

# *Epitaxial Fe/Si/Fe(001) Structure and Magnetism of a Unique System*

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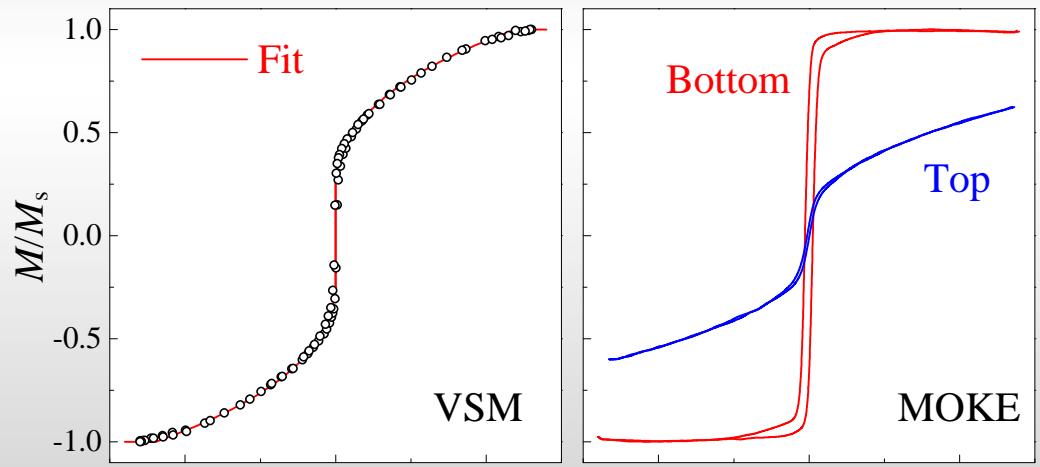
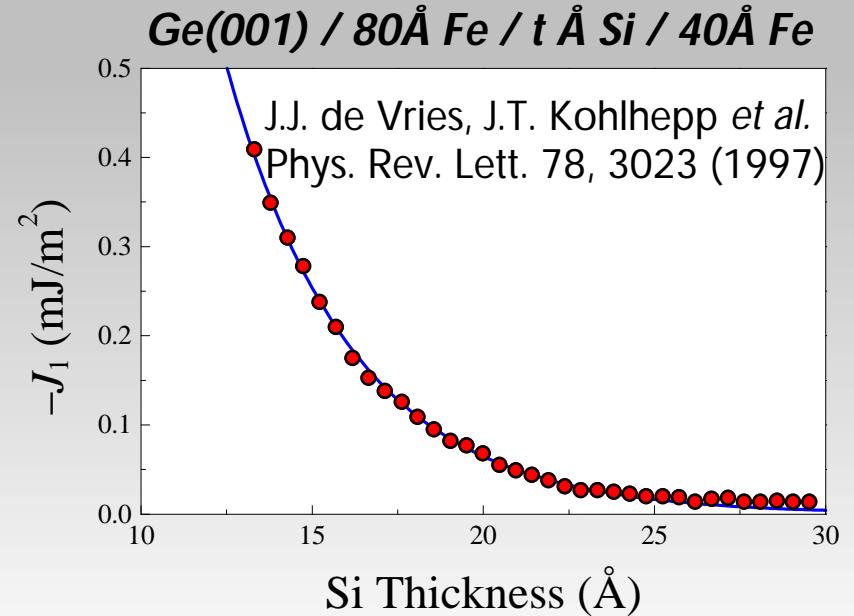
*Delft University of Technology, Interfacultair Reactor Instituut,  
Mekelweg 15, 2629 JB Delft, The Netherlands*

# Motivation (1)

- exploring new materials, combining ferromagnetic metals (Fe, Co, ...) and traditional semiconducting materials (Si, Ge, ...)
- new type of interlayer exchange coupling
- investigation of iron-silicide formation in well-defined systems

# Motivation (2)

- **bilinear (antiferromagnetic) coupling** in epitaxial Fe/Si/Fe(001) sandwiches and textured Si/Fe(110) multilayers observed
- formation of c-FeSi in the spacer
  - special band structure and density of states features
  - Fermi - surface of c-FeSi
- additionally, observation of a strong **biquadratic (orthogonal) coupling** however, origin was still unknown
- Fe/Si multilayers are not suitable for studying biquadratic coupling

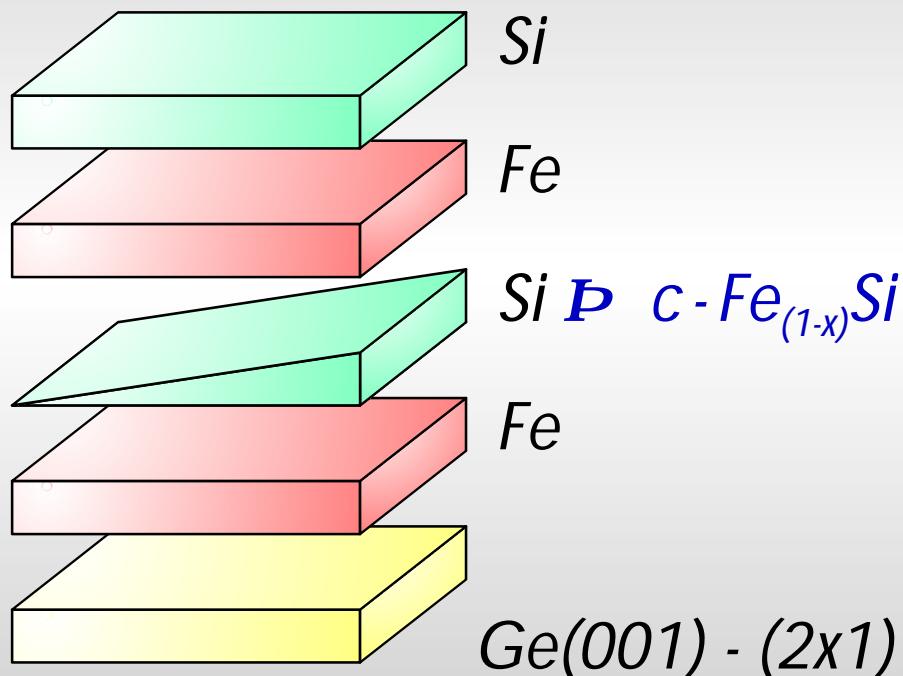


J.T. Kohlhepp et al.  
Phys. Rev. B 55, R696 (1997)

# Experimental

## *preparation*

MBE - grown sandwiches



## *characterization*

*in-situ:*

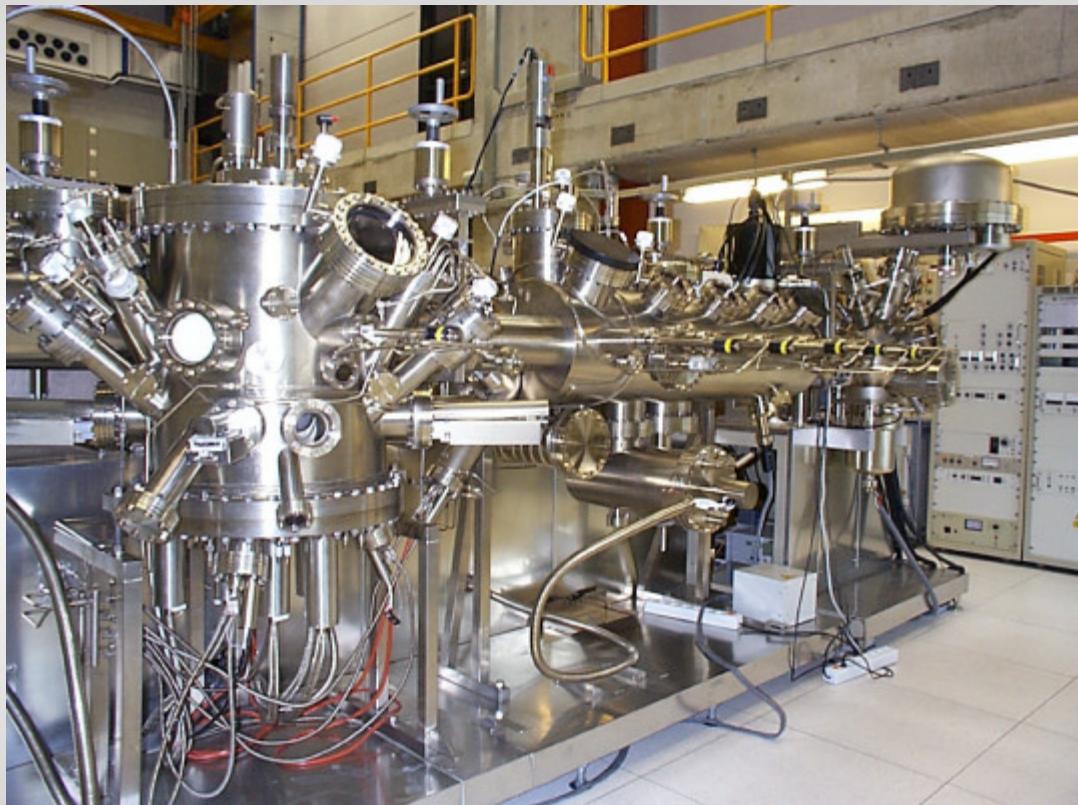
- AES
- XPS
- LEED
- STM

*ex-situ:*

- temperature dependent MOKE
- SQUID
- Mössbauer spectroscopy (CEMS)
- XRD
- Polarized Neutron Reflectometry (ISIS, Rutherford Appleton Labs)

# *VG Semicon V80M MBE system*

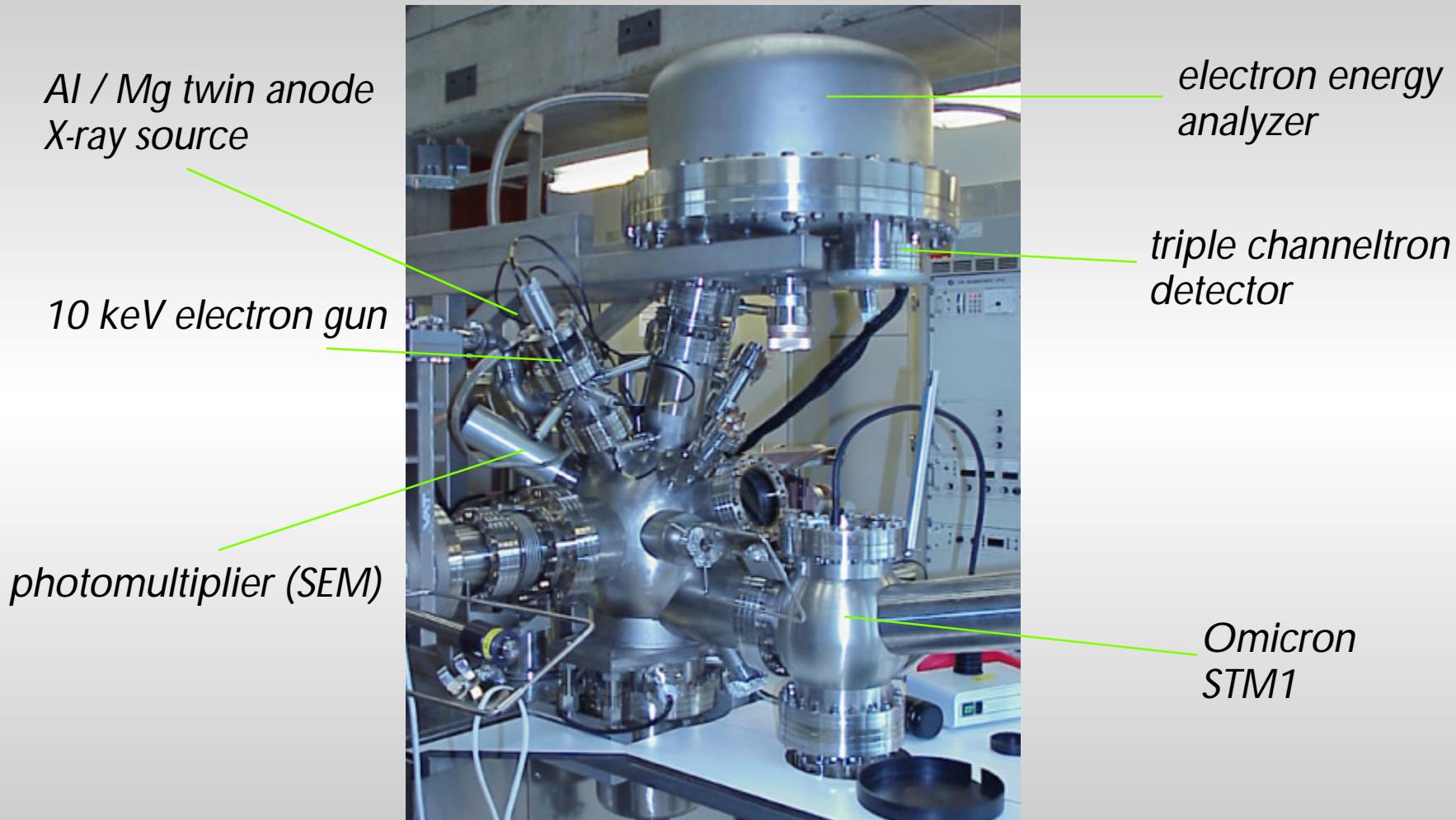
## *features:*



- base pressure  $< 2 \times 10^{-11}$  mbar
- separate growth, preparation and analysis chambers
- deposition chamber:
  - 3 e-guns, 4 Knudsen cells
  - quartz crystal monitors
  - RHEED
  - variable temperature 240-1100 K
- preparation chamber:
  - sputter cleaning (300-1100 K)
  - LEED
  - fast entry load lock

# VG Semicon V80M MBE system

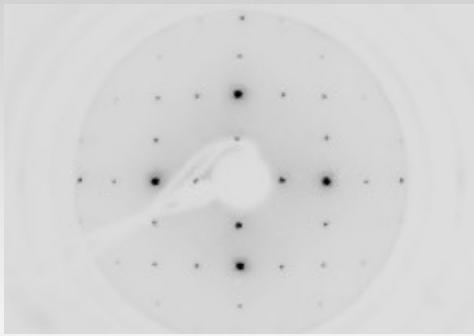
- surface analysis chamber (ESCA - lab) -



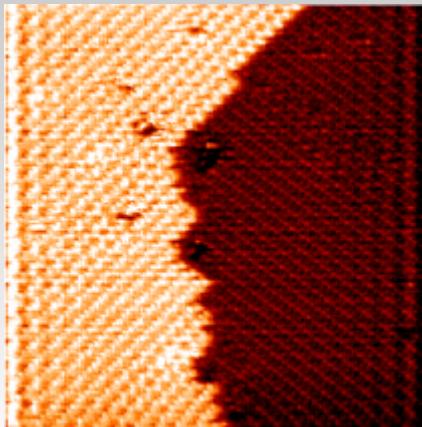
# Properties of Fe on Ge(001)

Ge(001)  
 $2 \times 1$

LEED

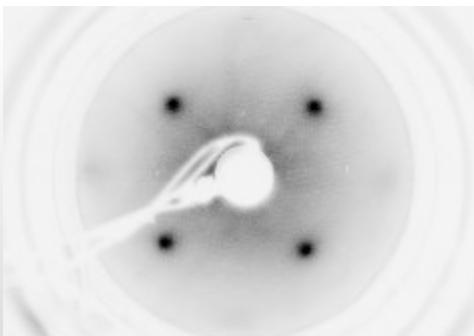


STM



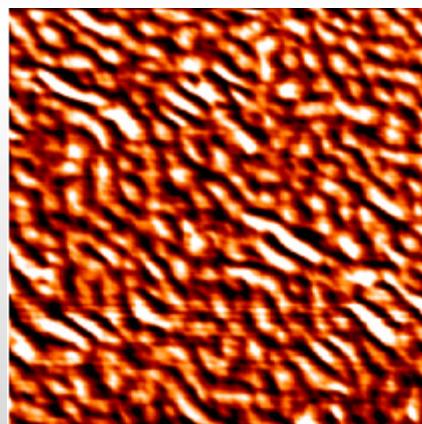
$25 \times 25 \text{ nm}^2$ , atomically flat

+ 60 Å Fe

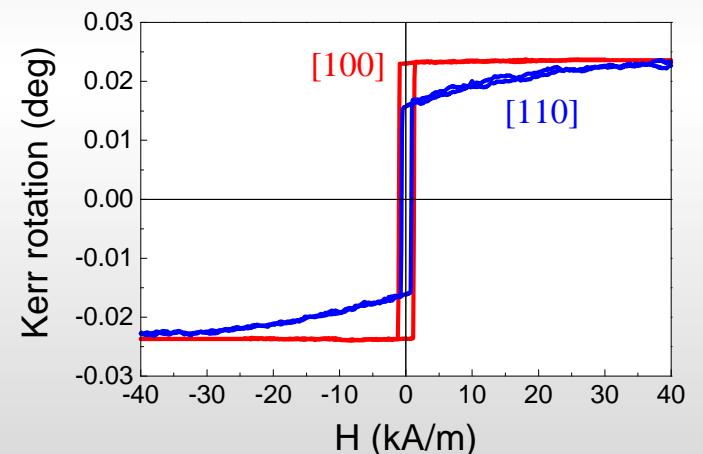


bcc(001) growth

$100 \times 100 \text{ nm}^2$   
0.24 nm rms roughness



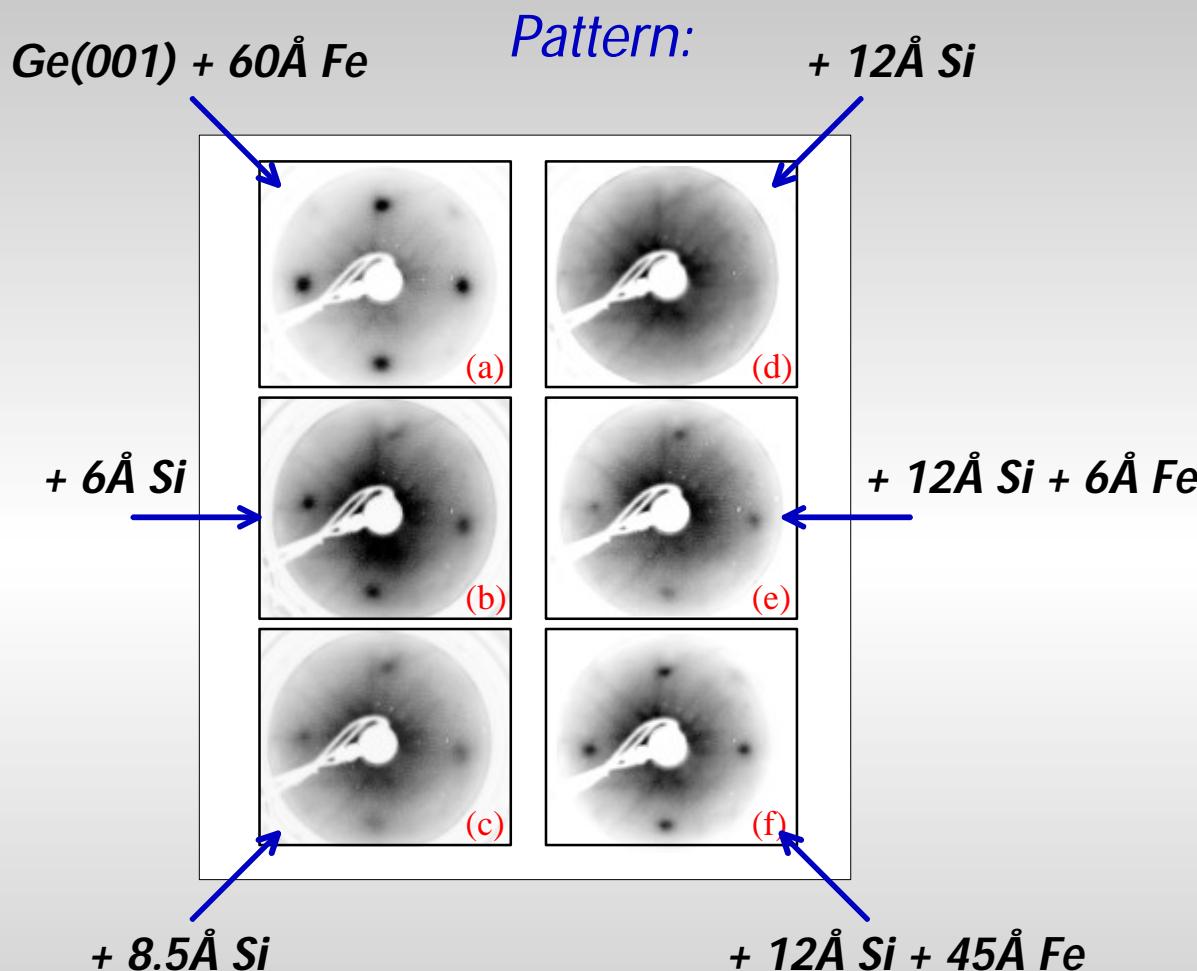
MOKE



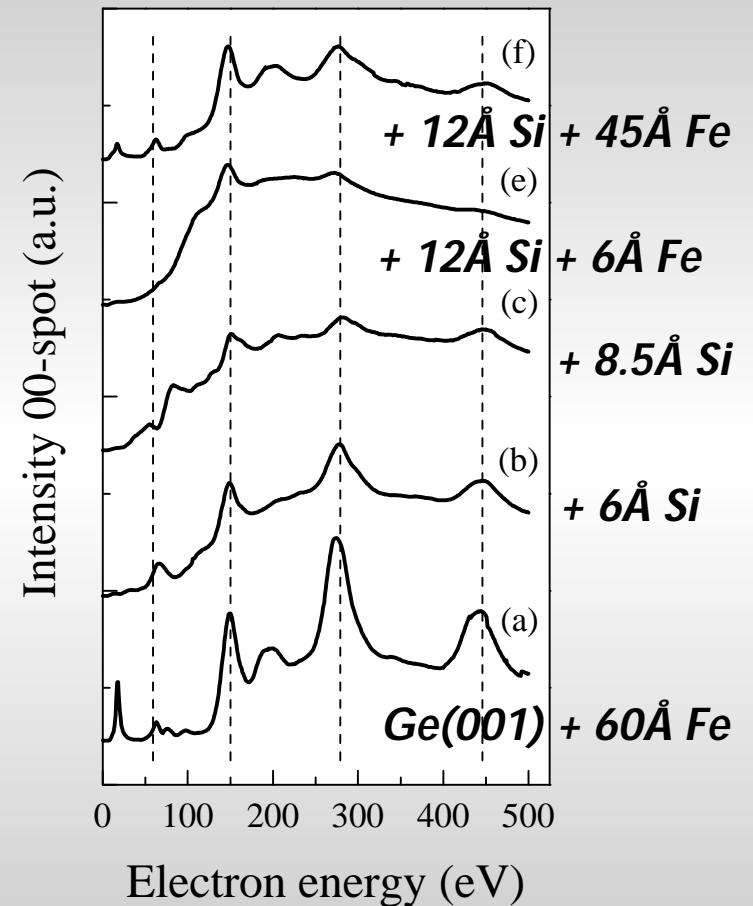
small coercivity  
fourfold cubic anisotropy

# Structure of Ge(001)/Fe/Si/Fe

## LEED studies:

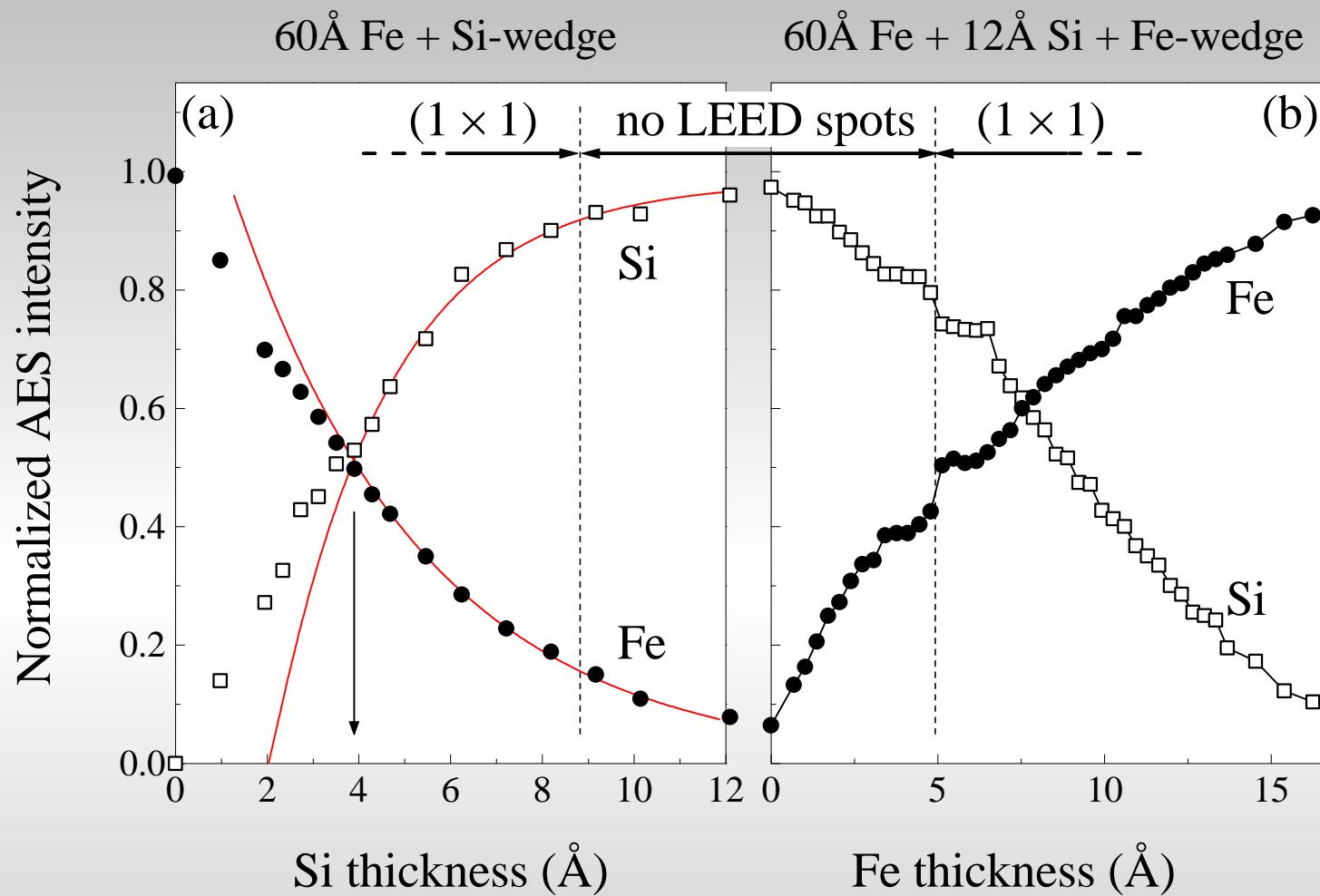


## I-V spectra:



P bcc-like (001)-structure maintained in stack

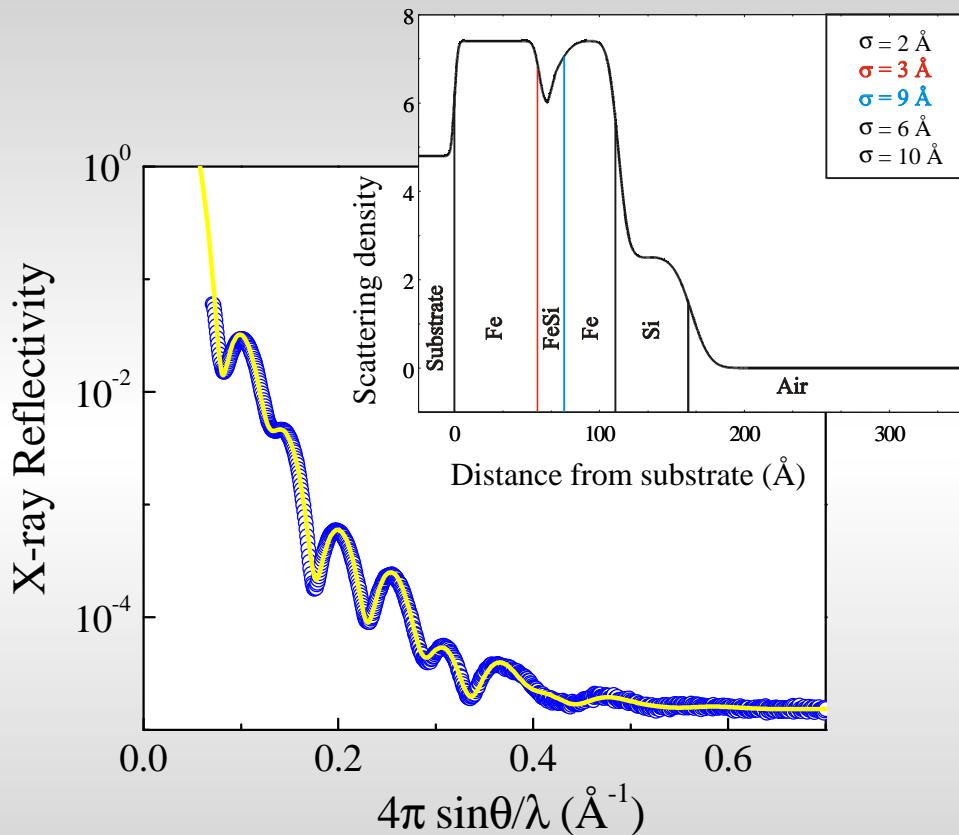
## AES studies:



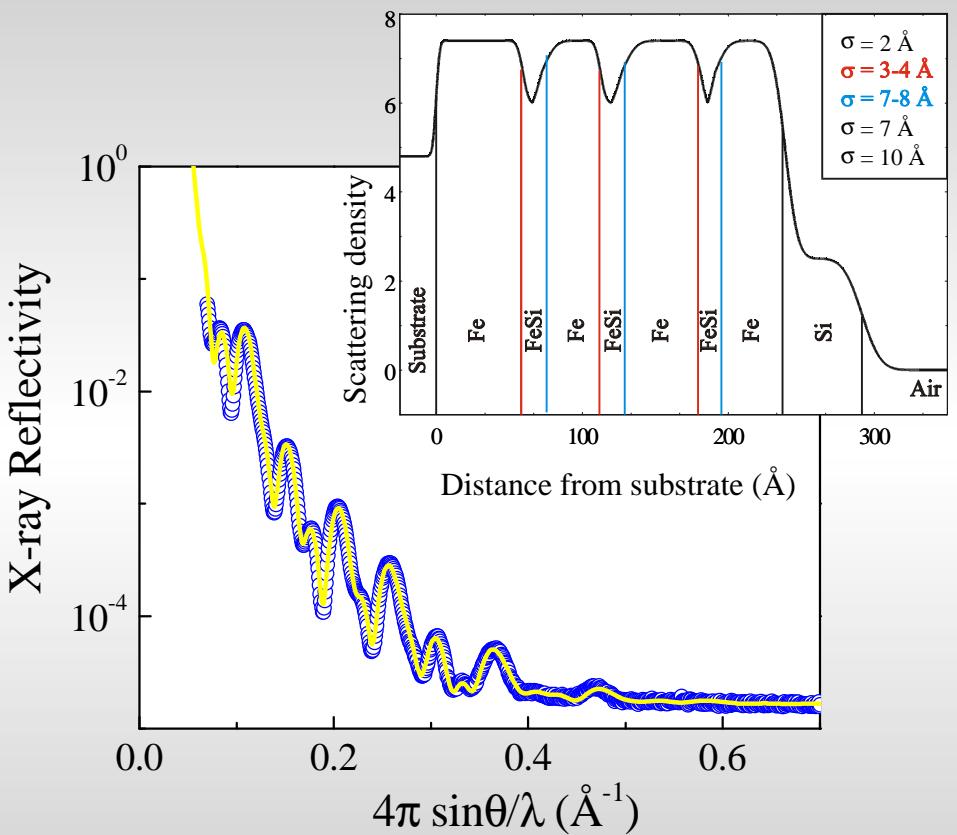
*P Fe diffuses from the bottom and top into the Si spacer  
accompanied by a reappearance of LEED spots*

## GIXR studies:

$\text{Ge}(001) / \text{60 \AA Fe} / 14 \text{ \AA Si} / \text{45 \AA Fe} / 40 \text{\AA Si}$

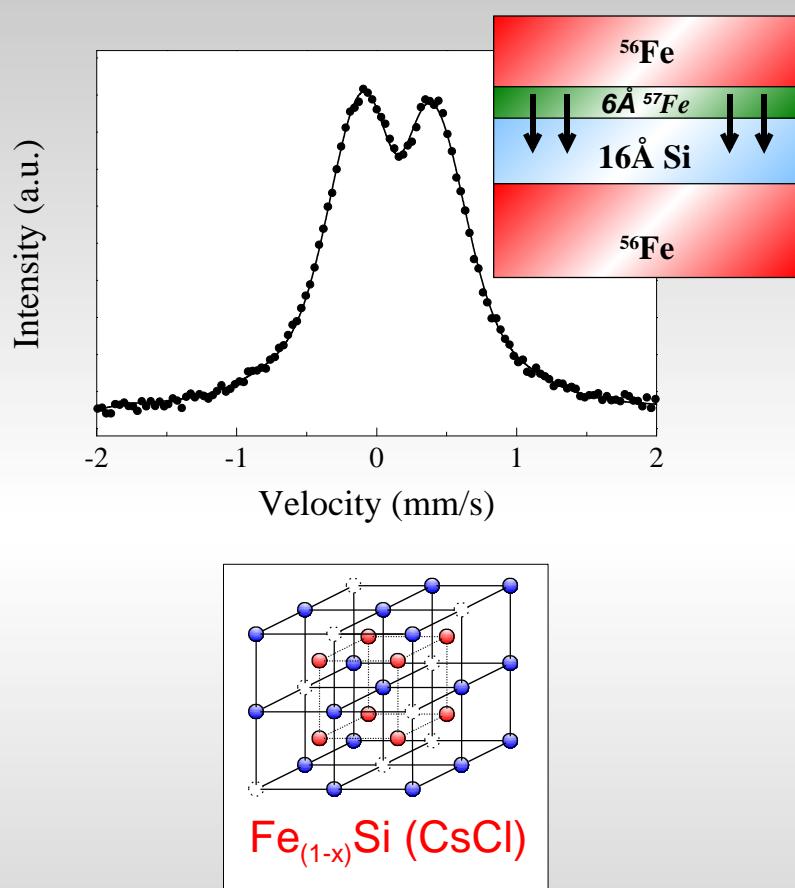


$\text{Ge}(001) / (\text{60 \AA Fe} / 14 \text{ \AA Si} / \text{45 \AA Fe} / 14 \text{\AA Si})_2 / 26 \text{\AA Si}$

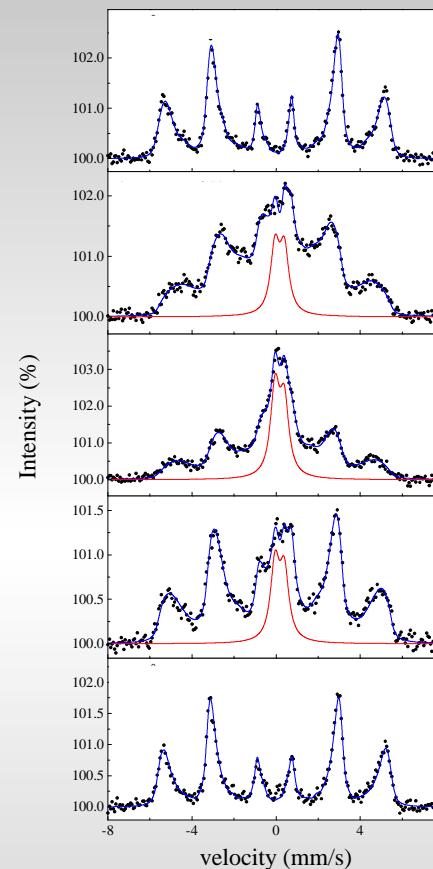


P *Si spacer completely gone; transformed to FeSi*

## CEMS studies:

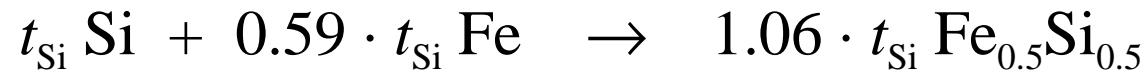
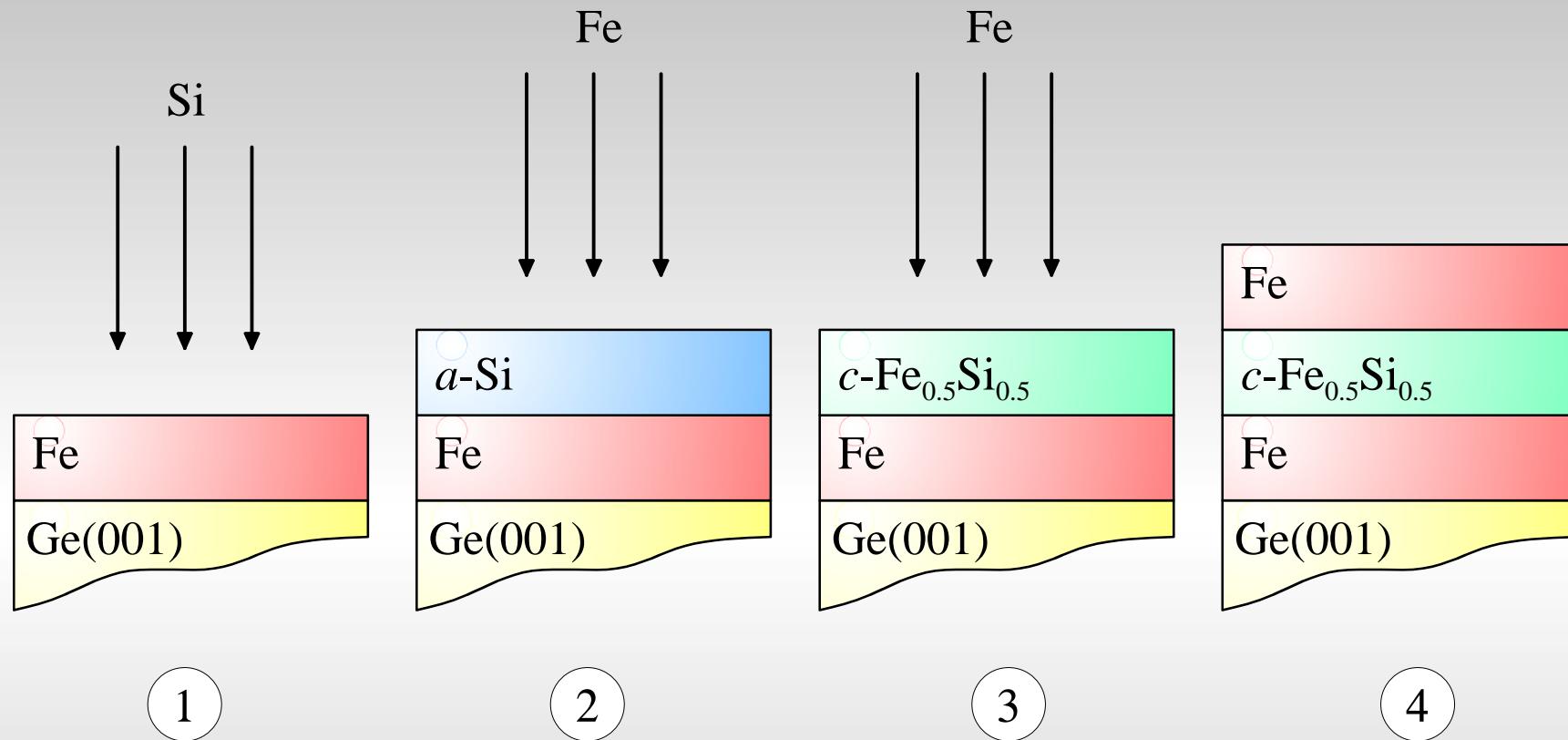


G.J. Strijkers, J.T. Kohlhepp et al.  
Phys. Rev. B 60, 9583 (1999)



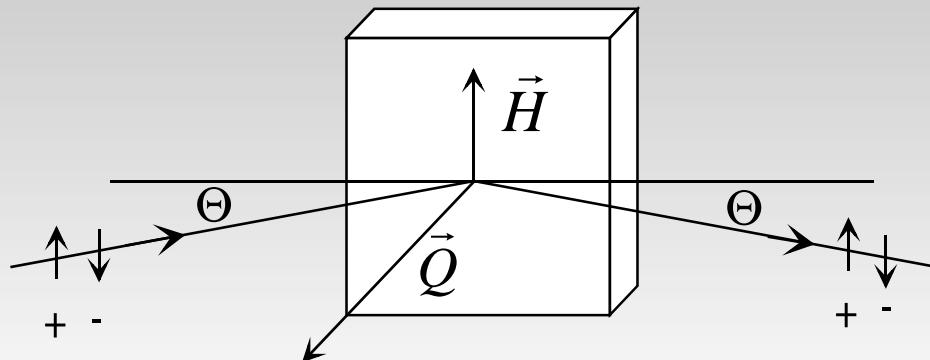
*P c- $\text{Fe}_{(1-x)}\text{Si}$  with metastable CsCl (B2) structure and  $x \gg 0.36$  is formed in the spacer*

# Summary: Iron-silicide formation (simplified)



# Magnetic Properties

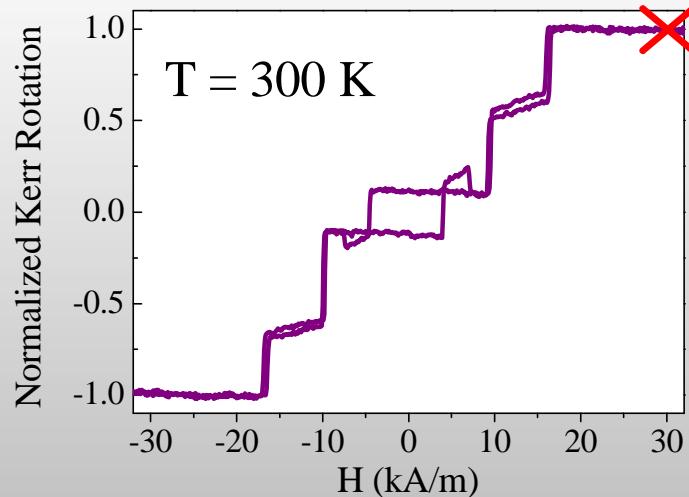
## Spin Polarized Neutron Reflectometry:



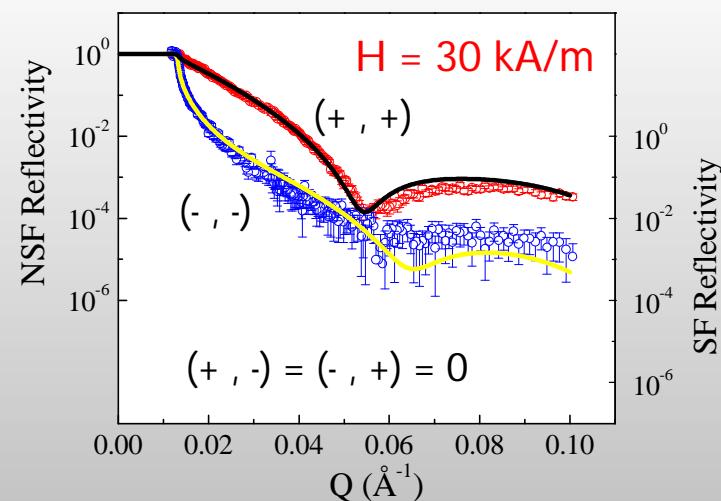
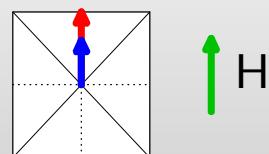
$$|\vec{Q}| = \frac{4p}{I} \sin \Theta$$

(+, +); (-, -) : Non Spin Flip (NSF) Reflectivity  
(+, -); (-, +) : Spin Flip (SF) Reflectivity

Ge(001) / 60 Å Fe / 14 Å Si / 45 Å Fe / 40 Å Si

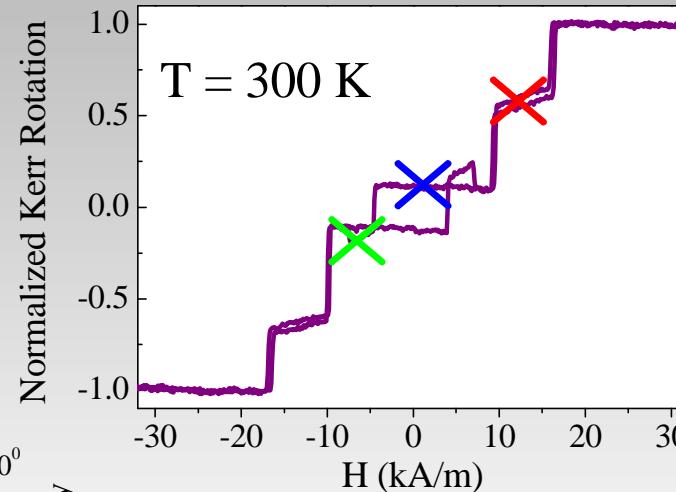
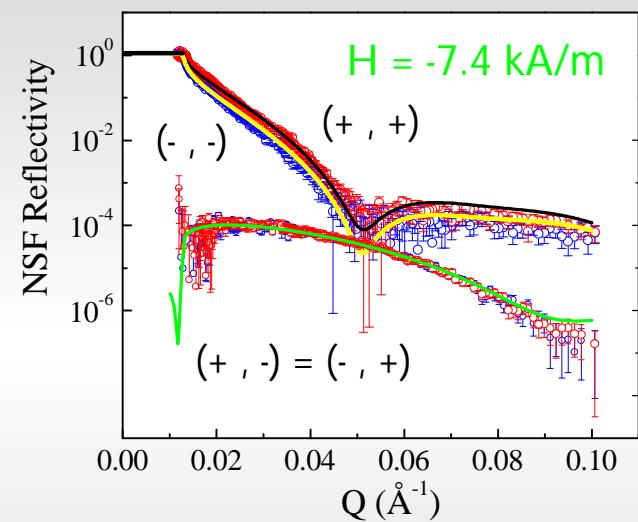
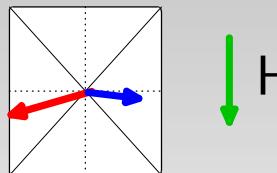


Saturation:

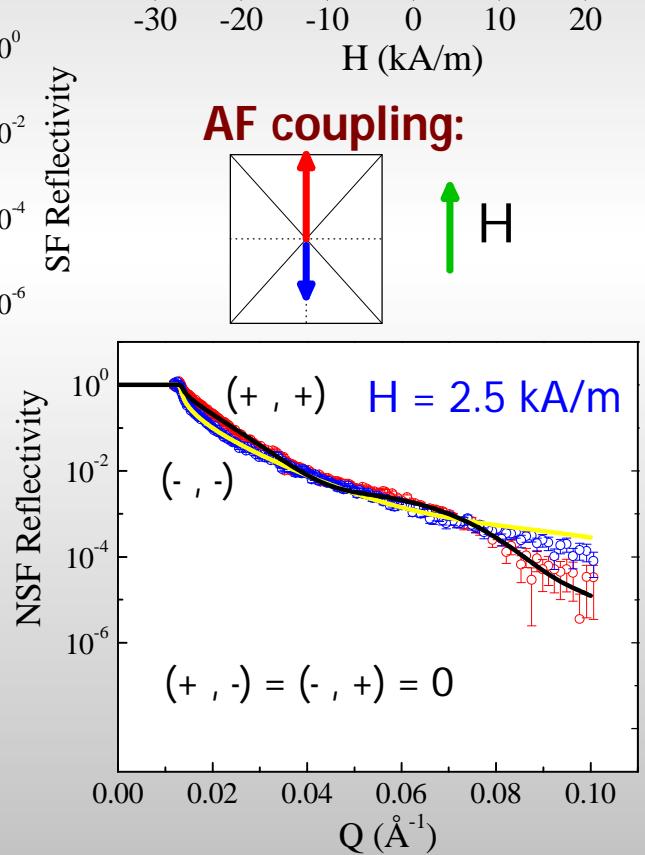
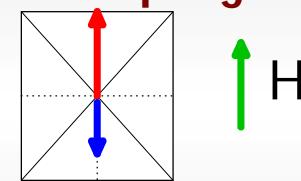


# Magnetization reversal details at RT:

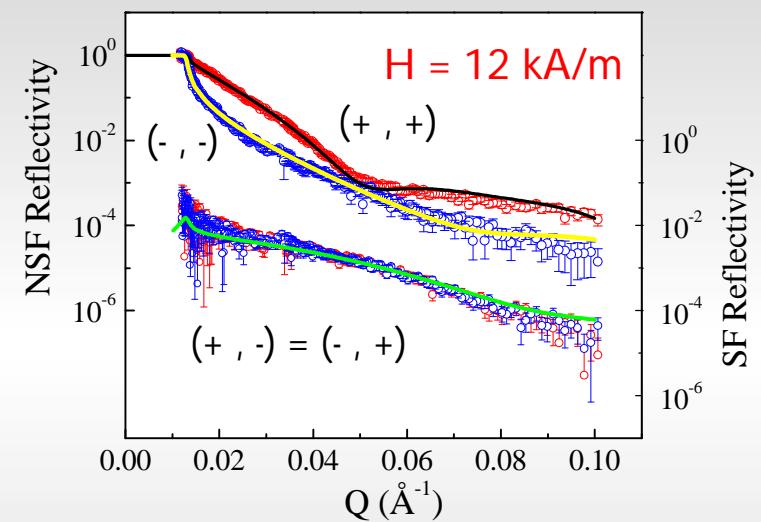
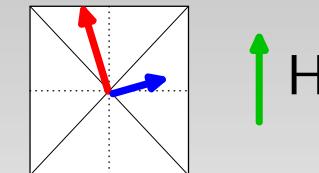
Spin flop:



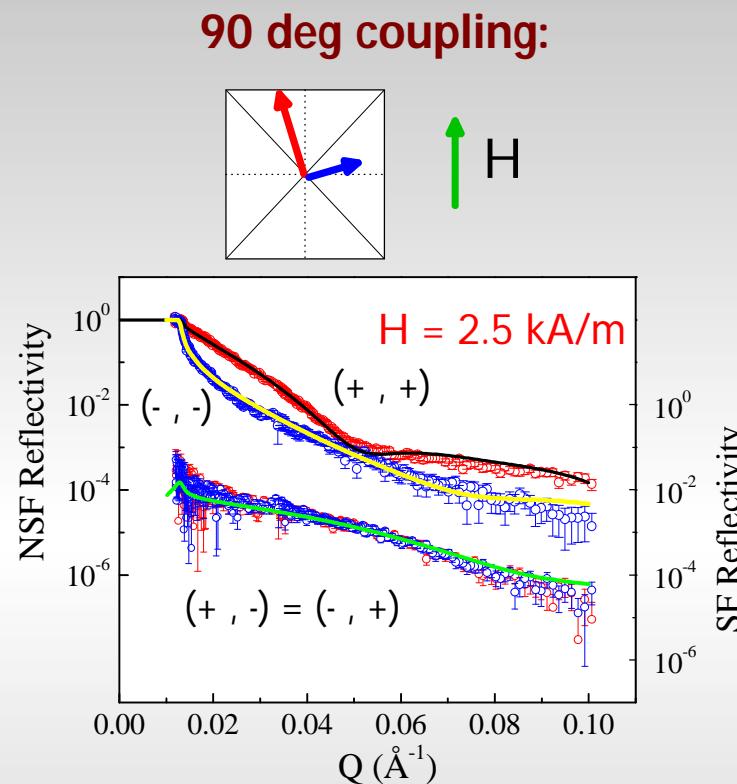
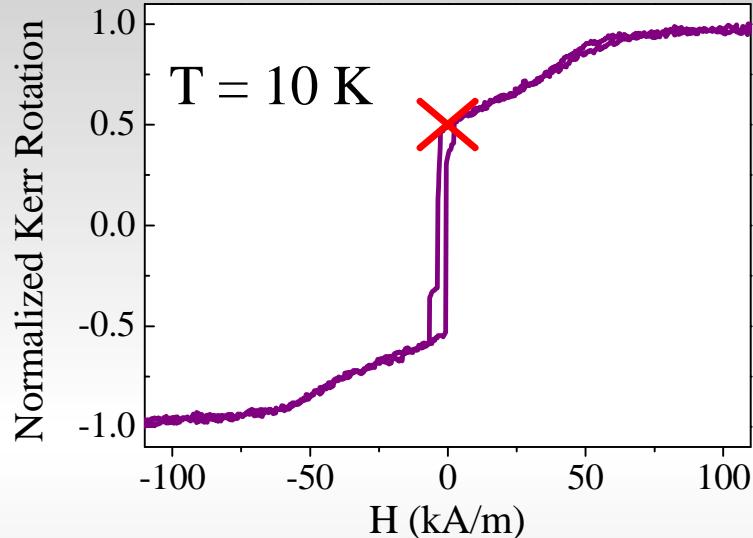
AF coupling:



90 deg coupling:



## *Magnetization reversal details at low T:*

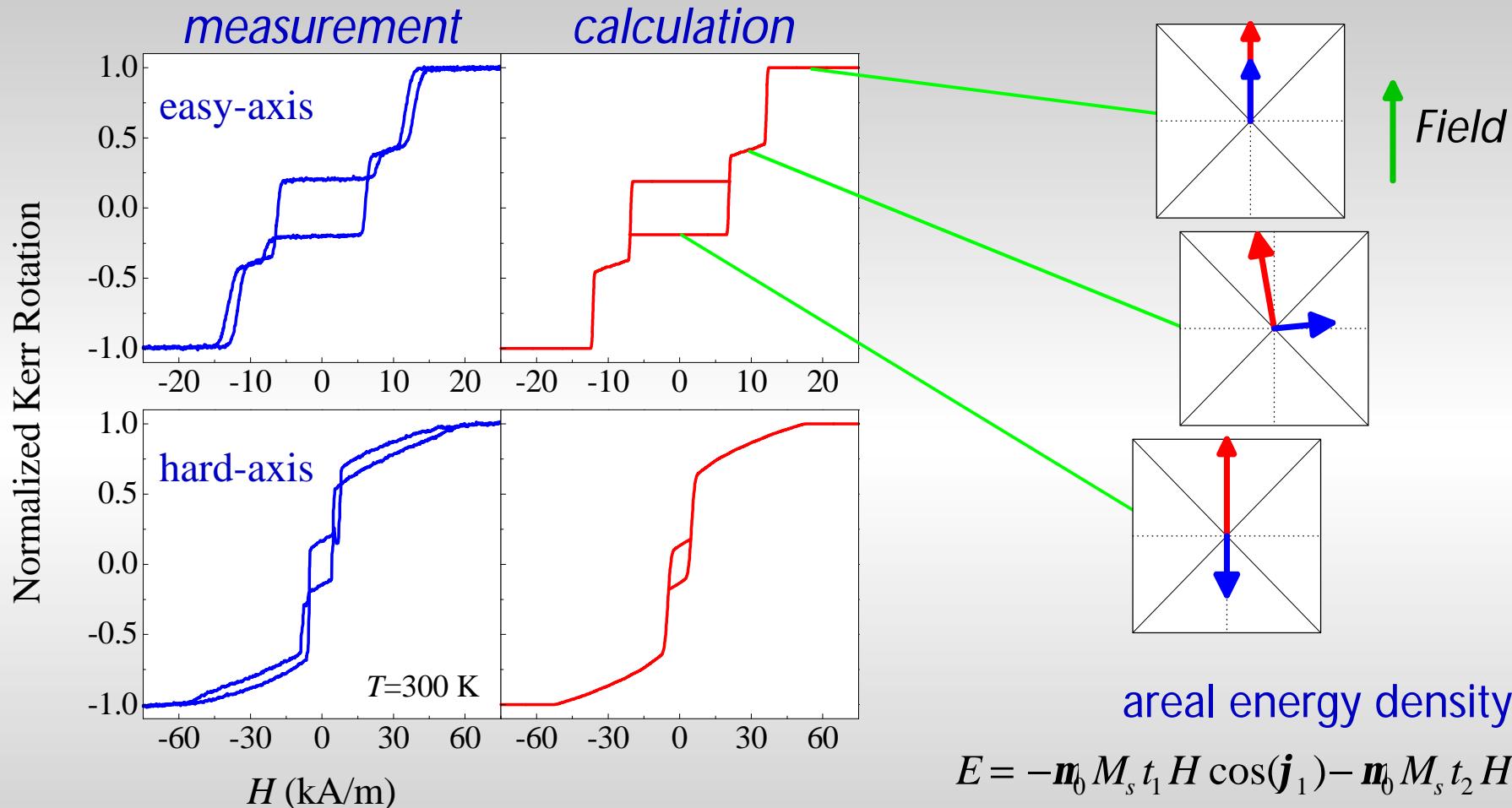


*P biquadratic coupling dominates at low T*

R.W.E. van de Kruijs, J.T. Kohlhepp *et al.*  
Phys. Rev. B (2002), accepted

# Magnetization loops and simulation (1)

Ge(001) / 115Å Fe / 13.7 Å Si / 90Å Fe / 30Å Si

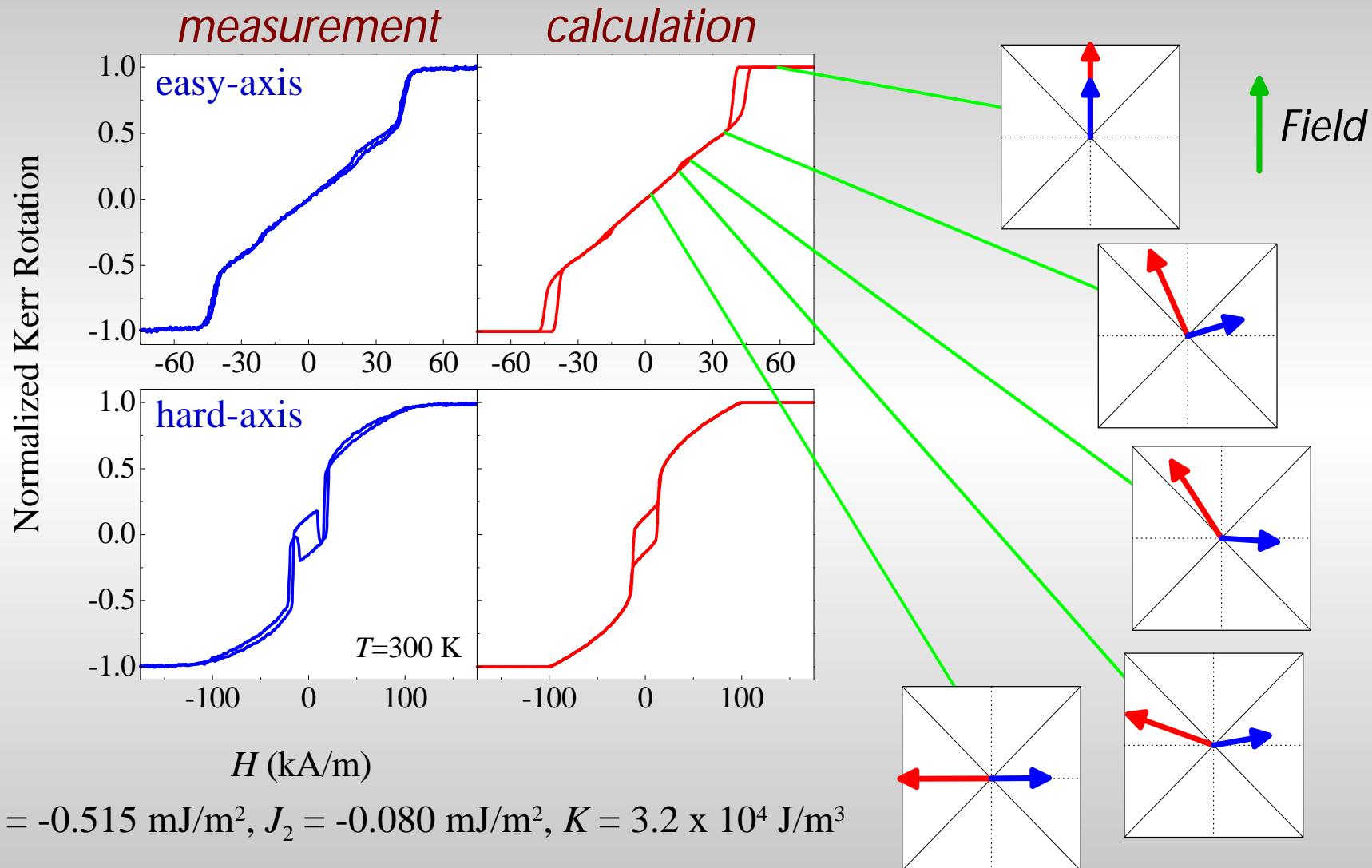


$$\begin{aligned} E = & -\mathbf{n}_0 M_s t_1 H \cos(\mathbf{j}_1) - \mathbf{n}_0 M_s t_2 H \cos(\mathbf{j}_2) \\ & + K t_1 \cos^2(\mathbf{j}_1) \sin^2(\mathbf{j}_1) + K t_2 \cos^2(\mathbf{j}_2) \sin^2(\mathbf{j}_2) \\ & - J_1 \cos(\mathbf{j}_1 - \mathbf{j}_2) + J_2 \cos^2(\mathbf{j}_1 - \mathbf{j}_2) \end{aligned}$$

$$J_1 = -0.160 \text{ mJ/m}^2, J_2 = -0.028 \text{ mJ/m}^2, K = 3.2 \times 10^4 \text{ J/m}^3$$

## Magnetization loops and simulation (2)

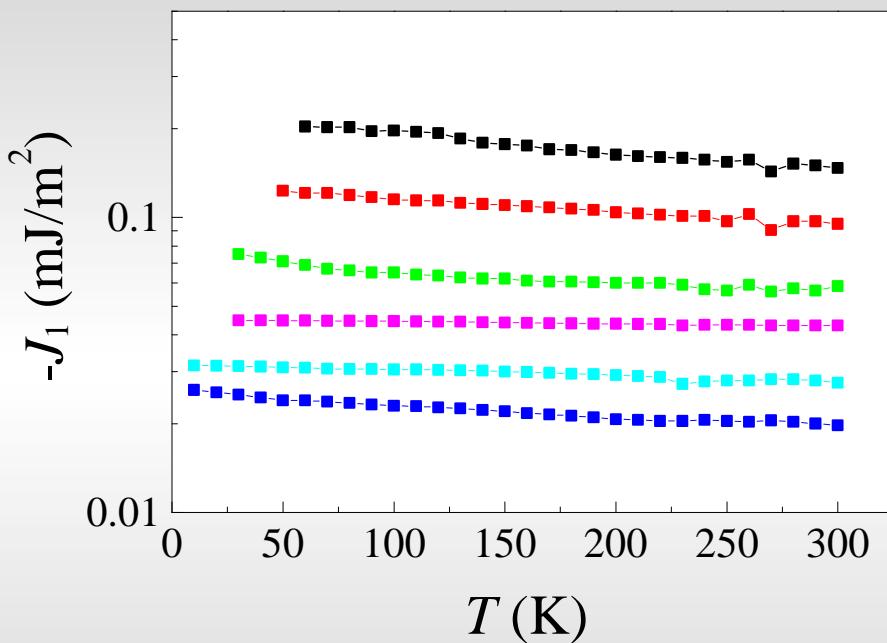
Ge(001) / 115 Å Fe / 12.4 Å Si / 90 Å Fe / 30 Å Si



# *Temperature dependence of coupling strengths*

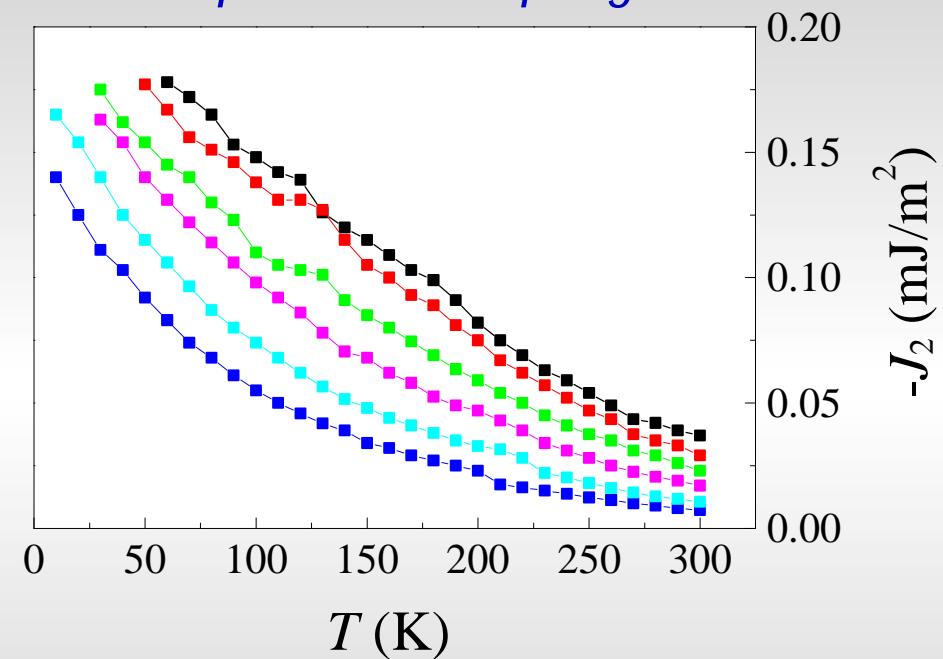
*Ge(001) / 60Å Fe /  $t$  Si / 45Å Fe / 30Å Si*

*bilinear coupling*



↑  
decreasing weakly  
with temperature

*biquadratic coupling*



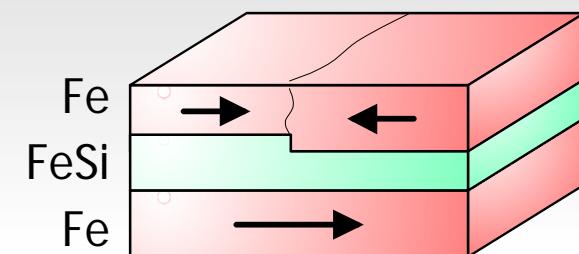
↑  
strong temperature  
dependence

## Possible origins of strong biquadratic coupling

✗ - intrinsic higher order term  $J_2(T) \propto 2 J_1(T)$

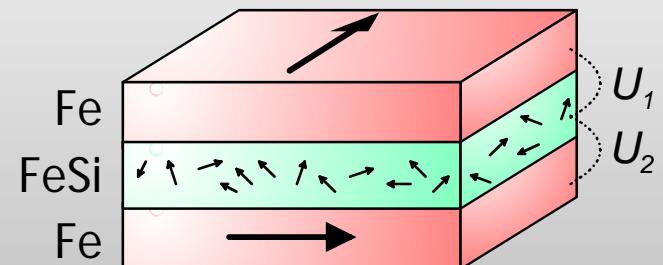
✗ - thickness fluctuations  $J_2(T) \propto (J_1(T))^2$

J.C. Slonczewski, Phys. Rev. Lett. 67, 3172 (1991)



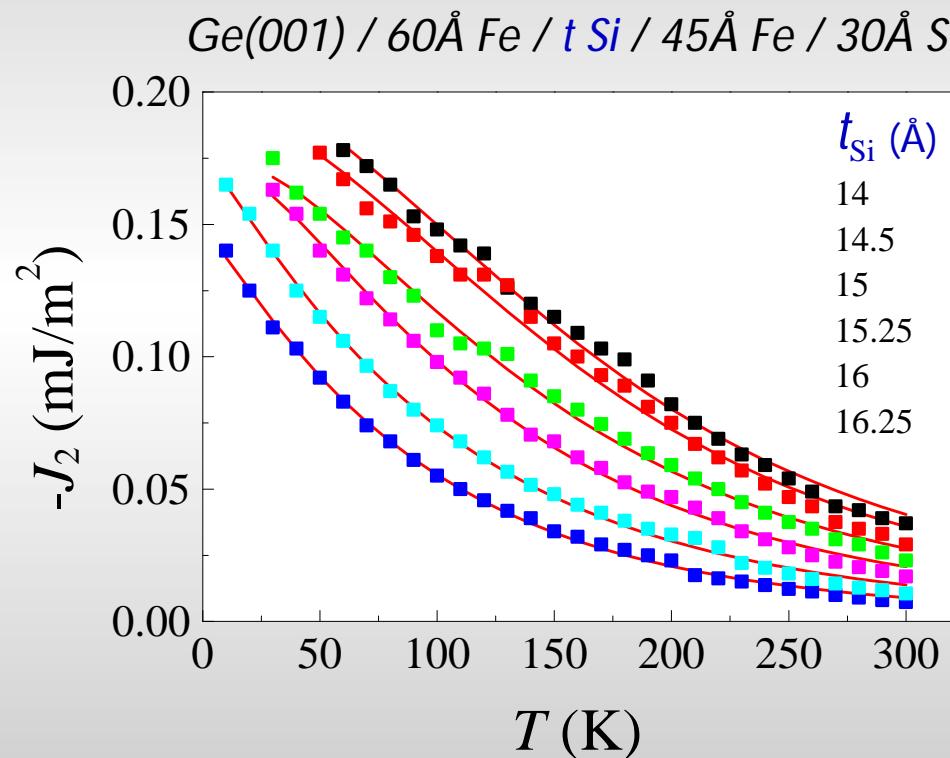
✓ - loose spins, strong temperature dependence of  $J_2$

J.C. Slonczewski, J. Appl. Phys. 73, 5957 (1993)



## Loose spins model

- strong temperature dependence of the biquadratic coupling can be described with the loose spins model

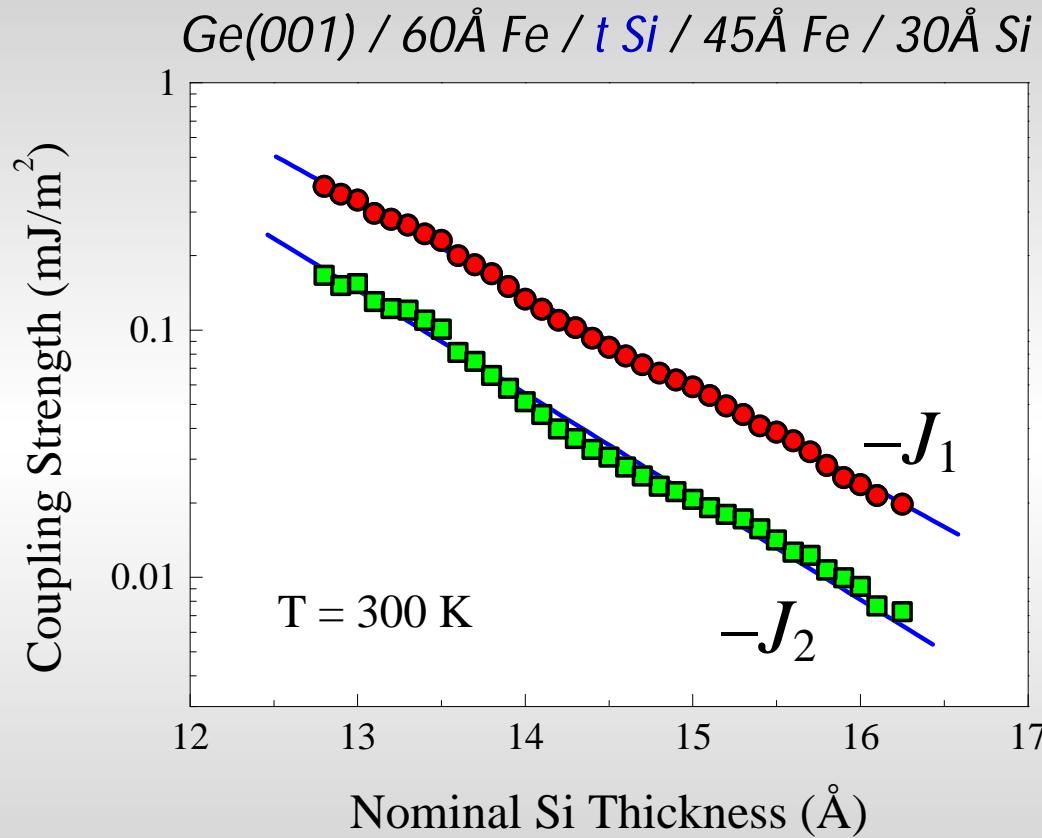


- concentration of 1% loose spins in the spacer
- interaction potentials  $U_1 = U_2$
- interaction potential  $U/k_B \approx 200 - 340$  K

G.J. Strijkers, J.T. Kohlhepp *et al.*  
J. Appl. Phys. 87, 5452 (2000)

## Qualitative proof of loose spins model

- bilinear and biquadratic coupling constants should have the same thickness dependence because of the identical interaction potential !



- indeed both  $J_1$  and  $J_2$  decay exponentially with the spacer thickness with approximately the same decay length  $\lambda$

## Quantitative proof of loose spins model

- theoretical intrinsic  $J_1$  under the assumption that loose spins are located at midplane:

$$J_1(t_{\text{Si}}, 0 \text{ K}) = a^{-2} e^{-t_{\text{Si}}/2l} U(t_{\text{Si}})$$

↑                           ↑  
lattice constant      interlayer coupling decay length

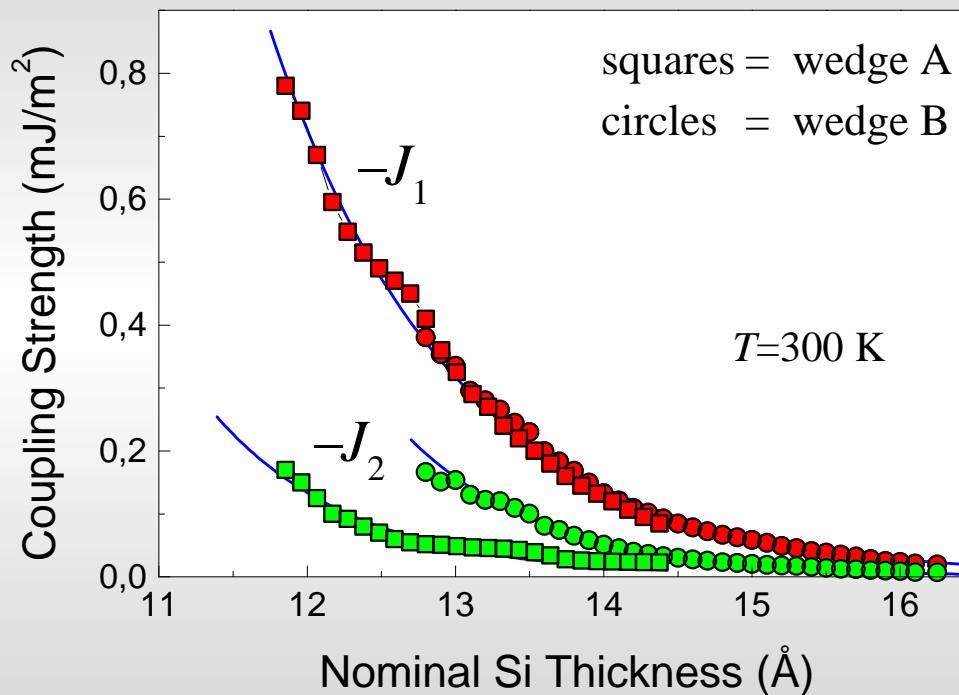
$t_{\text{Si}}$ (Å)	$J_1(\text{calc})$ (mJ/m <sup>2</sup> )	$J_1(\text{exp})$ (mJ/m <sup>2</sup> )
14.25	$0.104 \pm 0.018$	$0.126 \pm 0.010$
15.00	$0.072 \pm 0.009$	$0.075 \pm 0.008$
15.25	$0.058 \pm 0.011$	$0.046 \pm 0.004$
16.00	$0.035 \pm 0.008$	$0.033 \pm 0.005$
16.25	$0.028 \pm 0.005$	$0.025 \pm 0.003$

G.J. Strijkers, J.T. Kohlhepp *et al.*  
Phys. Rev. Lett. 84, 1812 (2000)

- good agreement between experimental and calculated values for  $J_1$ ;  
**but** this is only correct if there is no bilinear loose spin contribution !

## *Is there a bilinear loose spin contribution ?*

- Apparently not !  
In slightly different prepared samples (different concentrations of loose spins in the FeSi spacer)  $J_1(t_{\text{Si}})$  is unchanged but  $J_2(t_{\text{Si}})$  varies:



- $J_1$  is apparently the only contribution to the overall bilinear coupling; virtually no bilinear loose spin contribution is observed !

# Conclusions

- in MBE-grown Fe/Si/Fe trilayers a *metastable CsCl structure* is maintained throughout the stack; a crystalline iron-silicide with a *metastable CsCl structure* is formed in the spacer layer
- the magnetization behavior can be fully understood and described with bilinear and biquadratic coupling constants
- the *biquadratic* coupling in Fe/Si/Fe is caused by loose spins in the FeSi spacer layer

*$J_1$  and  $J_2$  are caused by the same interaction potential !*

- virtually no *bilinear loose spin* contribution is observed