

Results from RHIC

Selected results from hydrodynamics and applications of AdS/CFT at RHIC

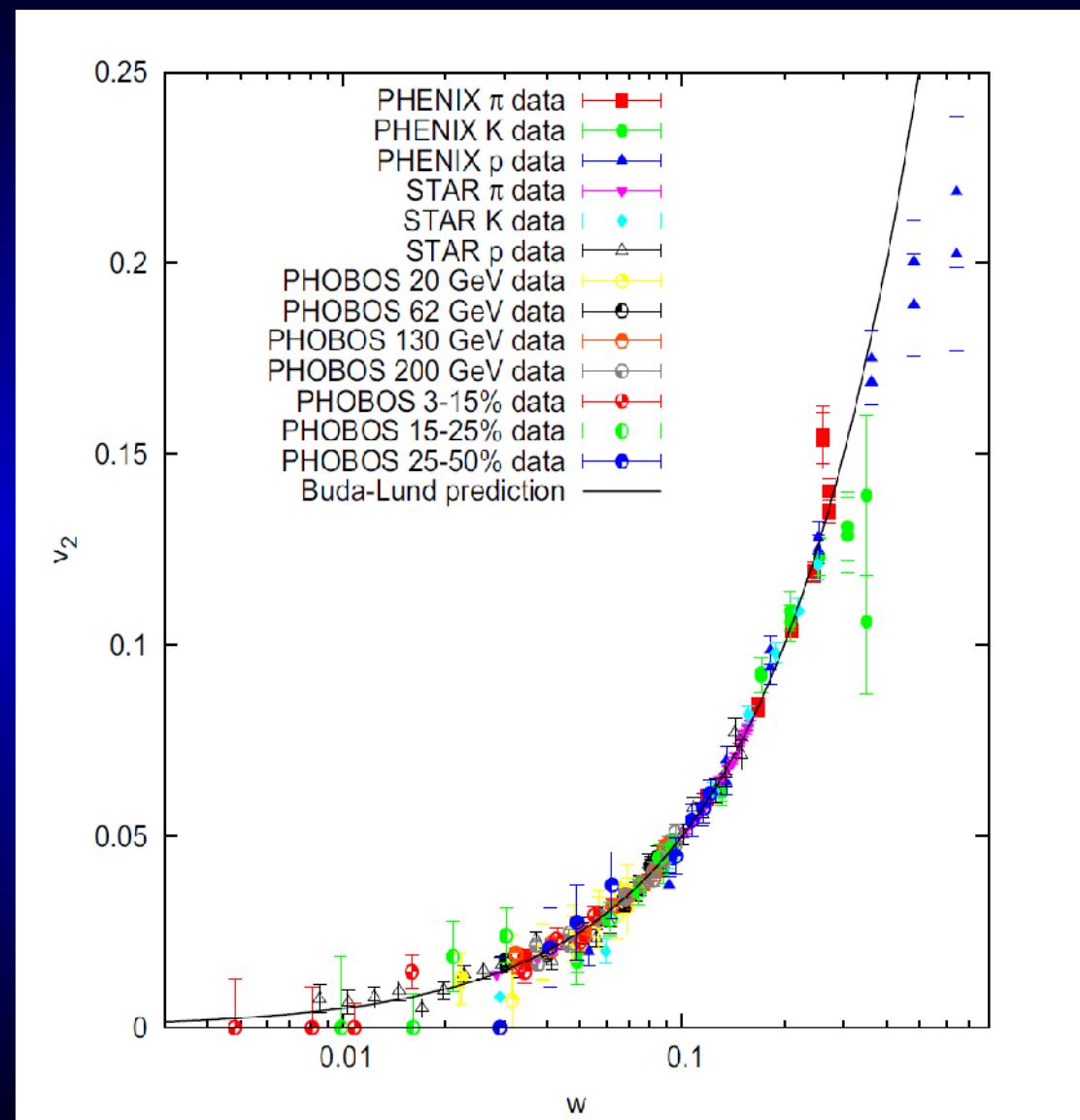
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based on S. Gubser, arXiv:0907.4808 [hep-th]

and works of R. Baier, M. Csanad, M.I. Nagy, S. Gubser, R. Vertesi and A. Ster

How well fluid dynamics works at RHIC?



Scaling predictions for (viscous) fluid dynamics

$$T'_x = T_f + m \dot{X}_f^2 ,$$

$$T'_y = T_f + m \dot{Y}_f^2 ,$$

$$T'_z = T_f + m \dot{Z}_f^2 .$$

- Slope parameters increase linearly with mass
- Elliptic flow is a universal function its variable w is proportional to transverse kinetic energy and depends on slope differences.

$$v_2 = \frac{I_1(w)}{I_0(w)}$$

$$w = \frac{k_t^2}{4m} \left(\frac{1}{T'_y} - \frac{1}{T_x} \right) ,$$

$$w = \frac{E_K}{2T_*} \varepsilon$$

Inverse of the HBT radii increase linearly with mass analysis shows that they are asymptotically the same

Relativistic correction: $m \rightarrow m_t$

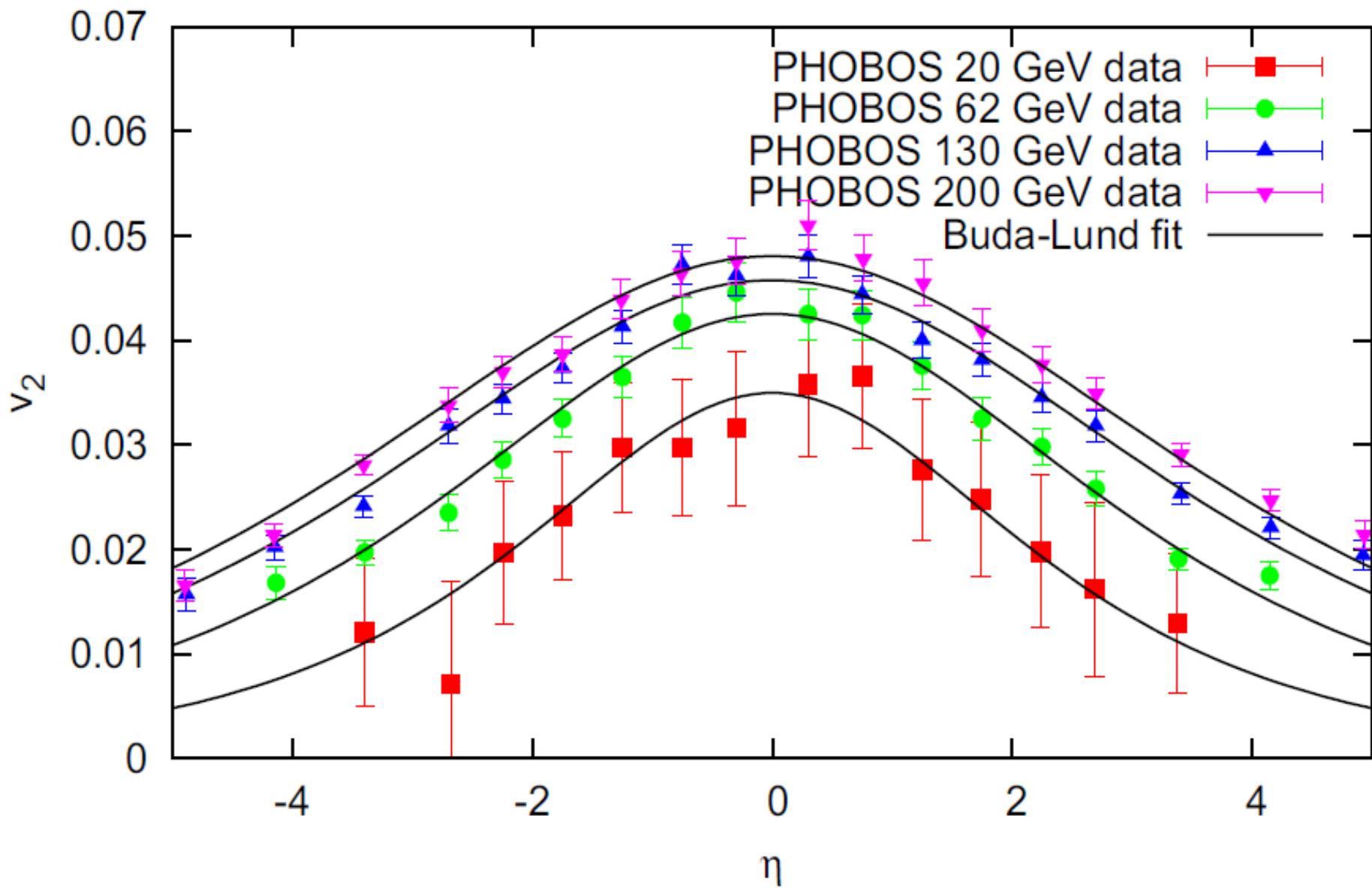
hep-ph/0108067,
nucl-th/0206051, visc. in prep.

$$R'_x{}^{-2} = X_f^{-2} \left(1 + \frac{m}{T_f} \dot{X}_f^2 \right) ,$$

$$R'_y{}^{-2} = Y_f^{-2} \left(1 + \frac{m}{T_f} \dot{Y}_f^2 \right) ,$$

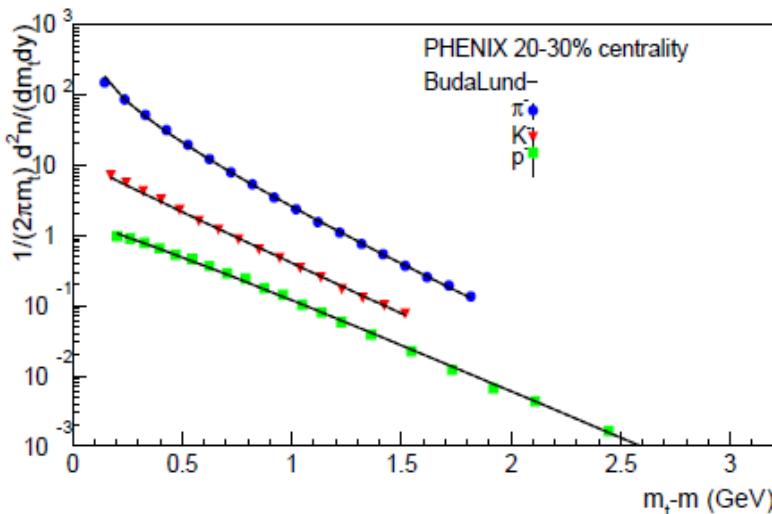
$$R'_z{}^{-2} = Z_f^{-2} \left(1 + \frac{m}{T_f} \dot{Z}_f^2 \right) .$$

Results

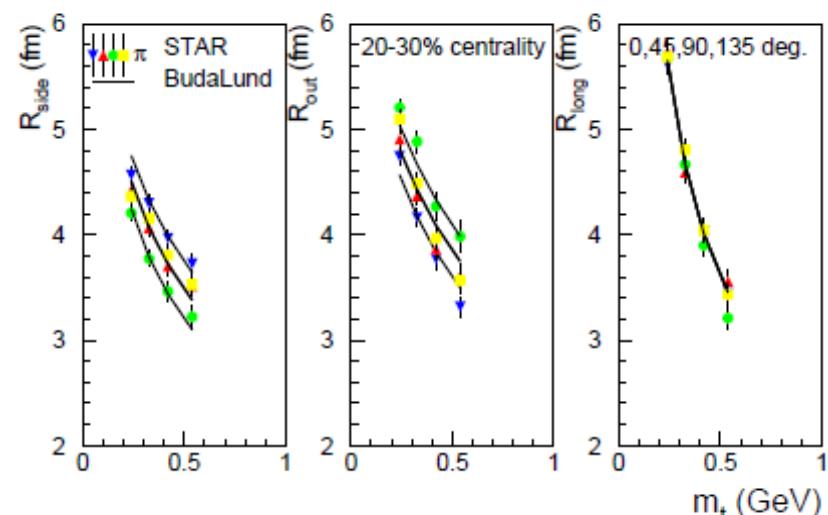
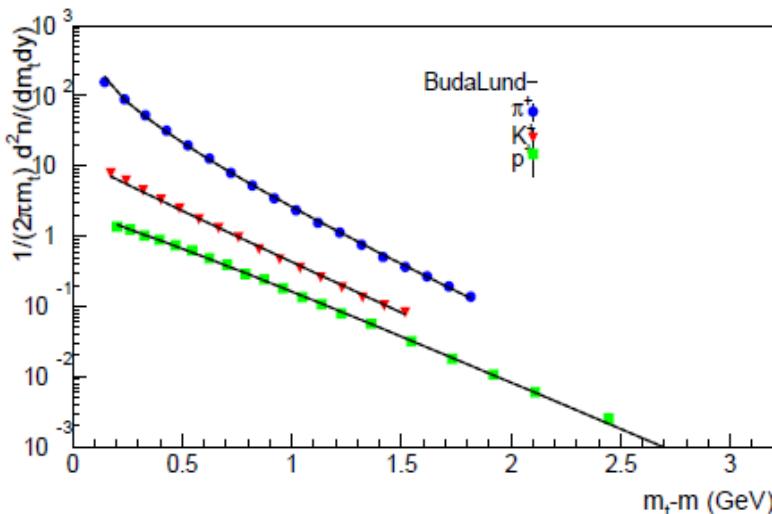
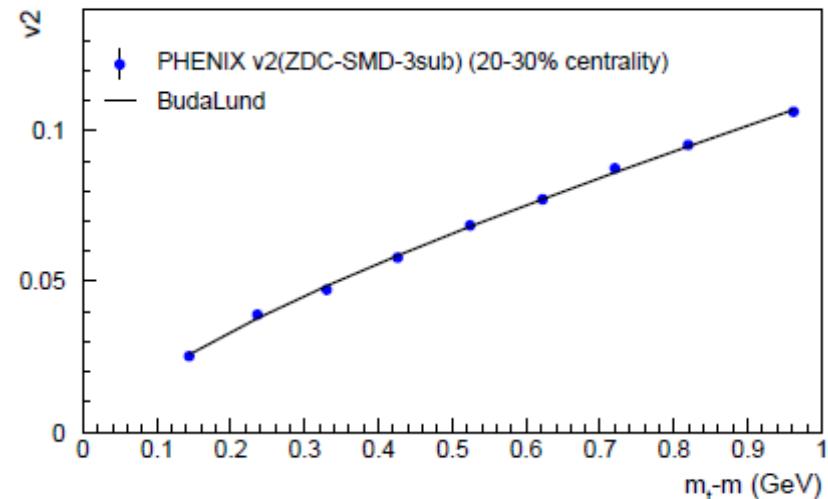


Simultaneous description of soft sector

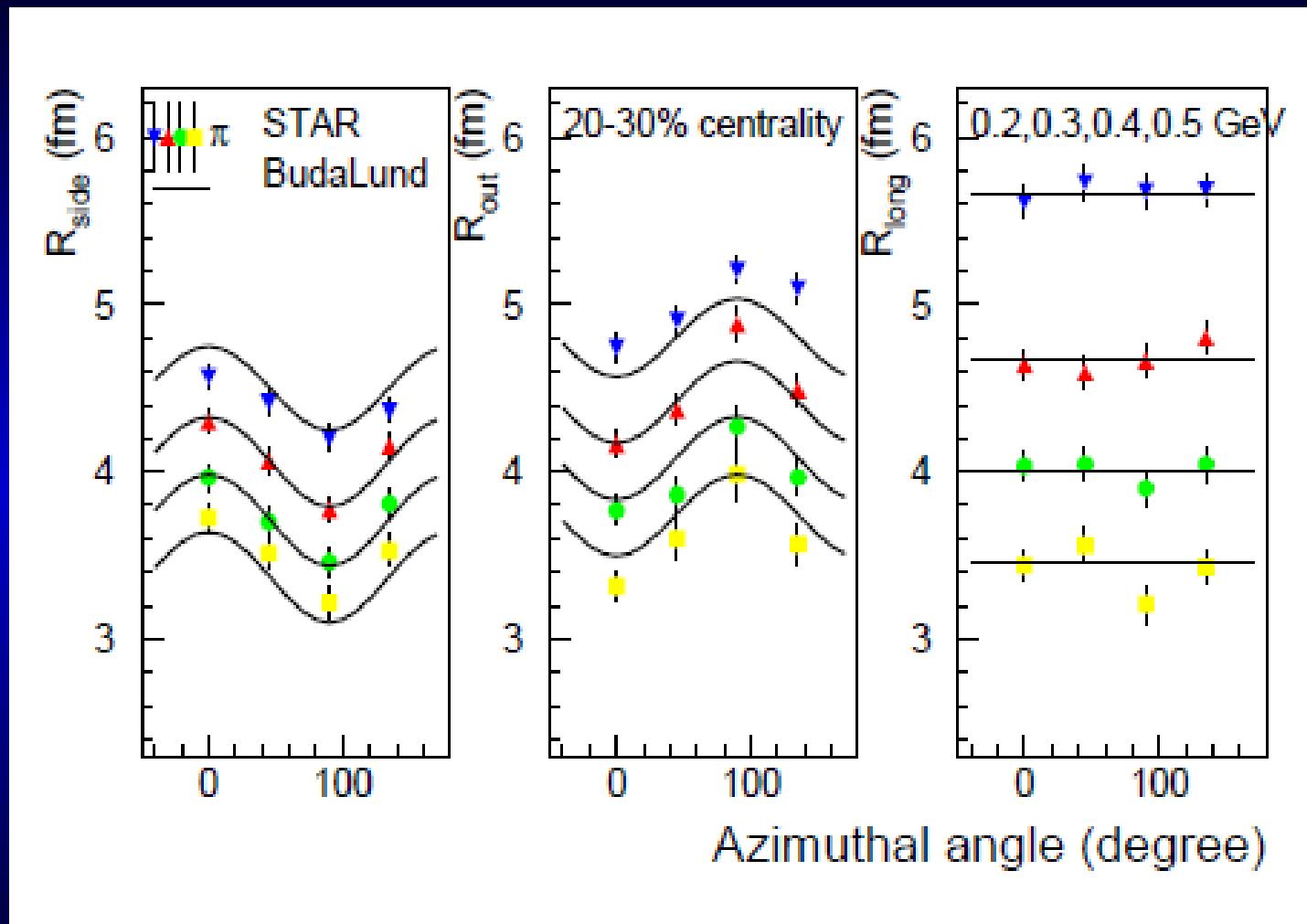
BudaLund hydro fits to 200 GeV Au+Au



BudaLund hydro fits to 200 GeV Au+Au



Simultaneous description of soft sector



Hydro works perfectly, utilizing scaling solutions, but ...

New, simple, exact solutions of rel. hydro

$$v = \tanh \lambda \eta,$$
$$p = p_0 \left(\frac{\tau_0}{\tau} \right)^{\lambda d \frac{\kappa+1}{\kappa}} \left(\cosh \frac{\eta}{2} \right)^{-(d-1)\phi_\lambda}$$

Possible cases (one row of the table is one solution):

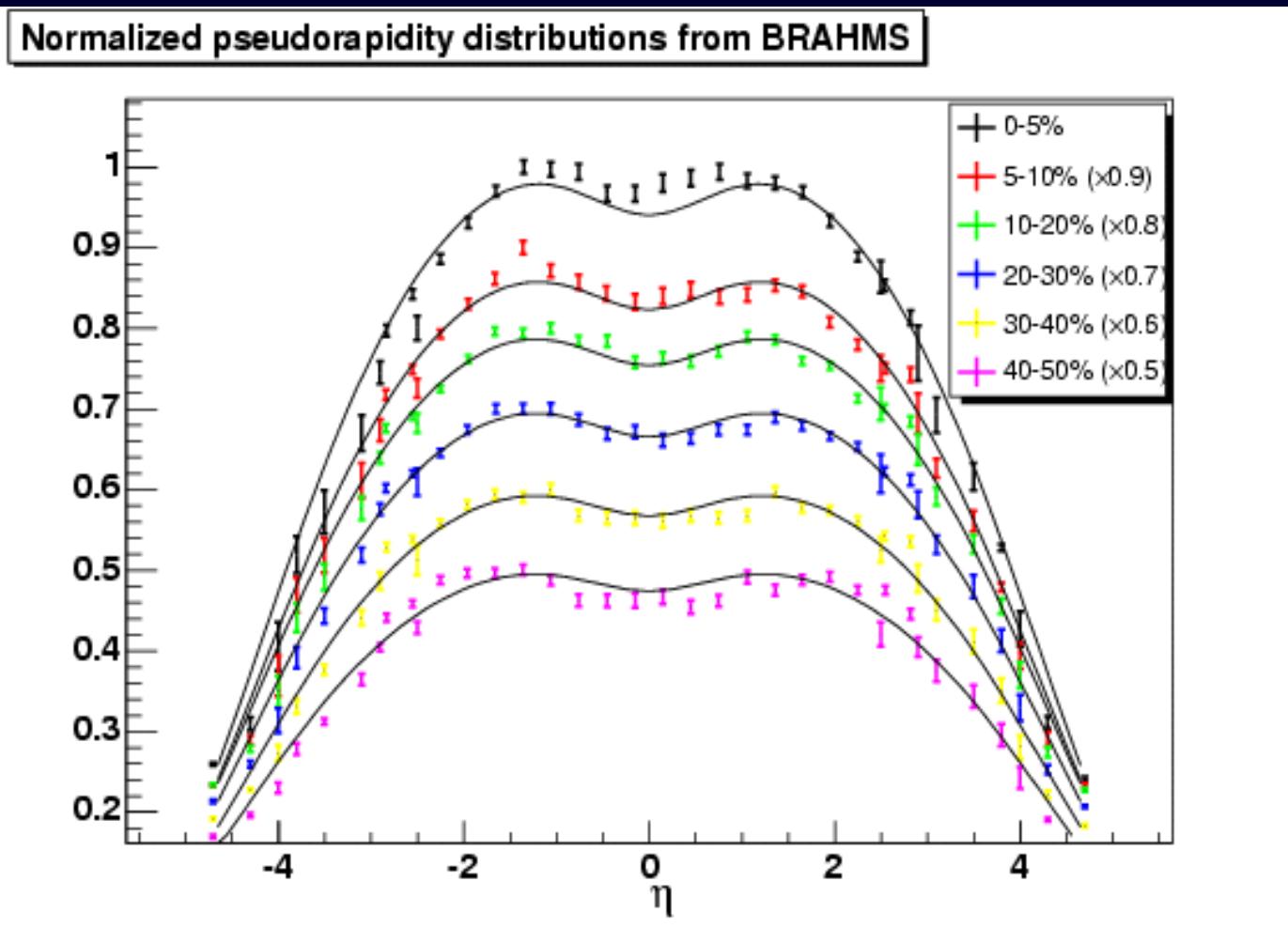
Case	λ	d	κ	ϕ_λ
a.)	2	$\in \mathbb{R}$	d	0
b.)	$\frac{1}{2}$	$\in \mathbb{R}$	1	$\frac{\kappa+1}{\kappa}$
c.)	$\frac{3}{2}$	$\in \mathbb{R}$	$\frac{4d-1}{3}$	$\frac{\kappa+1}{\kappa}$
d.)	1	$\in \mathbb{R}$	$\in \mathbb{R}$	0
e.)	$\in \mathbb{R}$	1	1	0

Nagy, Cs.T., Csanad: [nucl-th/0605070](#),
[arXiv:0709.3677v1](#), [arXiv:0805.1562](#)

- New, accelerating, d dimension
- d dimensional with $p=p(\tau, \eta)$
- (thanks T. S. Biro)
- Hwa-Bjorken, Buda-Lund type
- Special EoS, but general velocity

If $\kappa = d = 1$, general solution is obtained, for
ARBITRARY initial conditions. It is STABLE !

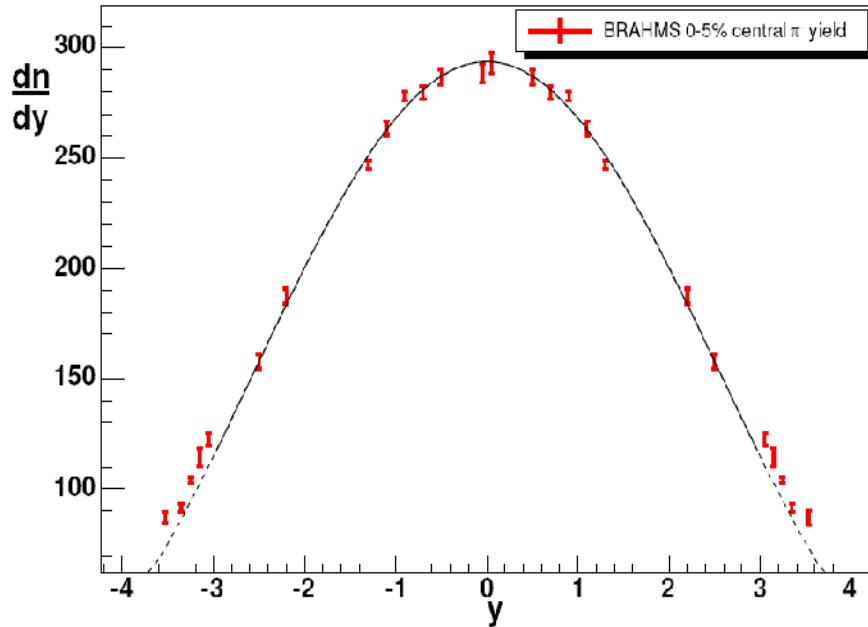
Pseudorapidity distributions



BRAHMS data fitted with the analytic formula of
Additionally: $y \rightarrow \eta$ transformation

BRAHMS rapidity distribution

$$\frac{dn}{dy} \approx \left. \frac{dn}{dy} \right|_{y=0} \cosh^{\pm \frac{\alpha}{2} - 1} \left(\frac{y}{\alpha} \right) e^{-\frac{m}{T_f} [\cosh^\alpha \left(\frac{y}{\alpha} \right) - 1]},$$



$$\lambda = \frac{\alpha - 1}{\alpha - 2}.$$

α	7.4 ± 0.13
$\left. \frac{dn}{dy} \right _{y=0}$	294 ± 1
χ^2/NDF	$30.6/14$
CL	0.6%
T_f (MeV)	200 (fixed)
m (MeV)	140 (fixed)
λ	1.18 ± 0.01 (derived)

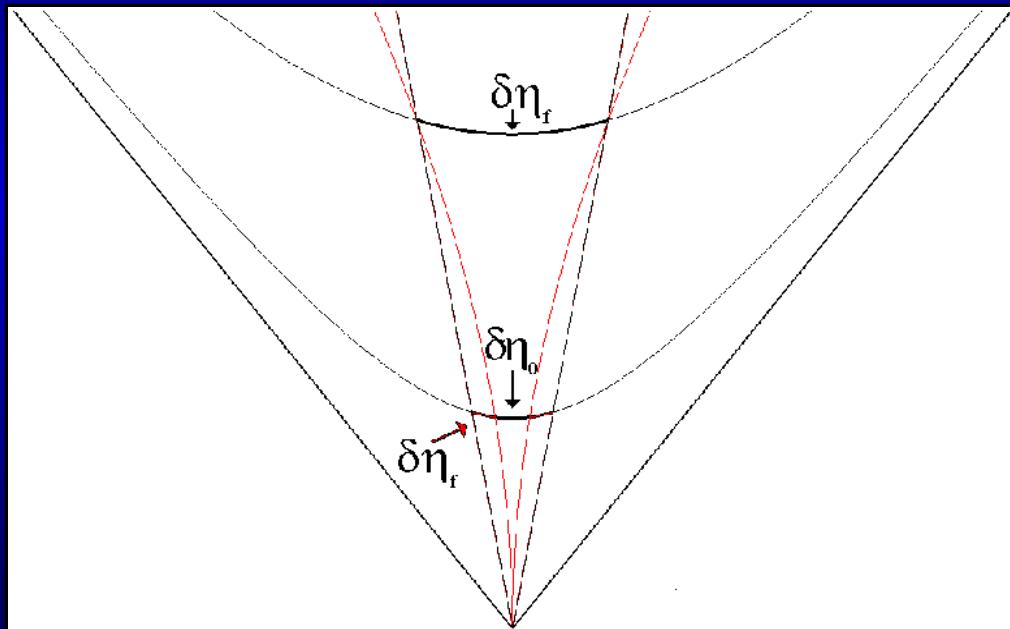
BRAHMS data fitted with the above analytic formula

Advanced energy density estimate

Fit result: $\lambda > 1$

Flows accelerate: \Rightarrow do work

\Rightarrow initial energy density is higher than Bjorken's



$$\frac{\varepsilon_c}{\varepsilon_{Bj}} = (2\lambda - 1) \left(\frac{\tau_f}{\tau_0} \right)^{\lambda-1}$$

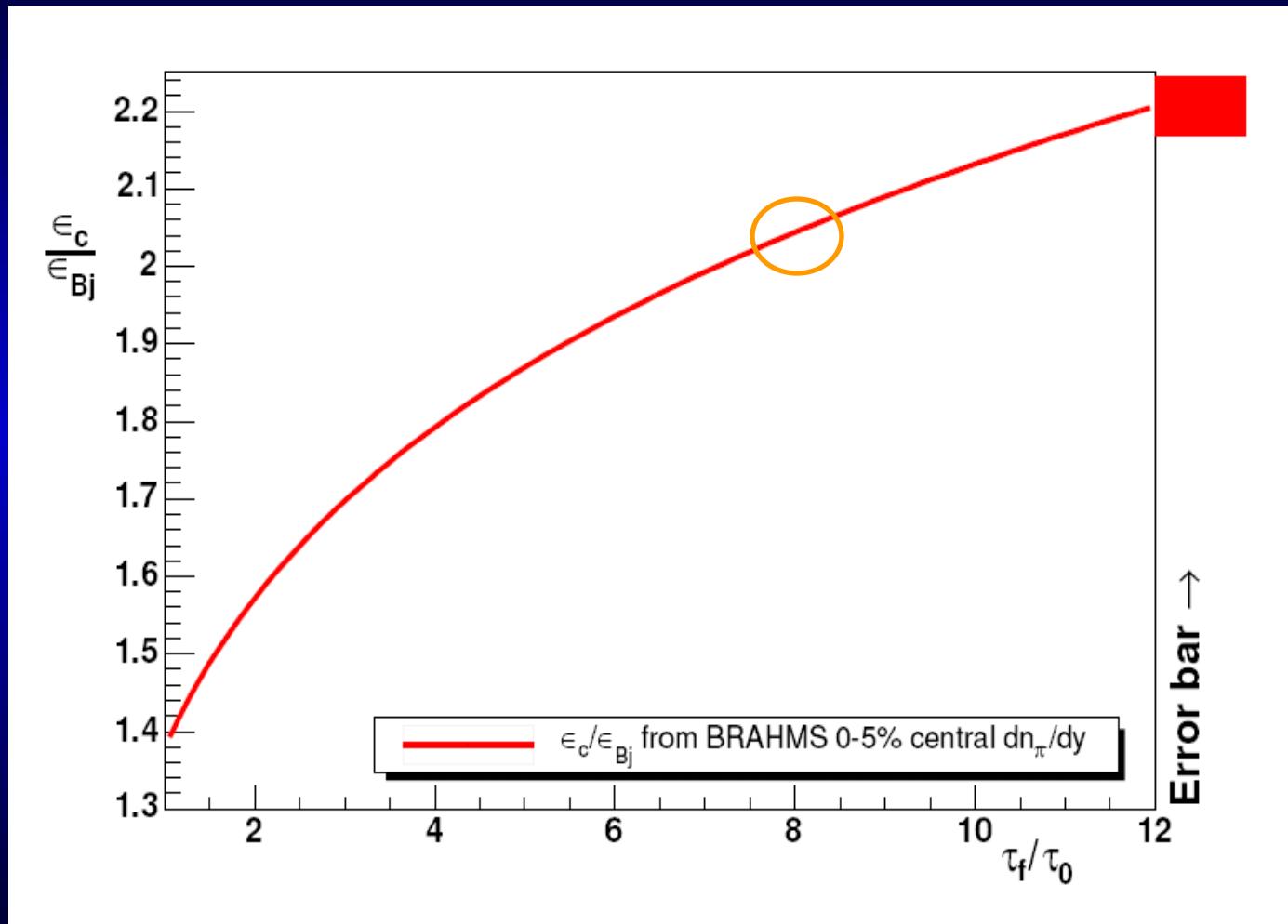
Corrections due to work and acceleration. Ref:

$$\varepsilon_{Bj} = \frac{\langle m_t \rangle}{(R^2 \pi) \tau_0} \frac{dn}{dy}$$

For $\lambda > 1$ (accelerating) flows, both factors > 1

Advanced energy density estimate

Correction depends on timescales, dependence is:



With a typical τ_f/τ_0 of $\sim 8-10$, one gets a correction factor of 2!

Conjecture: EoS dependence of ε_0

Four constraints

1) ε_{Bj} is independent of EoS ($\lambda = 1$ case)

2) $c_s^2 = 1$ case is solved for any $\lambda > 0.5$

$$\frac{\varepsilon_c}{\varepsilon_{Bj}} = (2\lambda - 1) \left(\frac{\tau_f}{\tau_0} \right)^{\lambda-1}$$

Corrections due to respect these limits.

3) c_s^2 dependence of $\varepsilon(\tau)$ is known

4) Numerical hydro results

Conjectured formula – given by the principle of Occam's razor:

$$\frac{\varepsilon_{c_s^2}}{\varepsilon_{Bj}} = (2\lambda - 1) \left(\frac{\tau_f}{\tau_0} \right)^{\lambda-1} \left(\frac{\tau_f}{\tau_0} \right)^{(\lambda-1)(1-c_s^2)}$$

Using $\lambda = 1.18$, $c_s = 0.35$, $\tau_f/\tau_0 = 10$, we get $e_{cs}/e_{Bj} = 2.9$

$\varepsilon_0 = 14.5 \text{ GeV/fm}^3$ in 200 GeV, 0-5 % Au+Au at RHIC

Perfect hydro conclusion

New exact, accelerating relativistic hydrodynamic sols.

Analytic (approximate) calculation of observables

Realistic rapidity distributions; BRAHMS data well described

No go theorem: same final states, different initial states

New estimate of initial energy density:

$\epsilon_c/\epsilon_{Bj} \sim 2$ @ RHIC

dependence on c_s estimated, $\epsilon_c/\epsilon_{Bj} \sim 3$ for $c_s = 0.35$

A lot to gain ...

advanced energy density estimate changes completely

the $s_{NN}^{1/2}$ or the centrality, where $\epsilon_{\text{crit},\text{IQCD}} \sim 1 \text{ GeV/fm}^3$ is reached

But again, how perfect is Quark Fluid?

Status at QM 2009 (April 09)

Adopted from A. Tang:

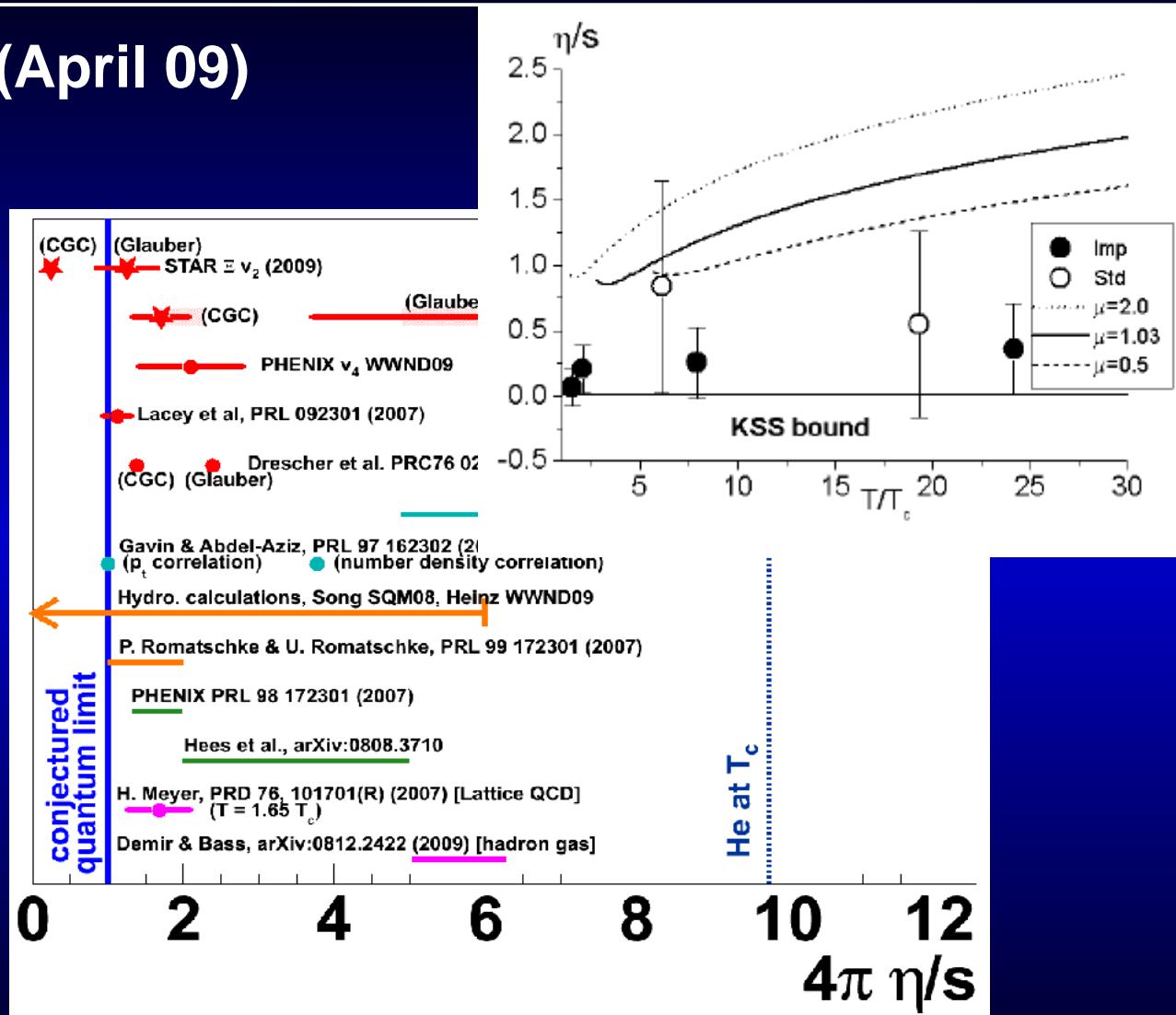
It is very low!

small η/s

large σ

large α_s

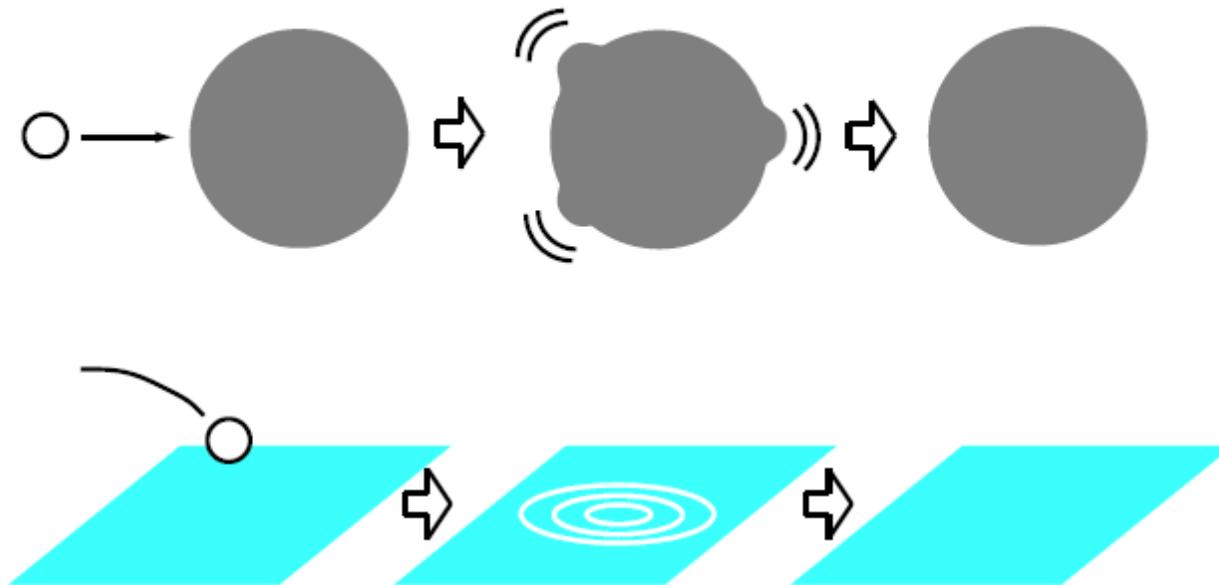
strong coupling limit
of QCD



CFT/AdS/gravity correspondence:
black holes at RHIC?

Perfect Fluid- Gravity duality

Illustration:



[from M. Natsuume]

quasinormal modes:

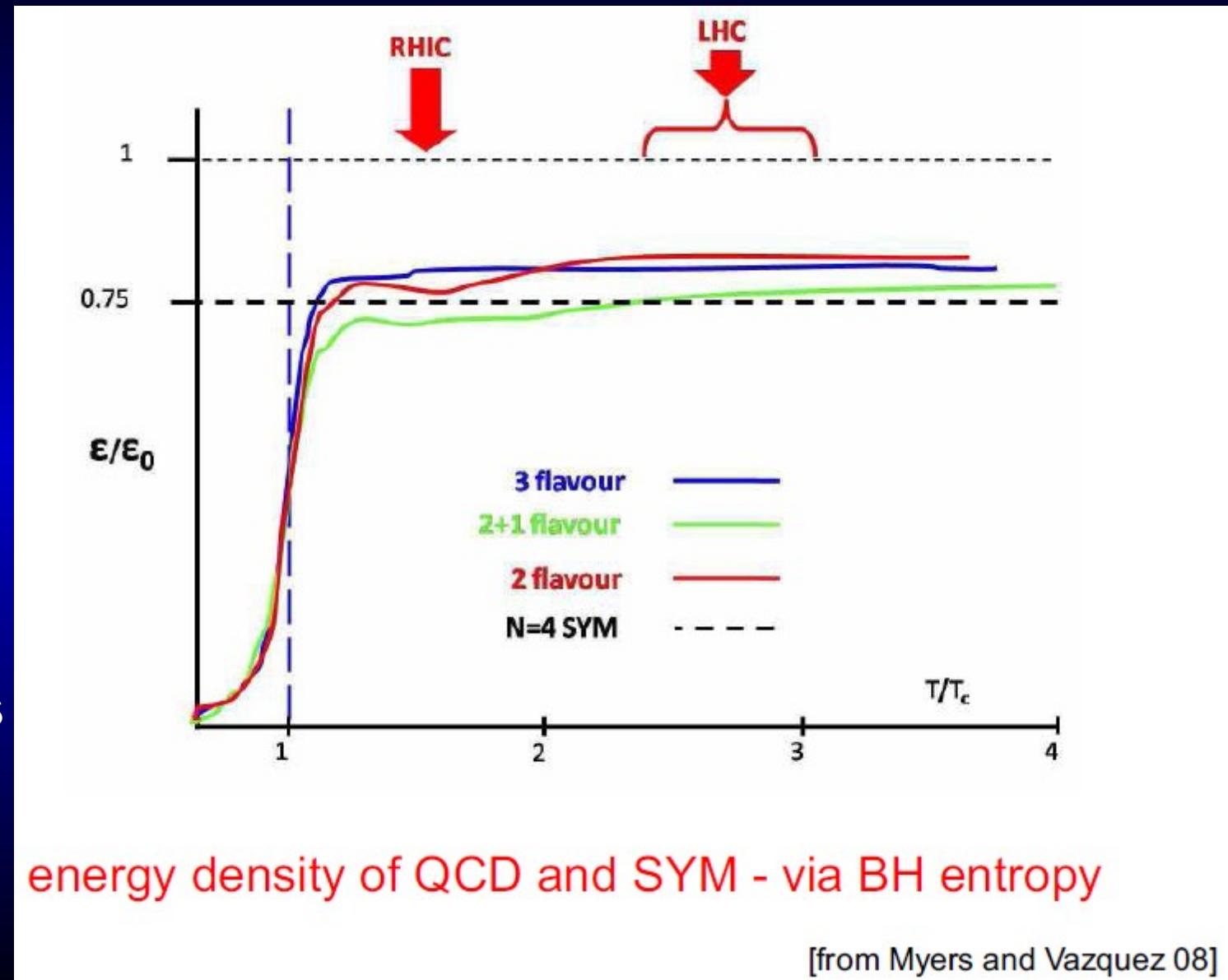
gravitational perturbation to a black hole
and to a hydrodynamic system

What about equation of state?

Equation
of State
from
AdS/CFT:

25% less
e than
that of
ideal gas
of massless
quanta

But perils!



What about shear viscosity?

near equilibrium: $\eta \simeq \epsilon \bar{v} \lambda_f$, *entropy density* $s \simeq \epsilon/m \rightarrow$

$$\frac{\eta}{s} \simeq m \bar{v} \lambda_f \simeq \hbar \frac{\text{mean free path}}{\text{deBroglie wavelength}}$$

- dilute system (QFT \rightarrow kinetic theory \rightarrow hydro):
scale $\lambda_f \rightarrow \frac{\eta}{s} \gg \hbar$, e.g. pQCD ($N_f = 0$)

$$\frac{\eta}{s} \simeq 3.8 \frac{1}{g^4 \ln(2.8/g)} \simeq O(1) \text{ for } g = 2.5$$

BUT with $\ln(2.8/g) \simeq O(1)$: $\frac{\eta}{s} \simeq 0.1 \rightarrow$ sensitive to constant under the log !

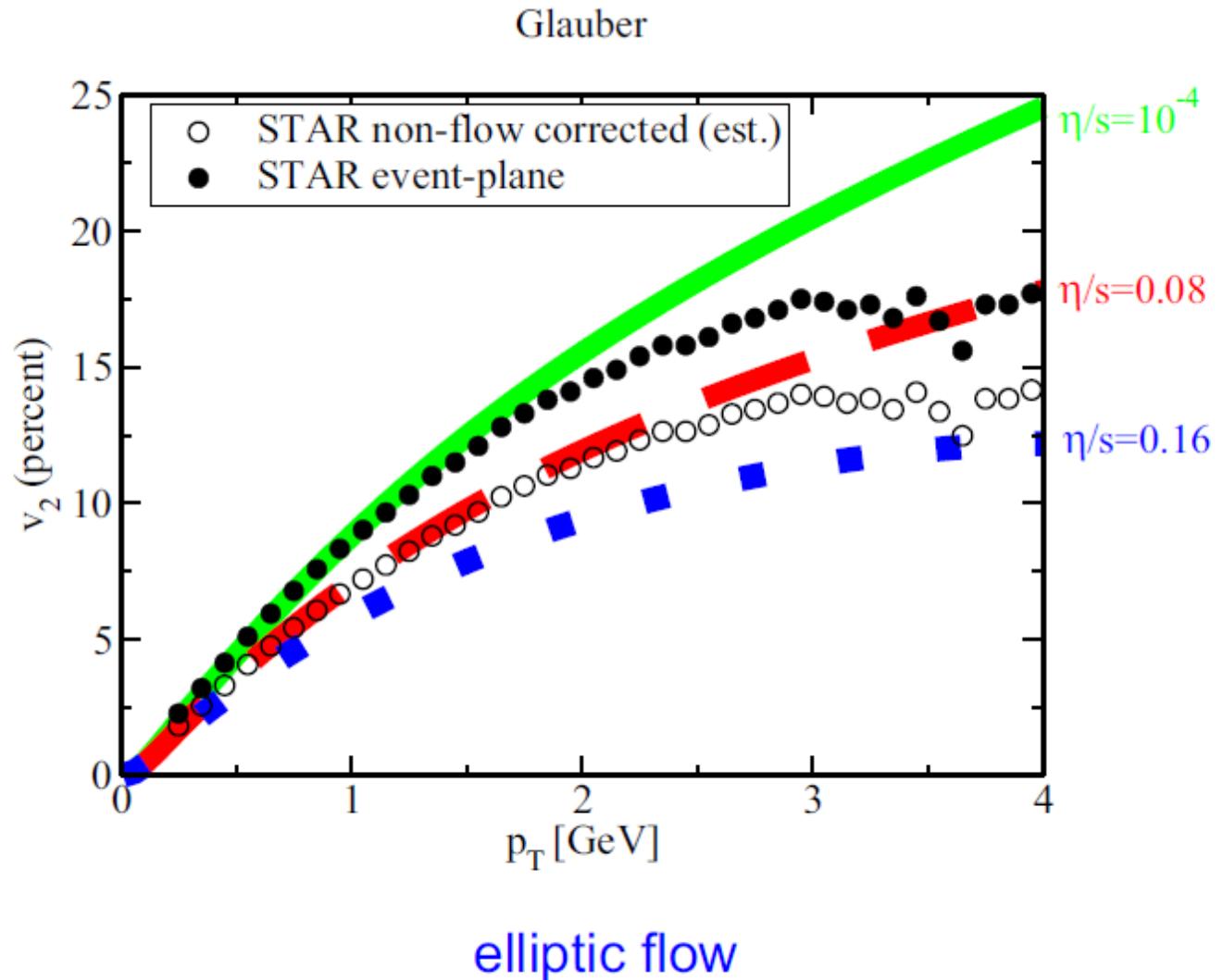
[Arnold, Moore and Yaffe 03]

- strongly coupled system (QFT \rightarrow hydro):
scale $1/T \rightarrow \frac{\eta}{s} = \frac{\hbar}{4\pi} \simeq 0.08$

[Policastro, Son and Starinets 01]

But perils!

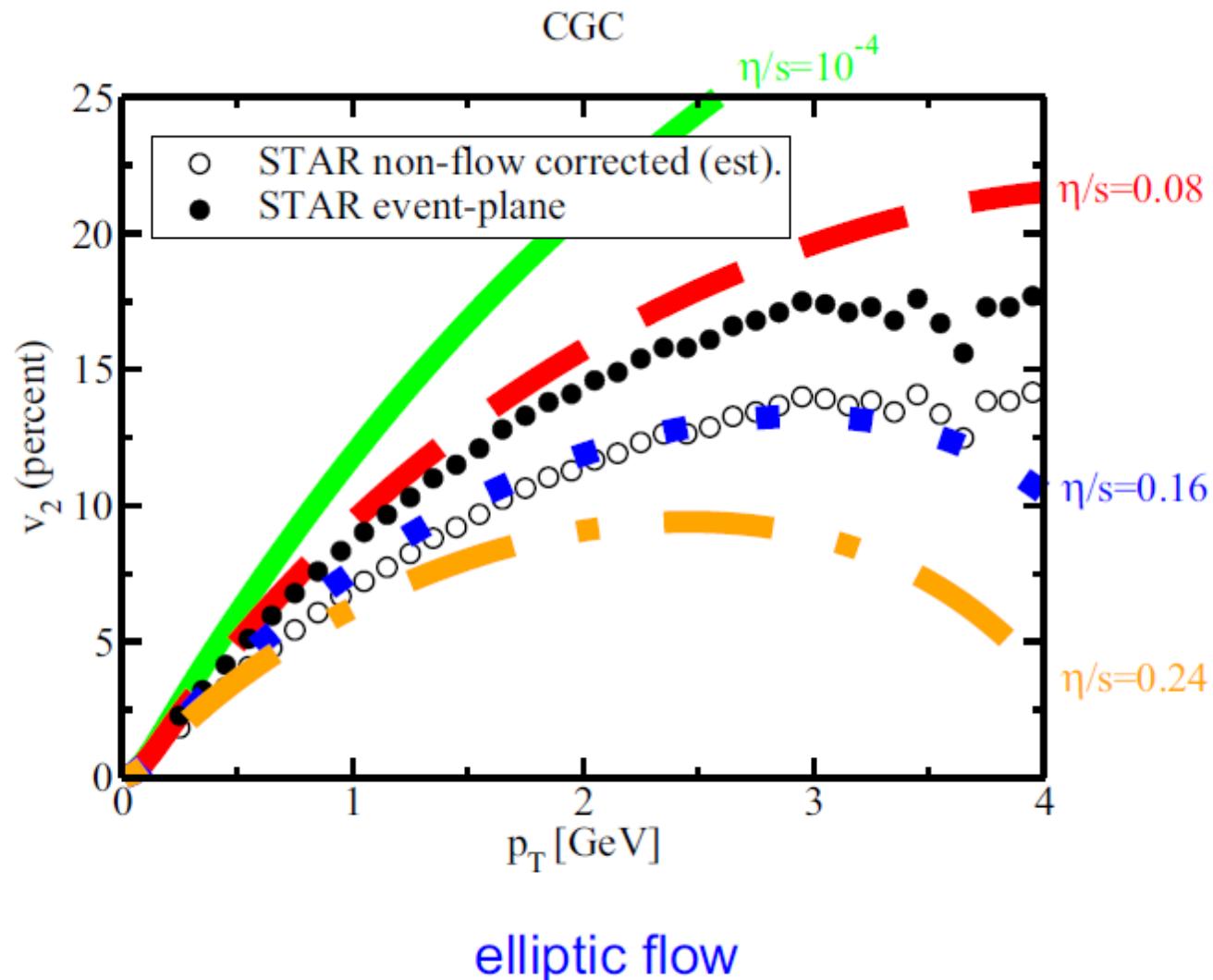
Conform hydro with shear viscosity 1



But perils!

[Luzum and Romatschke 08]

Conform hydro with shear viscosity 2



But perils!

[Luzum and Romatschke 08]

Comparision and some perils

	QCD	$\mathcal{N}=4$ SYM
$T=0$	$N_c = 3 = N_f$, confinement, discrete spectrum, scattering,	N_c large, N_f/N_c small, deconfined, conformal, supersymmetric,
	very different !!	
$T > T_c$	strongly-coupled plasma of gluons & fundamental matter deconfined, screening, finite corr. lengths, . . .	strongly-coupled plasma of gluons & adjoint and fundamental matter deconfined, screening, finite corr. lengths, . . .
	very similar !!	
$T \gg T_c$	runs to weak coupling	remains strongly-coupled
	very different !!	

QCD and $\mathcal{N} = 4$ *SYM* as a function of temperature

[from Myers and Vazquez 08]

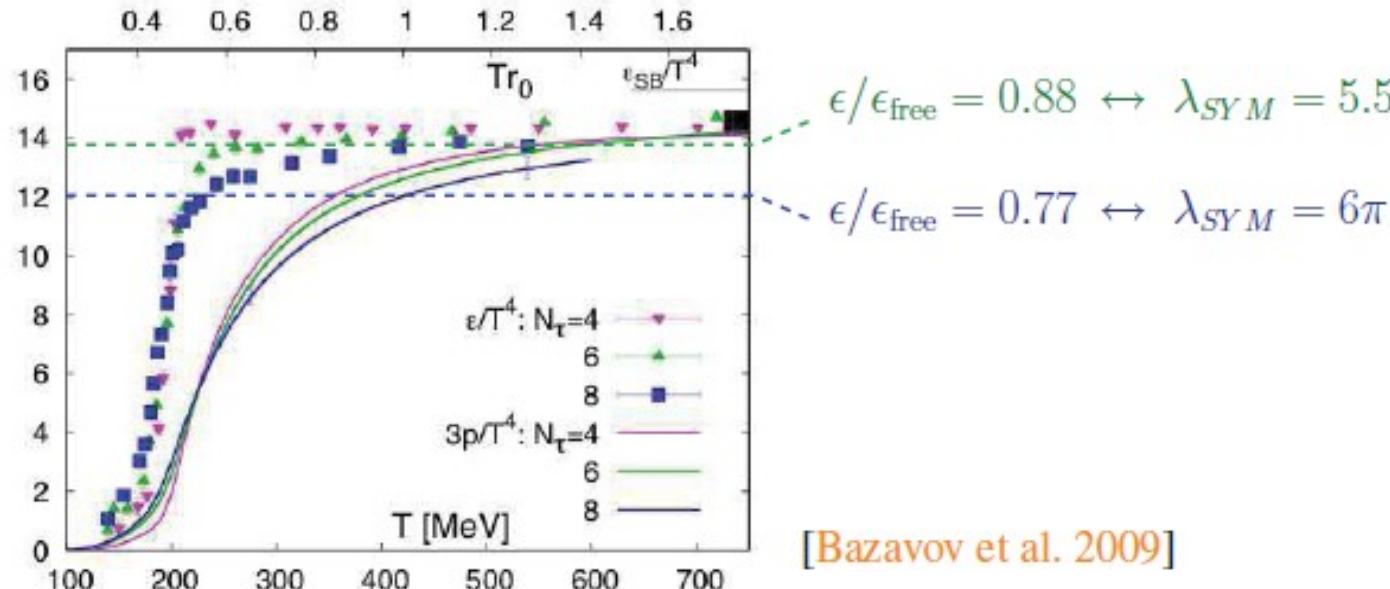
Some perils, following Gubser's talk



$\mathcal{N} = 4$ SYM has considerably more degrees of freedom as QCD: for $SU(3)$,

$$\epsilon_{\text{free}}^{\text{SYM}} \approx 39T^4$$

$$\epsilon_{\text{free}}^{\text{QCD}} \approx 16T^4 \quad (\text{3 massless flavors})$$



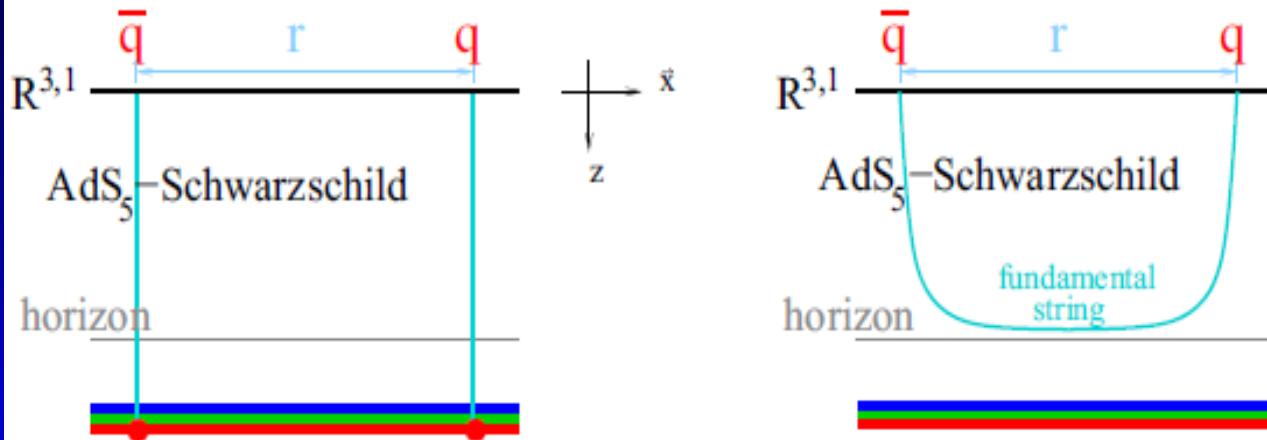
- Constant λ_{SYM} cannot match running λ_{QCD} except in a limited window. We're interested in $200 \text{ MeV} \lesssim T \lesssim 300 \text{ MeV}$.



QCD is significantly non-conformal in this window.

Static $q\bar{q}$ potential, following Gubser's

Static $q\bar{q}$ potential favors $\lambda \approx 5.5$.



- Fundamental charge arises because hanging string ends on a D3-brane.
- A color singlet pair charges can lower their energy by joining their hanging strings.

$$\text{SYM: } V_{q\bar{q}}(r) = -\frac{4\pi^2}{\Gamma(1/4)^4} \frac{\sqrt{\lambda}}{r} \times [1 - (\text{Debye screening})]$$

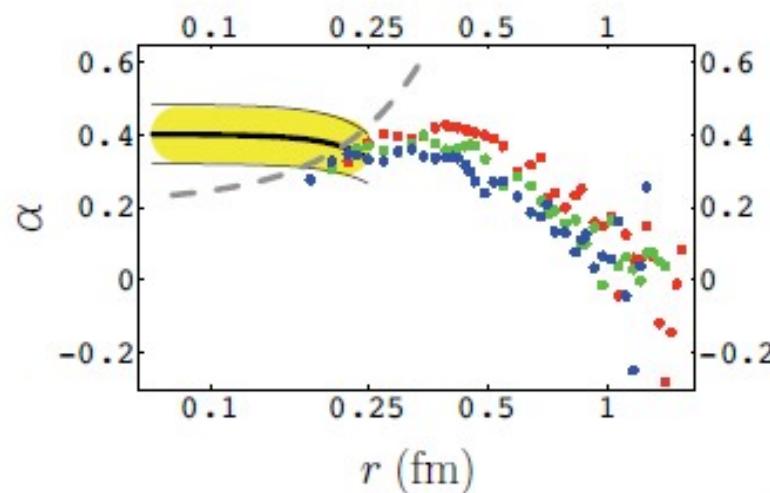
$$\text{QCD: } V_{q\bar{q}}(r) = \left(-\frac{4\alpha/3}{r} + \sigma r \right) \times [1 - (\text{Debye screening})]$$

Static qq potential redux

More precisely: $\alpha_{q\bar{q}}(r, T) = \frac{3}{4}r^2 \frac{\partial F_{q\bar{q}}}{\partial r}$ is the effective coupling.

$$T_{SYM} = 190 \text{ MeV} \leftrightarrow T_{QCD} \approx 240 \text{ MeV}$$

if we match $\epsilon_{SYM} = \epsilon_{QCD}$.



$T_{lattice} = 209, 233, 255 \text{ MeV}$
[Kaczmarek and Zantow 2005]

[Gubser 2006c]

- Curve shown corresponds to $\lambda = 5.5^{+2.5}_{-2}$ in SYM. Surprising because then $\alpha_{SYM} \approx 0.15$.



Match to string theory is conspicuously imperfect because SYM doesn't confine. And string theory curve isn't even understood for $r \gtrsim 0.25 \text{ fm}$; but see [Bak et al. 2007].

Static qq potential: fine with v2 data

$\lambda = 5.5$ significantly affects viscosity [Policastro et al. 2001; Buchel 2008]:

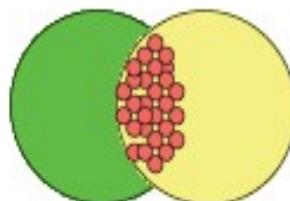
$$\frac{\eta}{s} = \frac{1}{4\pi} \left(1 + \frac{15\zeta(3)}{\lambda^{3/2}} + \dots \right) \approx 0.2 \quad \text{for } \lambda = 5.5.$$



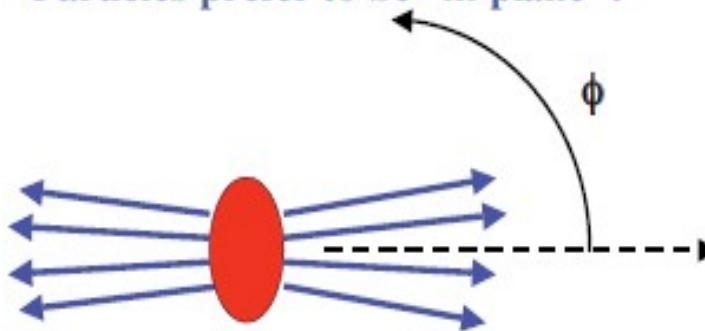
higher order terms could be important.

$\eta/s \approx 0.2$ is not far from elliptic flow data, with CGC initial eccentricity and a soft EOS.

Beam's eye view of a non-central collision:



Particles prefer to be “in plane”:



Anisotropic expansion is driven by pressure gradient, which is greater “in plane.”

Progress and perils

Bulk/Soft physics:

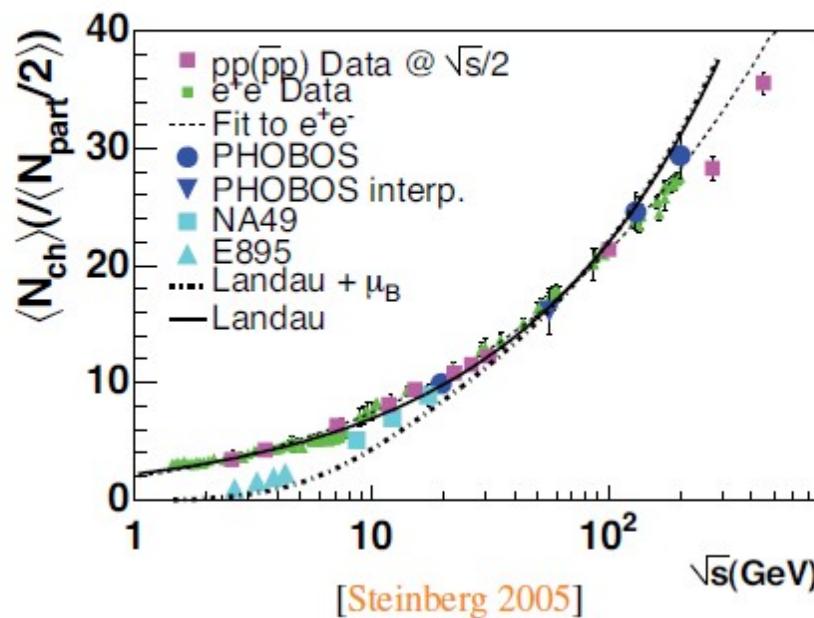
- evaluation of bulk viscosity
- dihadron correlations
- Mach cone like, non-hydro structures
- multiplicity distributions !!



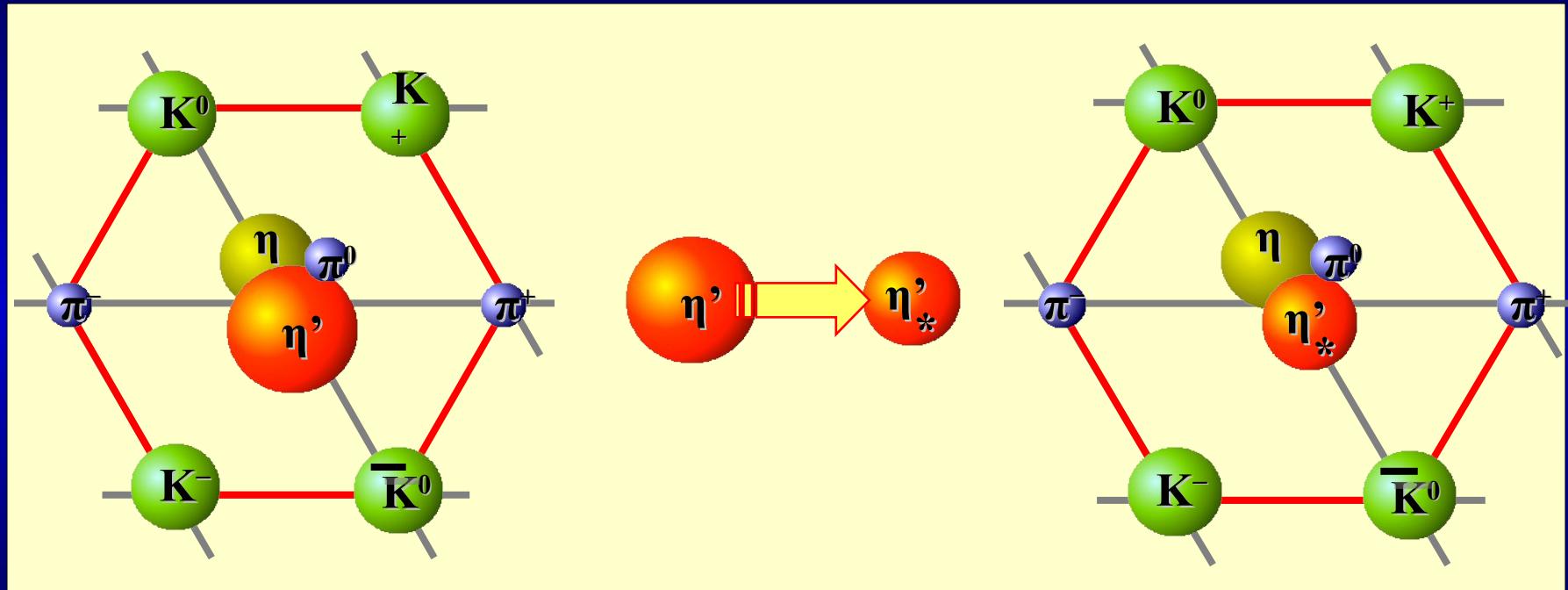
$N_{\text{charged}} \sim E^{2/3}$ is faster than Landau ($E^{1/2}$) and much faster than CGC ($\sim E^{0.3}$)

Hard physics:

See the talk
of G. Barnaföldi



Significant in-medium η' mass reduction in $\sqrt{s_{nn}} = 200$ GeV Au+Au collisions



R. Vértesi, T. Cs, J. Sziklai, arXiv:0905.2803 [nucl-th]

Motivation

- **Approximate SU(3) symmetry**

Spontaneous symmetry breaking \rightarrow 9 Goldstone bosons
 $U_A(1)$ -breaking terms $\rightarrow \eta'$ gains extra mass

Refs: T. Kunihiro, Phys.Rev.Lett. B218 363 (1989)
J.Kapusta, D.Kharzeev, L.McLerran, Phys.Rev. D53 5028 (1996)
Z.Huang, X.-N. Wang, Phys.Rev. D53 5034 (1996)

In hot medium, η' mass reduced to quark model mass

- **Signal: Enhanced η' production at low p_T**

$$\frac{N_{\eta'}^*}{N_{\eta'}} = \left(\frac{m_{\eta'}^*}{m_{\eta'}} \right)^\alpha e^{-\left(\frac{m_{\eta'}^* - m_{\eta'}}{T} \right)}$$

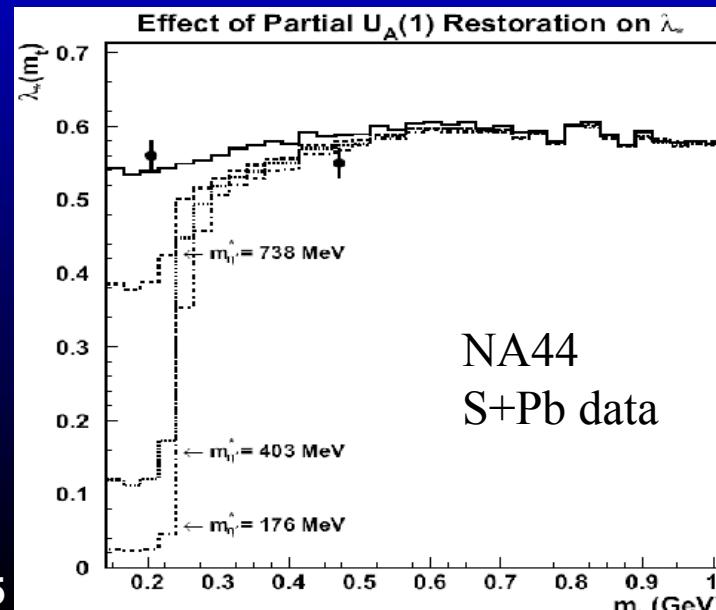
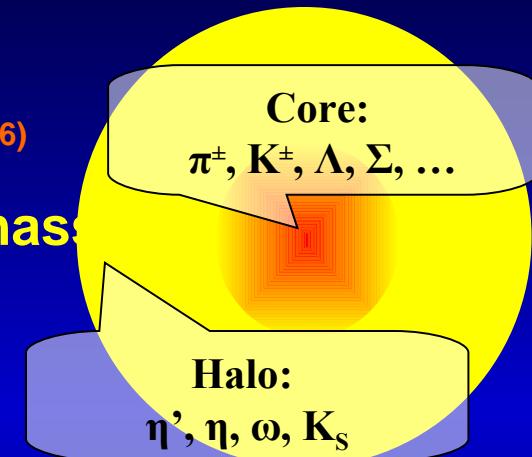
- **Observation channel: pion BEC**

η' has long lifetime, decays into pions

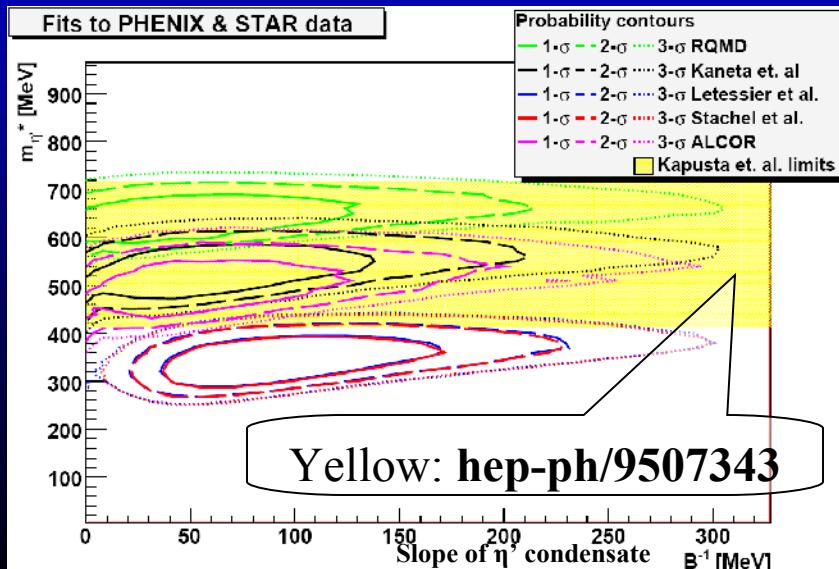
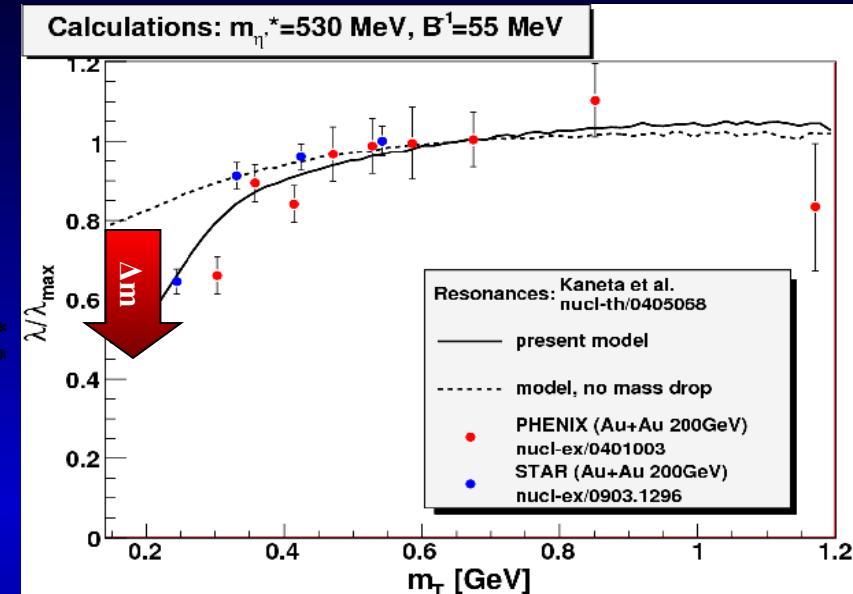
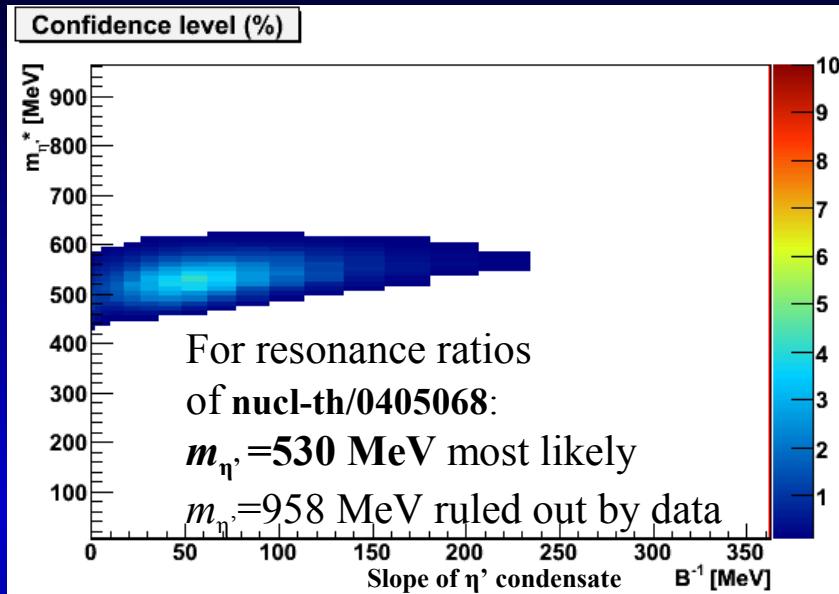
Core-halo picture: more η' 's enhance halo

Measurable through the parameter λ^*

$$\lambda^* = \left(\frac{N_{\text{core}}^{\pi^+}}{N_{\text{halo}}^{\pi^+} + N_{\text{core}}^{\pi^+}} \right)^2$$



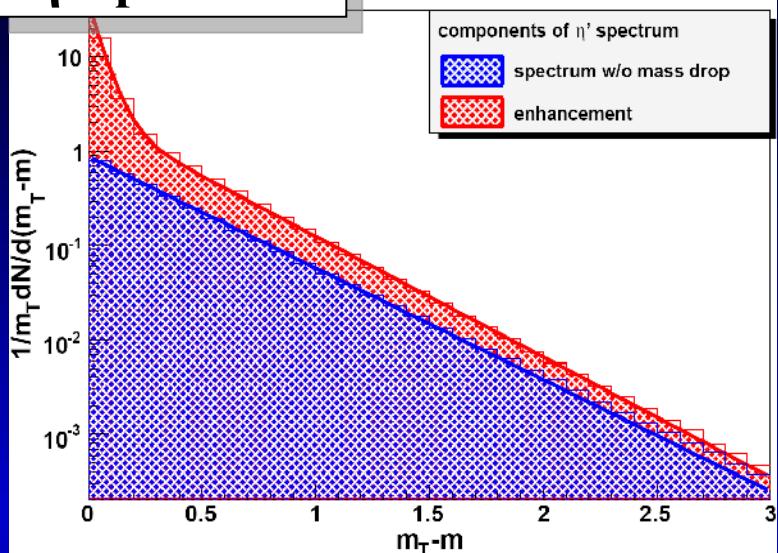
Simulation and Results



6 models for the particle ratios
(ALCOR, FRITIOF, Letessier et.al.,
Kaneta et.al., RQMD, Stachel et.al.)
Resonance decays: JETSET
Systematic studies for freezeout
parameters and prefactors

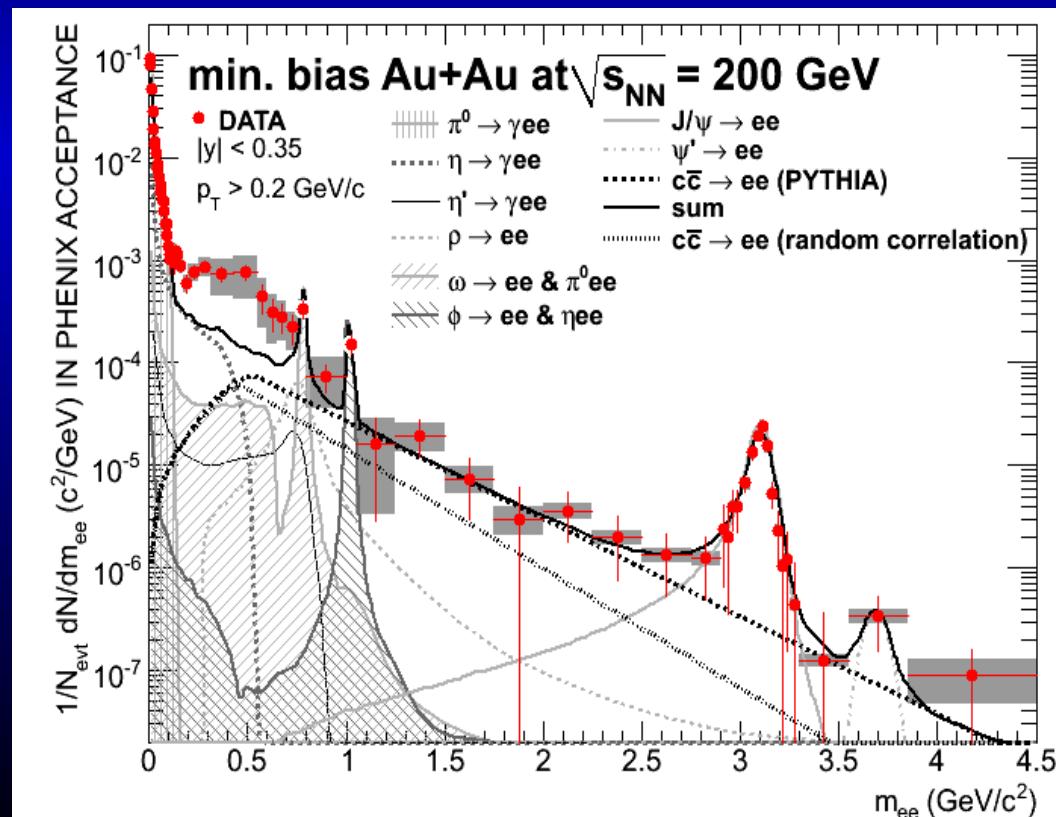
Cross-check plan: low p_T and dileptons

η' spectrum



Excess at $m_{\ell\ell} < 1$ GeV
Seen at SPS (CERES)
and RHIC (PHENIX)
Only in A+B reactions
Absent in p+p
Possible explanation:
 η' enhancement

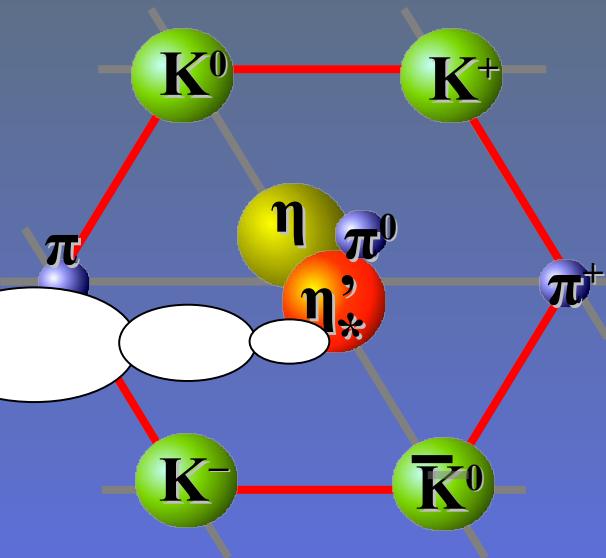
- Resonance ratios:
[nucl-th/0405068](#)
 - Enhancement factor ~ 24
 - Breaks m_T scaling for η'



Conclusion

$m_{\eta^*}^* < m_{\eta'} - 200 \text{ MeV}$
at the 99.9% confidence level

from PHENIX+STAR $\pi^+\pi^+$
correlation data + 6 models



Cross-check with dilepton spectrum needed

More λ^* data at low p_T is needed to reduce systematics

-> Chiral restoration at lower T as deconfinement?

Revitalize interest in chiral symmetry restoration

Summary II

How well hydro works at RHIC?

in $p_t < 1.5$ GeV: it works perfectly

scaling laws for spectra, v_2 , HBT radii predicted
but: quark number scaling of v_2 is exp. discovery

The hottest, most perfect and most opaque fluid - ever made

AdS/CFT correspondence: PRogress and PERILS

- (1) Exact results in hydro: see Bajnok's talk.
- (2) New scaling laws discovered
- (3) QCD: Possibly chiral symmetry restored at lower T than confinement.

Backup slides

Correlations for VARIOUS Quark Matters

Transition to hadron gas may be:

- (strong) 1st order
- second order (Critical Point, CP)
- cross-over
- from a supercooled state (scQGP)

Type of phase transition:

its correlation signature:

Strong 1st order QCD phase transition:

(Pratt, Bertsch, Rischke, Gyulassy)

$$R_{\text{out}} \gg R_{\text{side}}$$

Second order QCD phase transition:

→ Cs, S. Hegyi, T. Novák, W.A. Zajc) non-Gaussian shape
 $\alpha(\text{Lévy})$ decreases to 0.5

Cross-over quark matter-hadron gas transition:

(lattice QCD, Buda-Lund hydro fits) hadrons appear from
a region with $T > T_c$

Supercooled QGP (scQGP) -> hadrons:

(T. Cs, L.P. Csernai) pion flash ($R_{\text{out}} \sim R_{\text{side}}$)

same freeze-out for all
strangeness enhancement
no mass-shift of ϕ