

Two-photon double ionisation of helium

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Outline

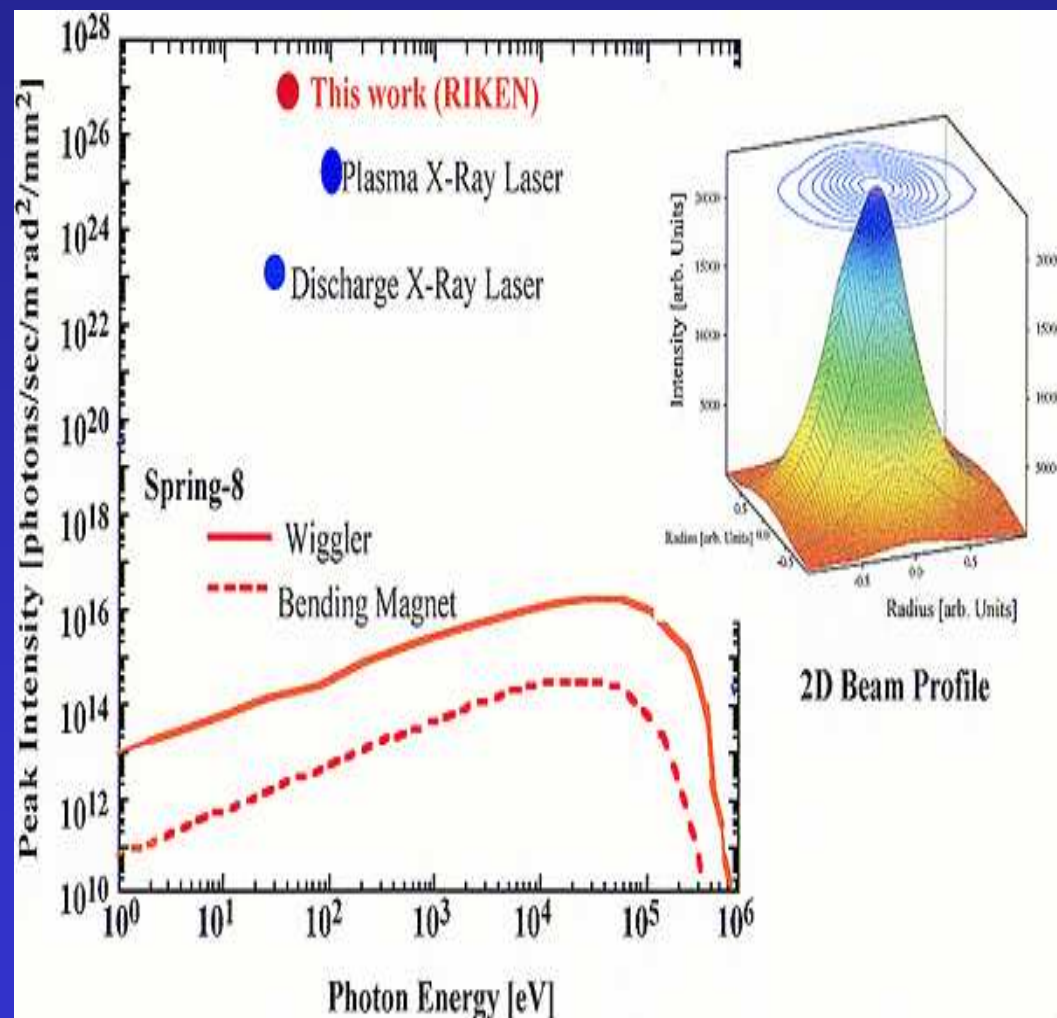
- Motivation
- The applied method
- Results for photoionisation
- Above-threshold double ionisation (ATDI)
- Summary and Outlook

Motivation

it is experimentally possible to create XUV laser pulses with HHG

- two-photon double ionisation can be studied
- electron-electron correlation can be investigated

<http://www.riken.jp/engn/rworld/research/lab/wako/laser-tech/result.htm>



The Hamiltonian of the system:

The time-dependent Schrödinger equation:

$$i\frac{\partial}{\partial t}\Psi(\vec{r}_1, \vec{r}_2, t) = (\hat{H}_{He} + \hat{V}(t))\Psi(\vec{r}_1, \vec{r}_2, t)$$

- $\hat{V}(t)$ laser-electron interaction
- $\Psi(\vec{r}_1, \vec{r}_2, t)$ Cl. wavefunction of helium

The unperturbed helium Hamiltonian:

$$\hat{H}(\vec{r}_1, \vec{r}_2)_{He} = -\frac{\nabla_1^2}{2} - \frac{\nabla_2^2}{2} - \frac{2}{r_1} - \frac{2}{r_2} + \frac{1}{|\vec{r}_1 - \vec{r}_2|}$$

- spin-spin, spin-orbit, and mass polarisation terms are neglected

The coupled-channel equation

Ansatz:

$$\Psi(\vec{r}_1, \vec{r}_2, t) = \sum_{j=1}^N a_j(t) \Phi_j(\vec{r}_1, \vec{r}_2) e^{-iE_j t}$$

Leads to a system of first-order-differential equations for the coefficients a_j :

$$\frac{da_k(t)}{dt} = -i \sum_{j=1}^N V_{kj}(t) e^{-i(E_j - E_k)t} a_j(t) \quad (k = 1 \dots N)$$

$$V_{kj}(t) = \langle \Phi_k | \hat{V}(t) | \Phi_j \rangle \quad \text{coupling matrix}$$

Initial conditions:

$$a_k(t \rightarrow -\infty) = \begin{cases} 1 & k = 1 \\ 0 & k \neq 1 \end{cases}$$

the final probability for each channel:

$$P_k(b, t \rightarrow \infty) = |a_k(t \rightarrow \infty)|^2$$

The wavefunction

Configuration interaction (CI) expansion of $\Phi_j(\vec{r}_1, \vec{r}_2)$ in terms of two-particle basis functions f_μ

$$\Phi_j(\vec{r}_1, \vec{r}_2) = \sum_{\mu} b_{\mu}^j f_{\mu}(\vec{r}_1, \vec{r}_2).$$

where $f_{\mu}(\vec{r}_1, \vec{r}_2)$ are symmetric ($S=0$) products of

1. **Slater-type orbitals:**

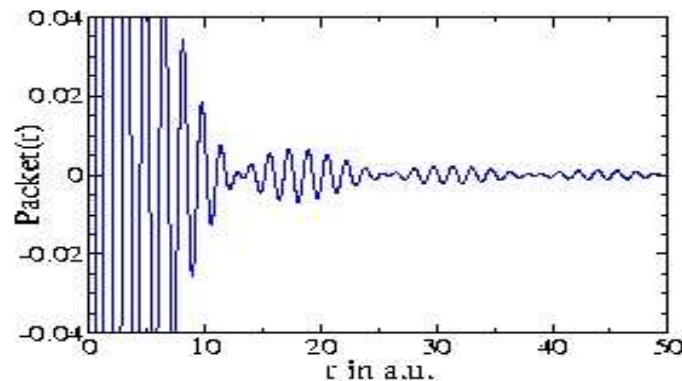
$$\chi_{n,l,m,\kappa}(\vec{r}) = C(n, \kappa) r^{n-1} e^{-\kappa r} Y_{l,m}(\theta, \varphi)$$

2. regular **Coulomb** wave packets:

$$\varphi_{k,l,m,\tilde{Z}}(\vec{r}) = N(k, \Delta k) \int_k^{k+\Delta k} R_l(\eta, \rho) dk' Y_{l,m}(\theta, \varphi)$$

- $\eta = \tilde{Z}/k'$, $\rho = k'r$, \tilde{Z} effective charge
- $N(k, \Delta k)$, $C(n, \kappa)$ normalisation constants

$Z_{\text{eff}} = 2, l = 0, k = 3.5 \text{ a.u.}, \Delta k = 0.5 \text{ a.u.}$



Angular Electron Distribution

we use the density operator to calculate the electron density after the pulse:

$$\rho(\vec{r}) = \left\langle \Psi(\vec{r}_1, \vec{r}_2, t \rightarrow \infty) \left| \sum_{i=1,2} \hat{\delta}(\vec{r} - \vec{r}_i) \right| \Psi(\vec{r}_1, \vec{r}_2, t \rightarrow \infty) \right\rangle$$

1.) we apply the Feshbach projection to extract the singly ionised contribution from the wavefunction

$$|\Psi_{ion}\rangle = (1 - \hat{P}_b - \hat{P}_{di})|\Psi(\vec{r}_1, \vec{r}_2, t \rightarrow \infty)\rangle$$

2.) integrate the radial and azimuthal angle to get the polar distribution of the ionised electron

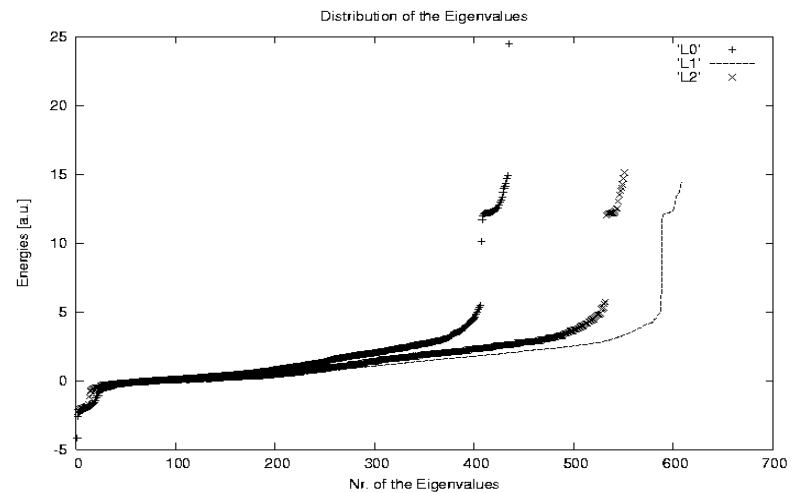
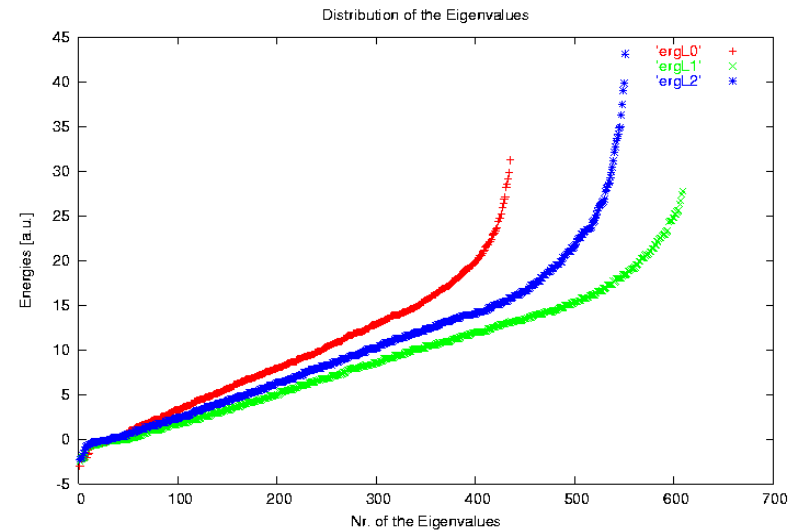
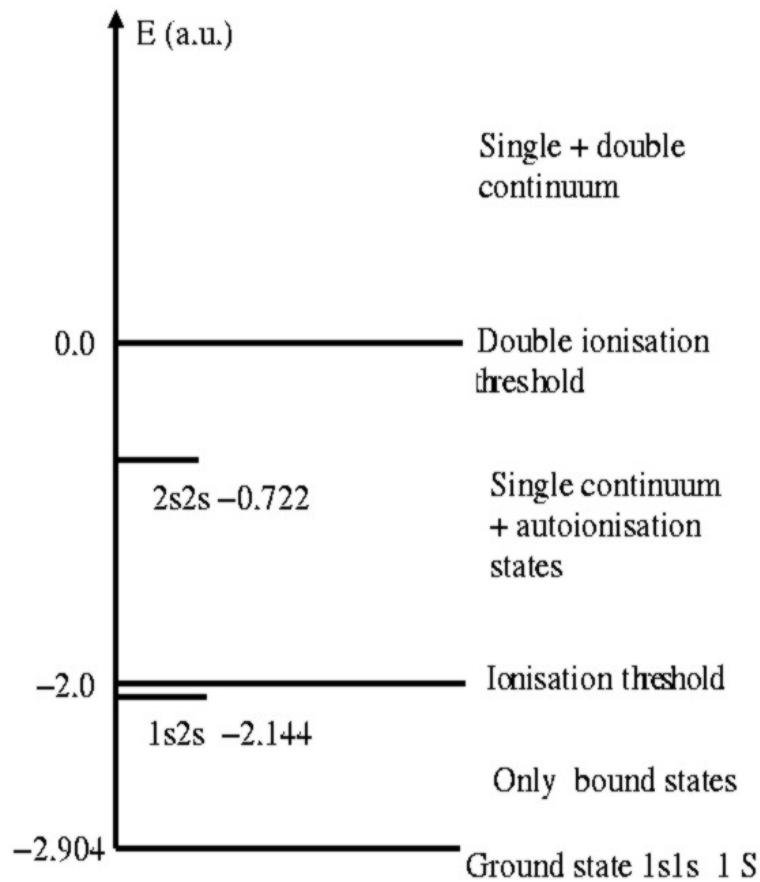
$$P(\theta) = \frac{1}{2\pi} \int_0^\infty \int_0^{2\pi} \langle \Psi_{ion} | \sum_{i=1,2} \hat{\delta}(\vec{r} - \vec{r}_i) | \Psi_{ion} \rangle r^2 dr d\varphi =$$

$$\frac{1}{\pi} \int_0^\infty \int_0^{2\pi} \int_{r_1} | \Psi_{ion}(r, \theta, \varphi, \vec{r}_1) |^2 d^3 r_1 r^2 dr d\varphi$$

the angular dependence of $P(\theta)$ is given by products of associated Legendre polynomials

Spectrum & Continuum States

Spectrum of the helium atom



Laser field

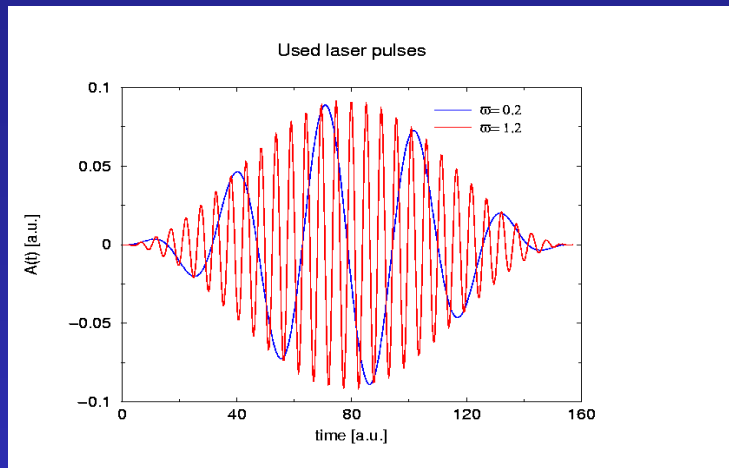
minimal coupling for the electromagnetic field to \hat{H}_{He}

$$\hat{H}(\vec{r}_1, \vec{r}_2, t) = \frac{(\vec{p}_1 - \vec{A}(\vec{r}_1, t)/c)^2}{2} + \frac{(\vec{p}_2 - \vec{A}(\vec{r}_2, t)/c)^2}{2}$$

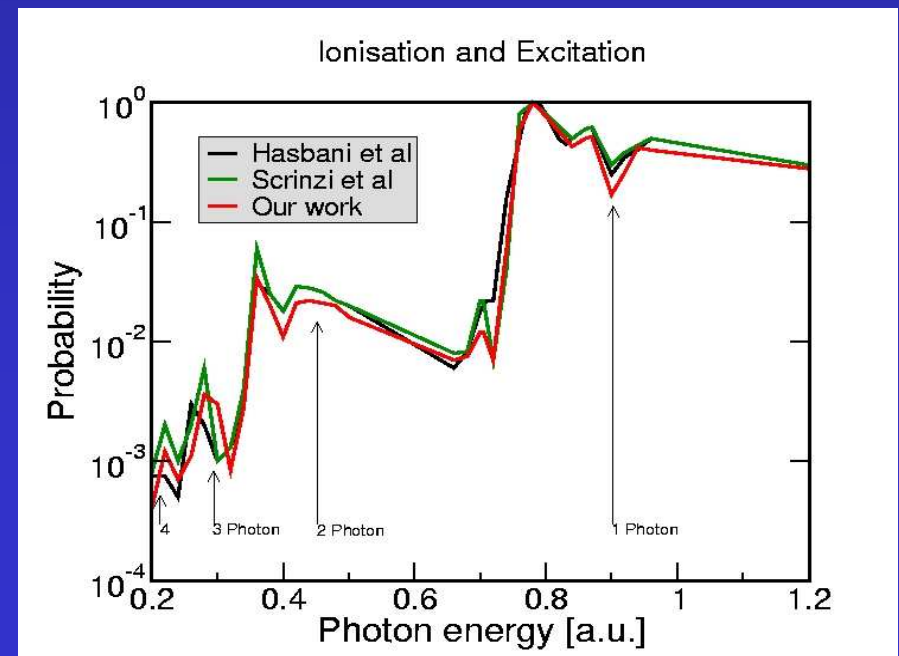
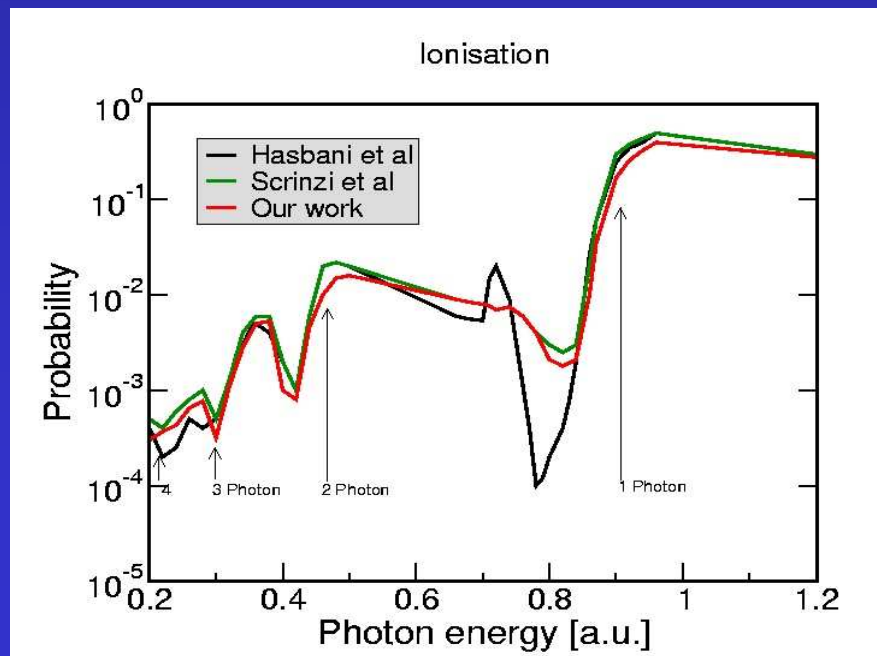
$$-\frac{2}{r_1} - \frac{2}{r_2} + \frac{1}{r_{12}} - \phi(\vec{r}_1, t) - \phi(\vec{r}_2, t) = \hat{H}_{He} + \hat{V}(t)$$

- dipole approximation: $\lambda_{laser} \gg r_{atom}$
- velocity gauge: $\hat{V}(\vec{r}, t) = \sum_{i=1,2} \vec{A}(t) \cdot \vec{p}_i$
- length gauge:
 $\hat{V}(\vec{r}, t) = - \sum_{i=1,2} \vec{E}(t) \cdot \vec{r}_i, \quad \vec{E}(t) = -\frac{1}{c} \frac{\partial}{\partial t} \vec{A}(t)$
- $\vec{E}(t) = E_0 \cdot F(t) \cdot \sin(\omega t) \cdot \vec{e}_z$ linearly polarised pulse
- envelope function: $F(t) = \sin^2\left(\frac{\pi t}{T}\right)$ T is the pulse duration
- or $F(t) = \exp\left(-\frac{(2\ln 2)t^2}{T^2}\right)$ T is the FWHM

Laser pulse & Photoionisation Results 1



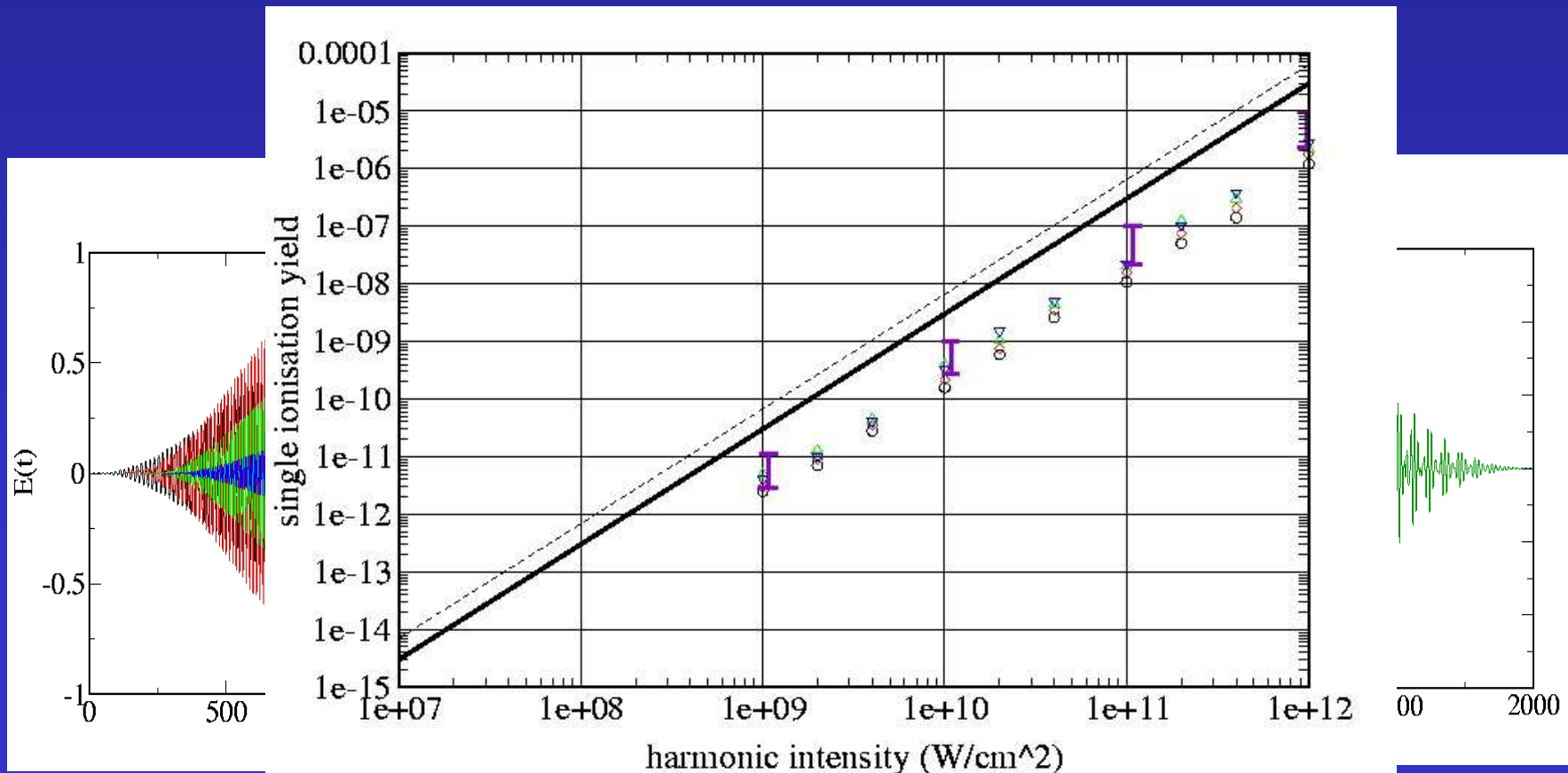
A. Scrinzi *et al.* PRA **58**, 1310 (1998)
R. Hasbani *et al.* JPB **33**, 2101 (2000)
I. F. Barna *et al.* Eur. Phys. J. D **27**, 287 (2003)



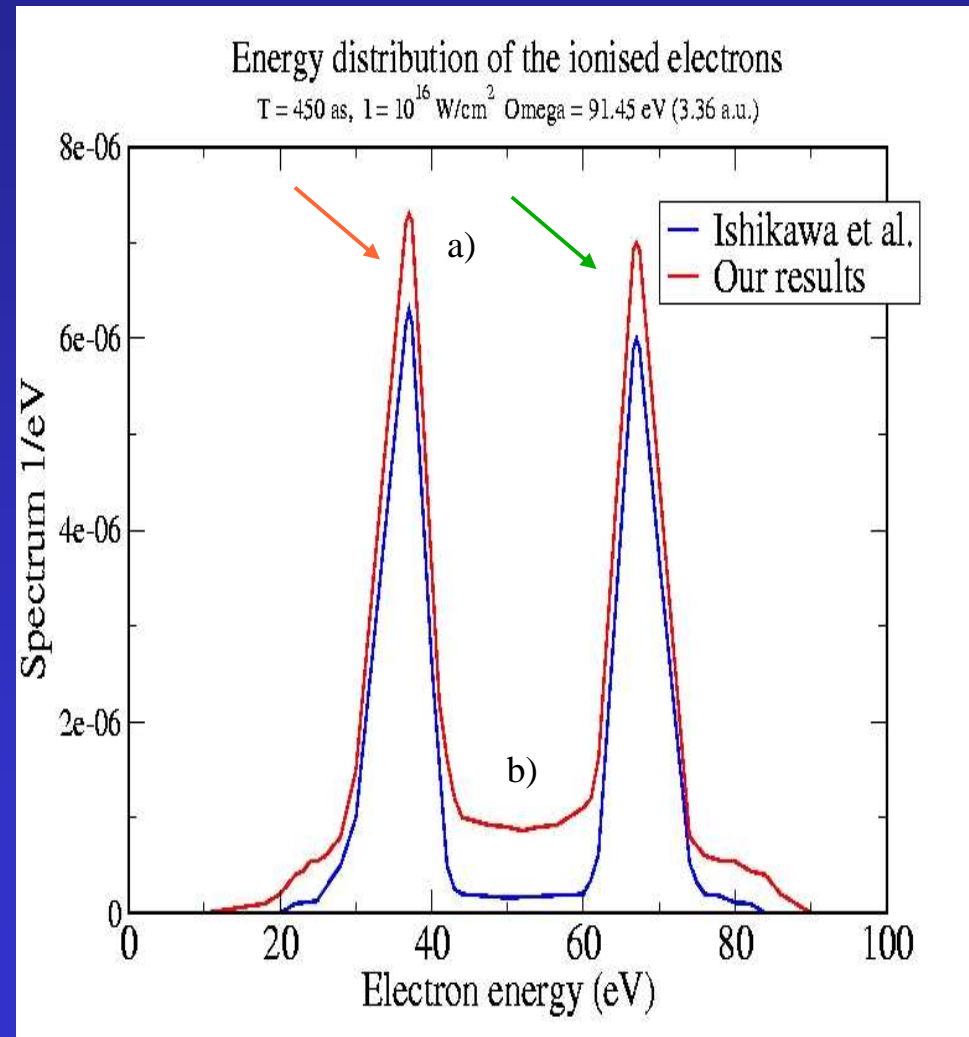
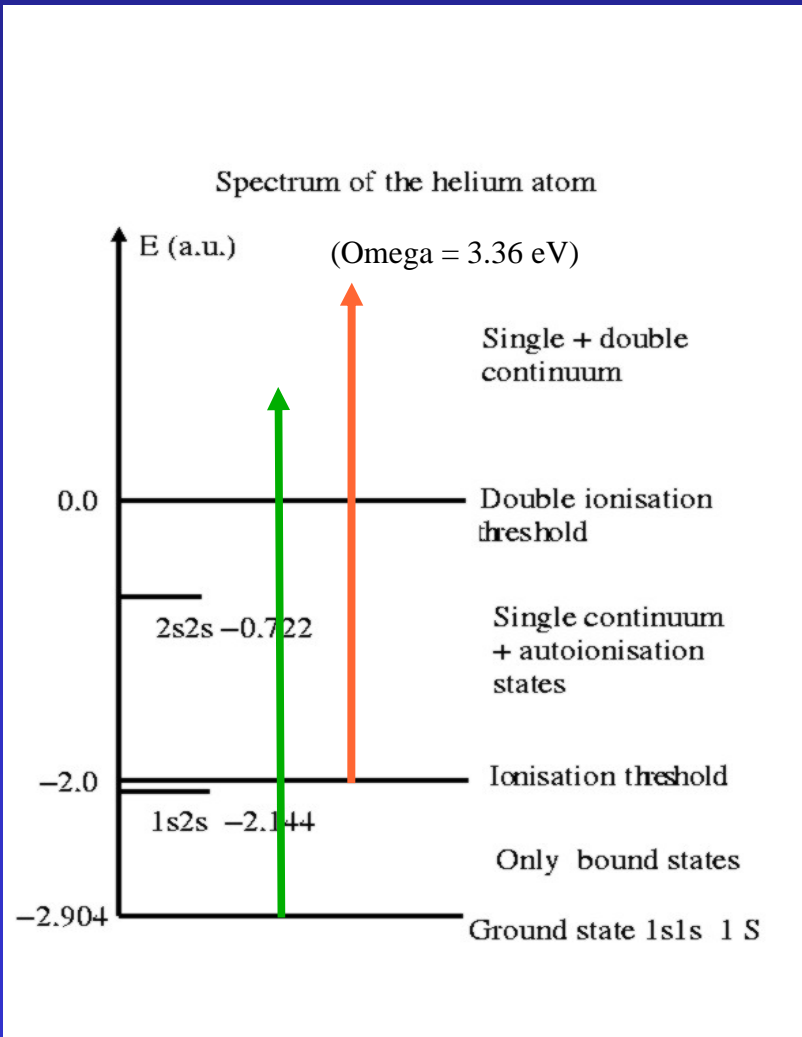
Photoionisation Results 2

Two-Photon Ionisation of He through a superposition of Higher Harmonics

N. A. Papadogiannis *et al.* *PRL* 90 133902-1 (2003)



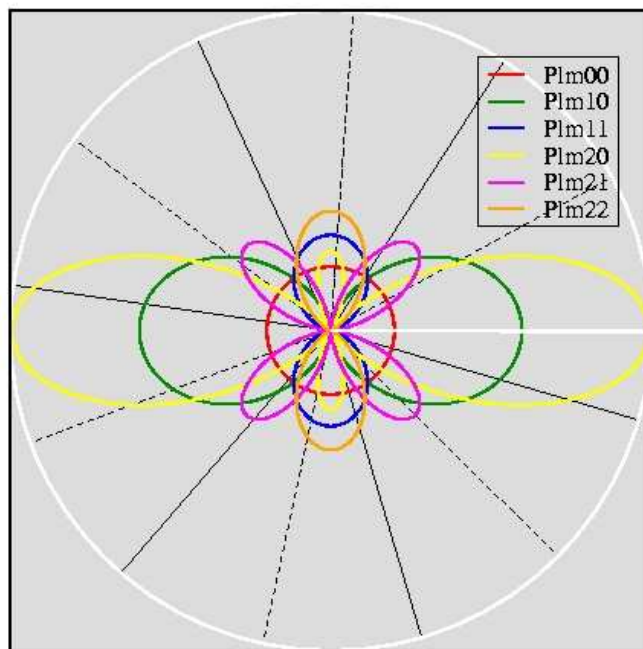
Above-threshold double ionisation (ATDI)



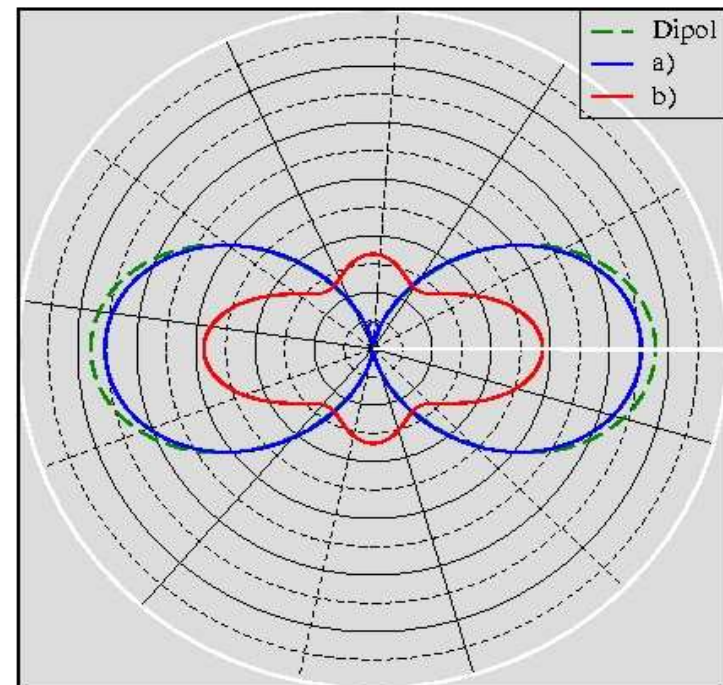
Angular distribution of the ejected electrons

$L = 0$ ss+pp+dd, $L = 1$ sp+pd, $L = 2$ sd+pp+dd

Associated Legendre Polynomials



Results



Summary and Outlook

we presented our CC method for photoionisation

showed some results for photoionisation

presented the angular distribution of the ejected electrons in above-threshold double ionisation

further work is in progress to clear out the dark points

*Thank you for
your attention!*

