photoionization
- Two-Photon Single-Ionization
- Summary

Motivation

it is experimentally possible to create XUV laser pulses with HHG

- two-photon ionization can be studied
- electron-electron correlation can be investigated

<http://www.riken.jp/engn/rworld/research/lab/wako/laser-tech/result.htm>



Our coupled-channel method

- The applied Coupled-Channel Method was originally developed for ion-He collision (details in I.F. Barna, *Ionization of helium in relativistic heavy-ion collisions*, Doctoral Thesis, University of Giessen (2002) „Giessener Elektronische Bibliothek“

<http://geb.uni-giessen.de/geb/volltexte/2003/1036>

I.F. Barna, N. Grün and W. Scheid, *Eur. Phys. J. D* **25**, (2003) 239

I.F. Barna, K. Tókési and J. Burgdörfer, *J. Phys. B* **38**, (2005) 1

- Later modified to study photoionization and coherent control

I.F. Barna, J.M. Rost, *Eur. Phys. J. D* **27**, (2003) 287

I.F. Barna, *Eur. Phys. J. D* **33**, (2005) 307

I.F. Barna, J. Wang and J. Burgdörfer, *Phys. Rev. A* **73**, 023402 (2006)

The Hamiltonian of the system:

The time-dependent Schrödinger equation:

$$i\frac{\partial}{\partial t}\Psi(\vec{r}_1, \vec{r}_2, t) = (\hat{H}_{He} + \hat{V}(t))\Psi(\vec{r}_1, \vec{r}_2, t)$$

- $\hat{V}(t)$ laser-electron interaction
- $\Psi(\vec{r}_1, \vec{r}_2, t)$ CI. wavefunction of helium

The unperturbed helium Hamiltonian:

$$\hat{H}(\vec{r}_1, \vec{r}_2)_{He} = -\frac{\vec{\nabla}_1^2}{2} - \frac{\vec{\nabla}_2^2}{2} - \frac{2}{r_1} - \frac{2}{r_2} + \frac{1}{|\vec{r}_1 - \vec{r}_2|}$$

- spin-spin, spin-orbit, and mass polarisation terms are neglected

The coupled-channel equation

Ansatz:

$$\Psi(\vec{r}_1, \vec{r}_2, t) = \sum_{j=1}^N a_j(t) \Phi_j(\vec{r}_1, \vec{r}_2) e^{-iE_j t}$$

Leads to a system of first-order-differential equations for the coefficients a_j :

$$\frac{da_k(t)}{dt} = -i \sum_{j=1}^N V_{kj}(t) e^{-i(E_j - E_k)t} a_j(t) \quad (k = 1 \dots N)$$

$$V_{kj}(t) = \langle \Phi_k | \hat{V}(t) | \Phi_j \rangle \quad \text{coupling matrix}$$

Initial conditions:

$$a_k(t \rightarrow -\infty) = \begin{cases} 1 & k = 1 \\ 0 & k \neq 1 \end{cases}$$

the final probability for each channel:

$$P_k(b, t \rightarrow \infty) = |a_k(t \rightarrow \infty)|^2$$

The wavefunction

Configuration interaction(CI) expansion of $\Phi_j(\vec{r}_1, \vec{r}_2)$ in terms of two-particle basis functions f_μ

$$\Phi_j(\vec{r}_1, \vec{r}_2) = \sum_{\mu} b_{\mu}^j f_{\mu}(\vec{r}_1, \vec{r}_2).$$

where $f_{\mu}(\vec{r}_1, \vec{r}_2)$ are symmetric ($S=0$) products of

1. **Slater**-type orbitals:

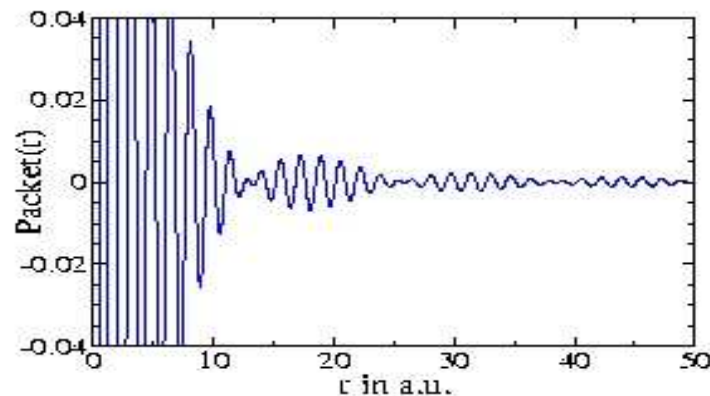
$$\chi_{n,l,m,\kappa}(\vec{r}) = C(n, \kappa) r^{n-1} e^{-\kappa r} Y_{l,m}(\theta, \varphi)$$

2. regular **Coulomb** wave packets:

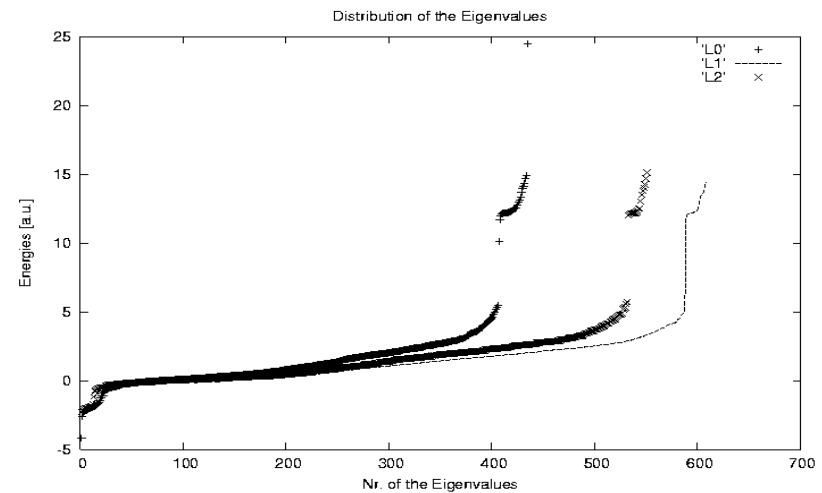
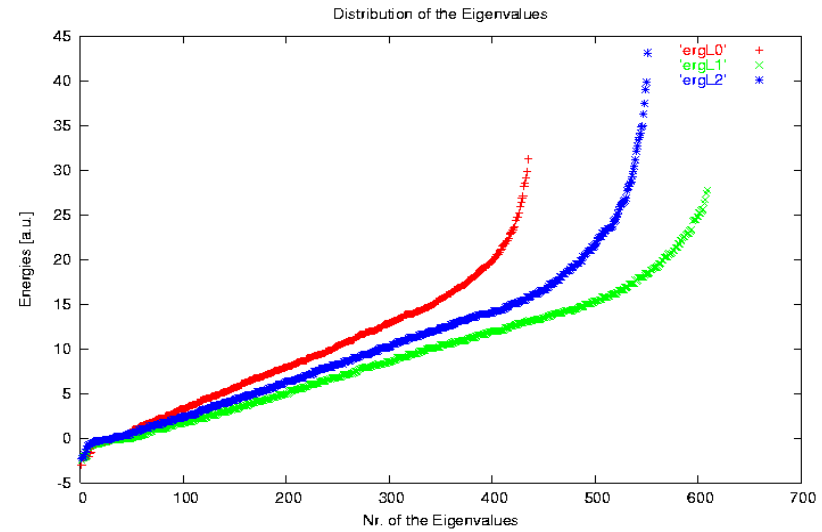
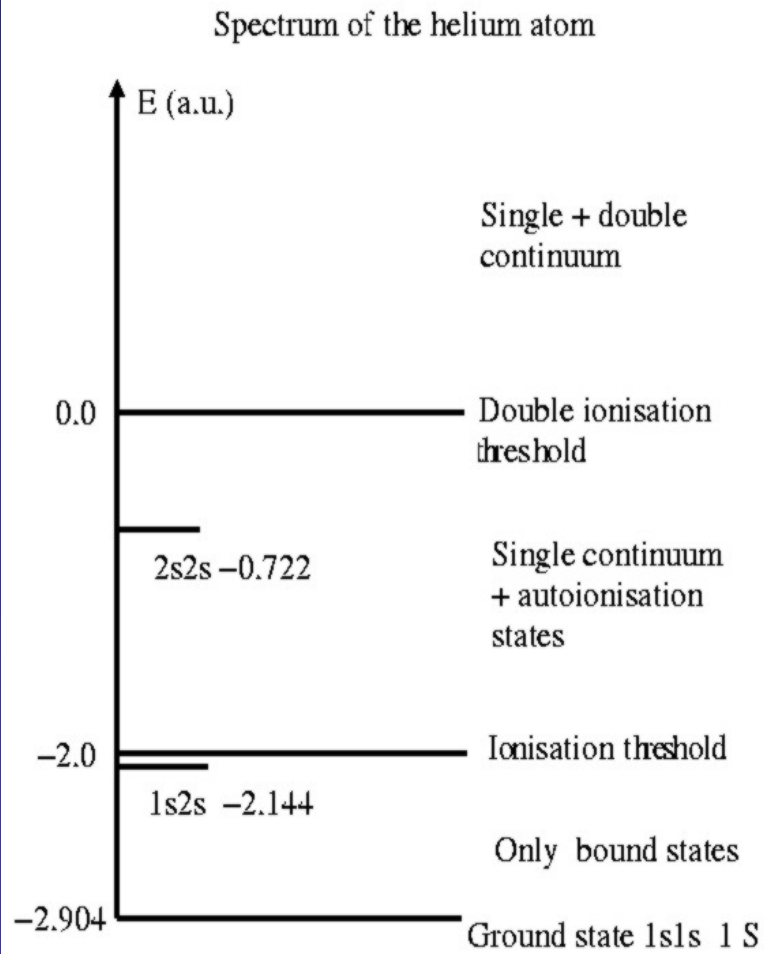
$$\varphi_{k,l,m,\tilde{Z}}(\vec{r}) = N(k, \Delta k) \int_k^{k+\Delta k} R_l(\eta, \rho) dk' Y_{l,m}(\theta, \varphi)$$

- $\eta = \tilde{Z}/k'$, $\rho = k'r$, \tilde{Z} effective charge
- $N(k, \Delta k)$, $C(n, \kappa)$ normalisation constants

$Z_{\text{eff}} = 2, l = 0, k = 3.5 \text{ a.u.}, \Delta k = 0.5 \text{ a.u.}$



Spectrum & Continuum States



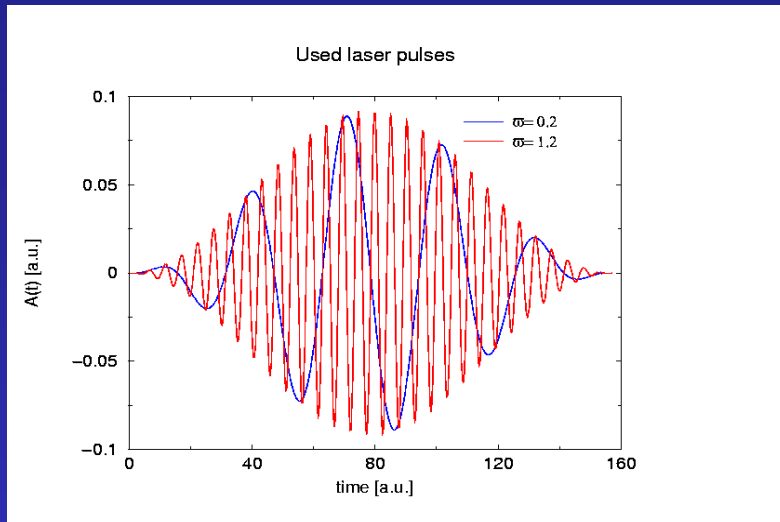
Laser field

minimal coupling for the electromagnetic field to \hat{H}_{He}

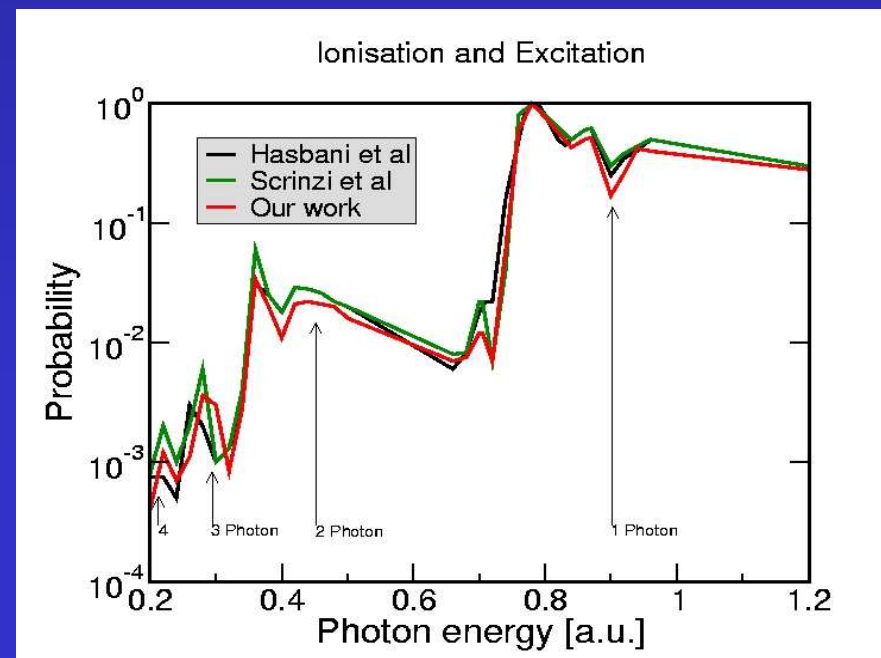
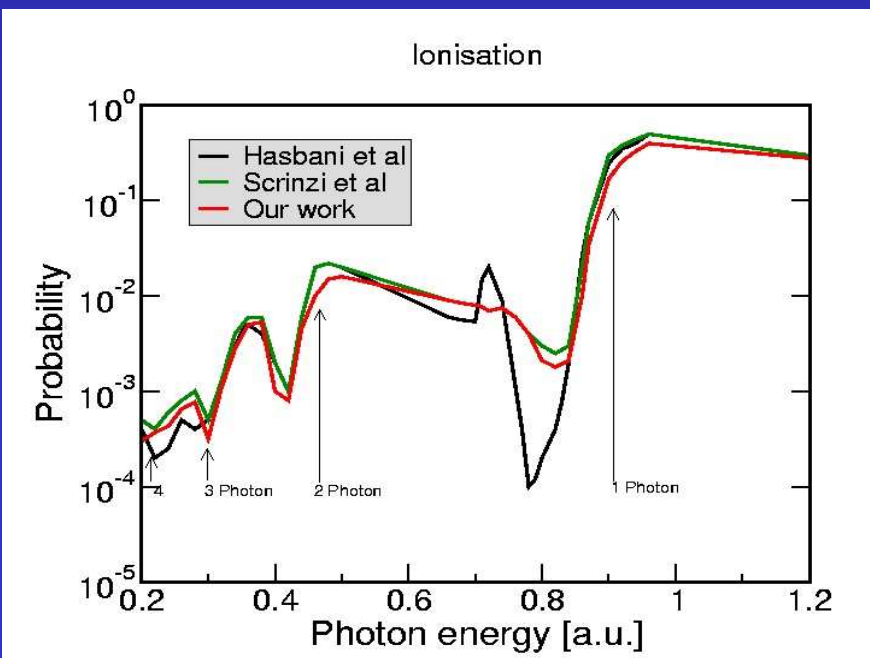
$$\hat{H}(\vec{r}_1, \vec{r}_2, t) = \frac{(\vec{p}_1 - \vec{A}(\vec{r}_1, t)/c)^2}{2} + \frac{(\vec{p}_2 - \vec{A}(\vec{r}_2, t)/c)^2}{2}$$
$$-\frac{2}{r_1} - \frac{2}{r_2} + \frac{1}{r_{12}} - \phi(\vec{r}_1, t) - \phi(\vec{r}_2, t) = \hat{H}_{He} + \hat{V}(t)$$

- **dipole approximation:** $\lambda_{laser} \gg r_{atom}$
- **velocity gauge:** $\hat{V}(\vec{r}, t) = \sum_{i=1,2} \vec{A}(t) \cdot \vec{p}_i$
- **length gauge:**
 $\hat{V}(\vec{r}, t) = - \sum_{i=1,2} \vec{E}(t) \cdot \vec{r}_i, \quad \vec{E}(t) = -\frac{1}{c} \frac{\partial}{\partial t} \vec{A}(t)$
- $\vec{E}(t) = E_0 \cdot F(t) \cdot \sin(\omega t) \cdot \vec{e}_z$ linearly polarised pulse
- **envelope function:** $F(t) = \sin^2\left(\frac{\pi t}{T}\right)$ T is the pulse duration
- or $F(t) = \exp\left(-\frac{(2\ln 2)t^2}{T^2}\right)$ T is the FWHM

Laser pulse & Photoionization



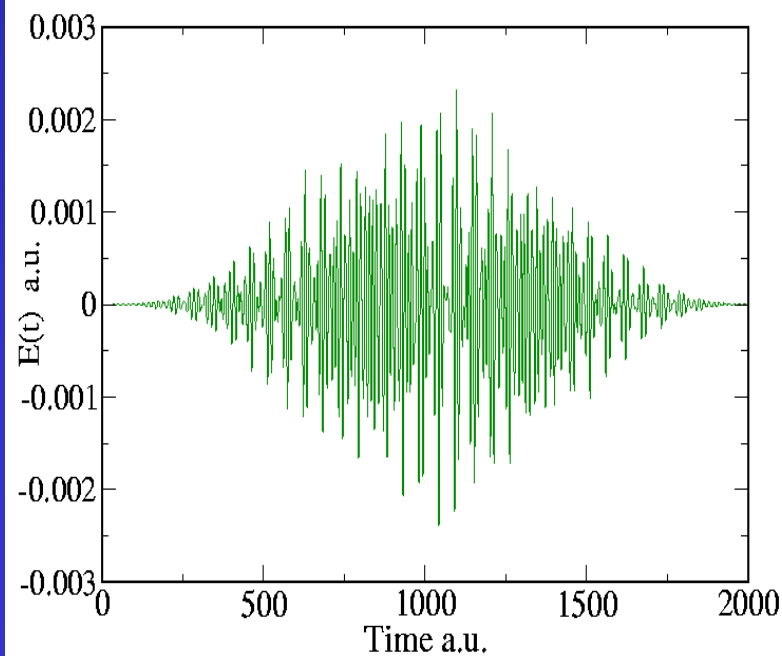
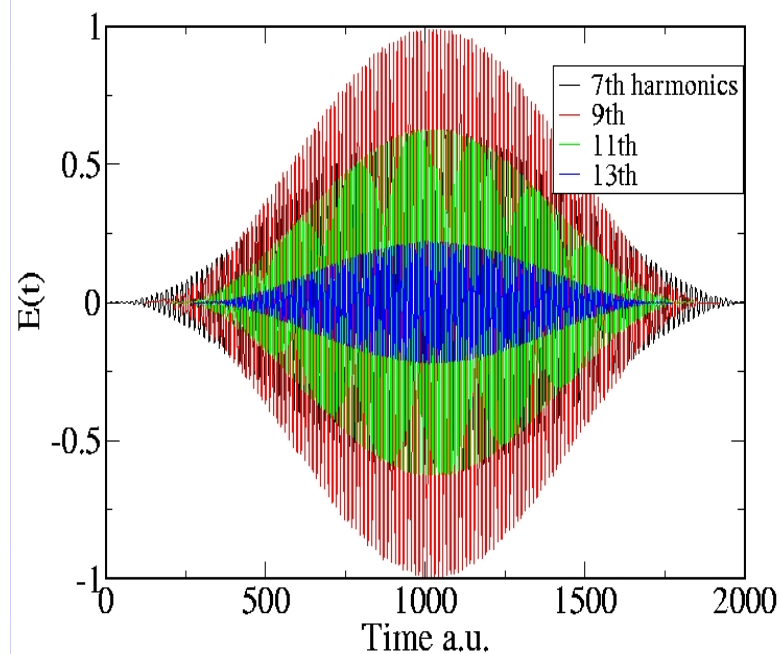
A. Scrinzi *et al.* PRA **58**, 1310 (1998)
R. Hasbani *et al.* JPB **33**, 2101 (2000)
I. F. Barna *et al.* Eur. Phys. J. D **27**, 287 (2003)



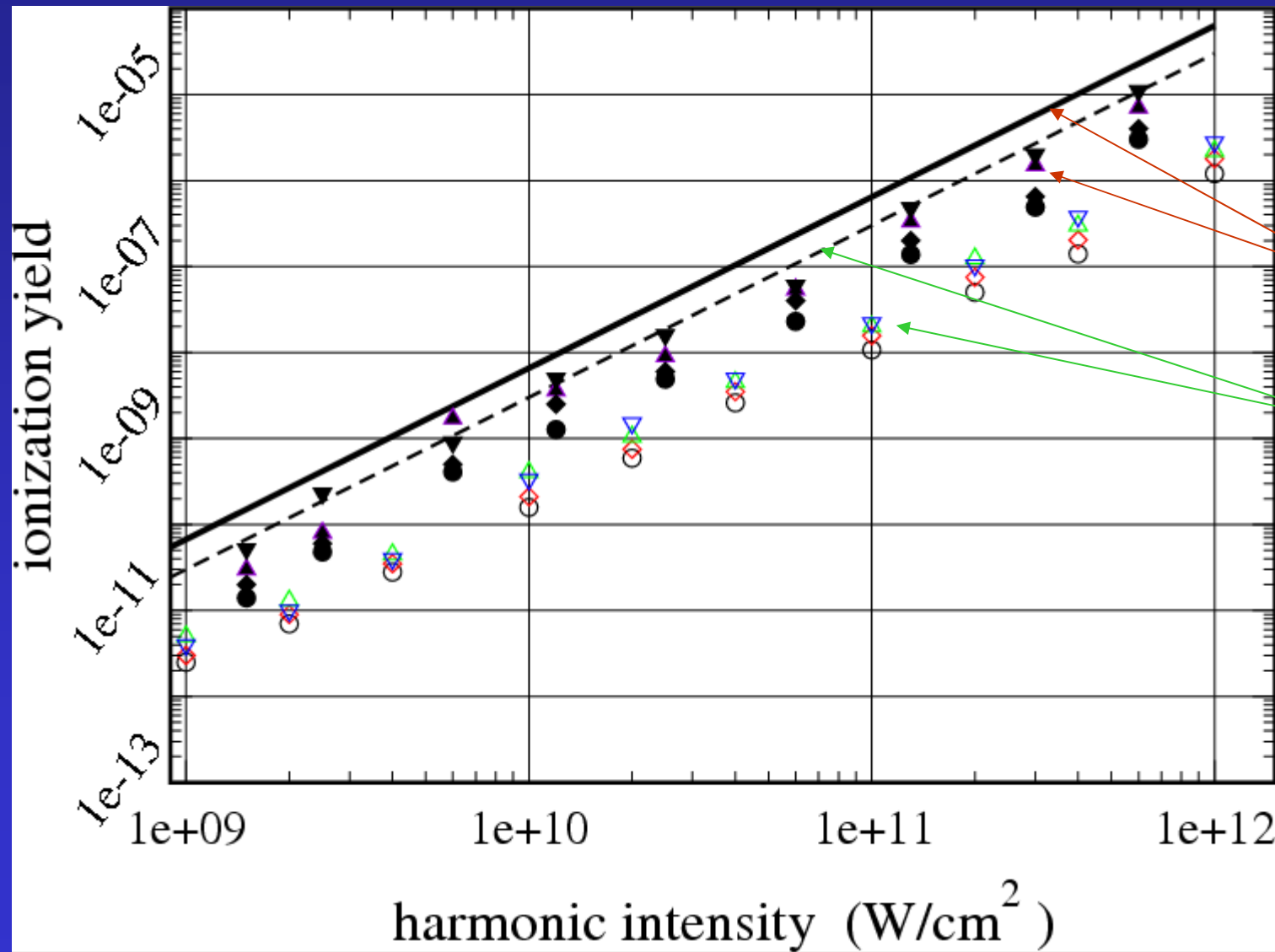
Two-Photon Ionization of He through a superposition of Higher Harmonics

Experiment & Theory (Coupled Channel method, Spline Base)

N. A. Papadogiannis *et al.* *PRL* 90 133902-1 (2003)



Results



our results

Papadogiannis
et al, PRL 2003

Summary

we presented our ab-initio coupled-channel method

showed some results for photoionization in short XUV laser pulses

showed some results for photoionization in a superposition of short XUV laser pulses

it seems to be that:

atomic orbit bases are inferior to spline bases