

# Hadronic modes in the Quark-Gluon Plasma

*hep-ph/0505080*

Massimo Mannarelli<sup>1,2,3</sup> and Ralf Rapp<sup>1</sup>

[massimo@lns.mit.edu](mailto:massimo@lns.mit.edu)

1) Cyclotron Institute and Physics Department, Texas A&M University, College Station, TX 77843-3366

2) Center for Theoretical Physics, Laboratory for Nuclear Science and Department of Physics

Massachusetts Institute of Technology, Cambridge, MA 02139

3) INFN Bruno Rossi Fellowship

# Outline of the talk

- Introduction
  - QCD Phase Diagram
  - Mesonic states above  $T_c$
- Our model
  - Dirac-Brueckner approach
  - Medium effects
  - Numerical results
- Conclusions and outlook

# Introduction

- **Lattice QCD (IQCD):**  $T_c \sim 170$  MeV and  $\epsilon_c \sim 1\text{GeV}/\text{fm}^3$
- **pQCD:** For  $T \gg T_c$  relevant degrees of freedom quarks and gluons with a screened (rather than antiscreened) color interaction
- **Experiment:** At RHIC reached  $\epsilon > \epsilon_c$  and rapid thermalization of the fireball and collective behaviors observed
- **Phenomenology:** Hydrodynamics describes with good accuracy the properties of the matter obtained at RHIC (not for very high  $p_t$ )

# Introduction

- **Lattice QCD (IQCD):**  $T_c \sim 170$  MeV and  $\epsilon_c \sim 1\text{GeV}/\text{fm}^3$
- **pQCD:** For  $T \gg T_c$  relevant degrees of freedom quarks and gluons with a screened (rather than antiscreened) color interaction
- **Experiment:** At RHIC reached  $\epsilon > \epsilon_c$  and rapid thermalization of the fireball and collective behaviors observed
- **Phenomenology:** Hydrodynamics describes with good accuracy the properties of the matter obtained at RHIC (not for very high  $p_t$ )

Hydro requires **thermalization times below 1 fm/c.**

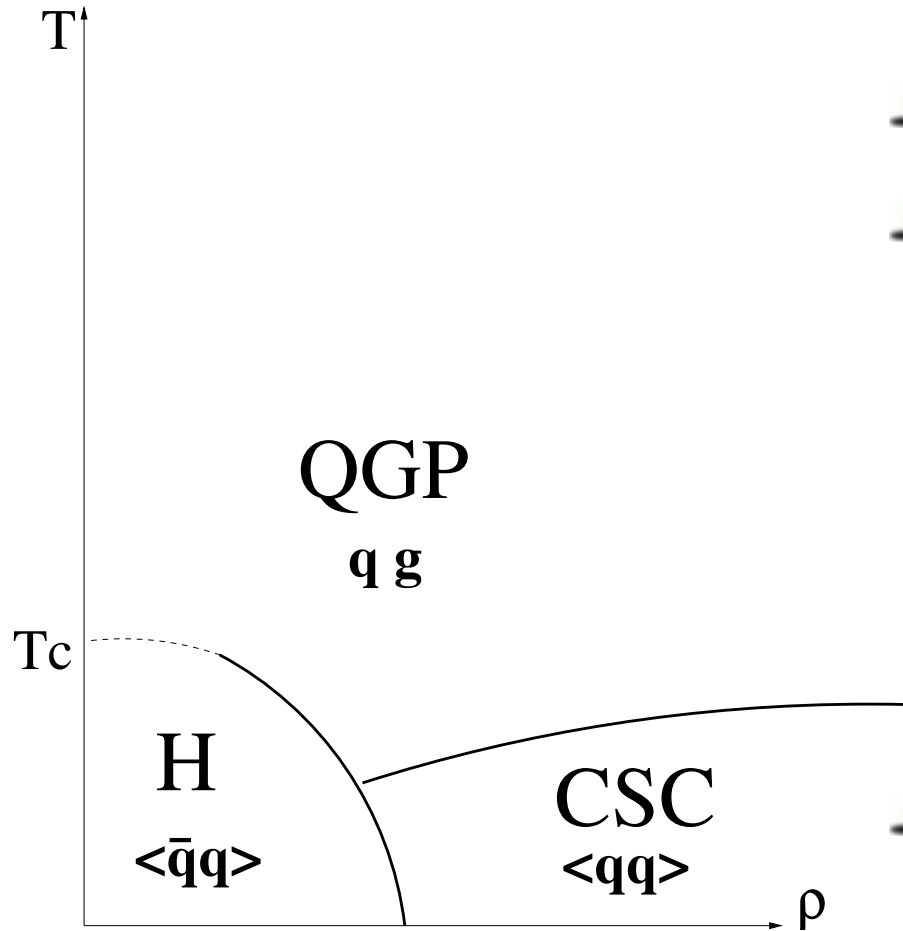
# Introduction

- **Lattice QCD (IQCD):**  $T_c \sim 170$  MeV and  $\epsilon_c \sim 1$  GeV/fm<sup>3</sup>
- **pQCD:** For  $T \gg T_c$  relevant degrees of freedom quarks and gluons with a screened (rather than antiscreened) color interaction
- **Experiment:** At RHIC reached  $\epsilon > \epsilon_c$  and rapid thermalization of the fireball and collective behaviors observed
- **Phenomenology:** Hydrodynamics describes with good accuracy the properties of the matter obtained at RHIC (not for very high  $p_t$ )

Hydro requires **thermalization times below 1 fm/c.**

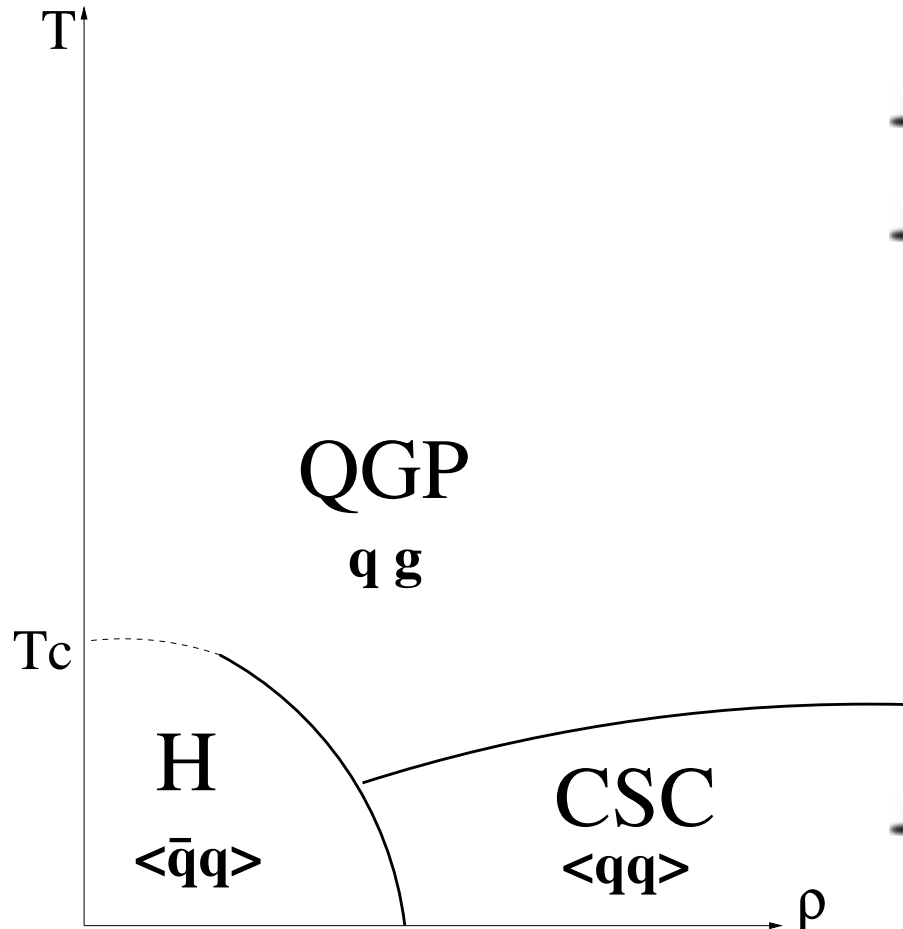
- One mechanism of thermalization: scattering of quarks in resonant states (for heavy quarks see Rapp after this talk in room 0.83)

# QCD phase diagram



- No isospin axis
- Three possible phases
  - H Hadronic
  - QGP Quark-gluon plasma
  - CSC Color Superconductor
- $T_c \sim 170$  MeV from lattice simulations

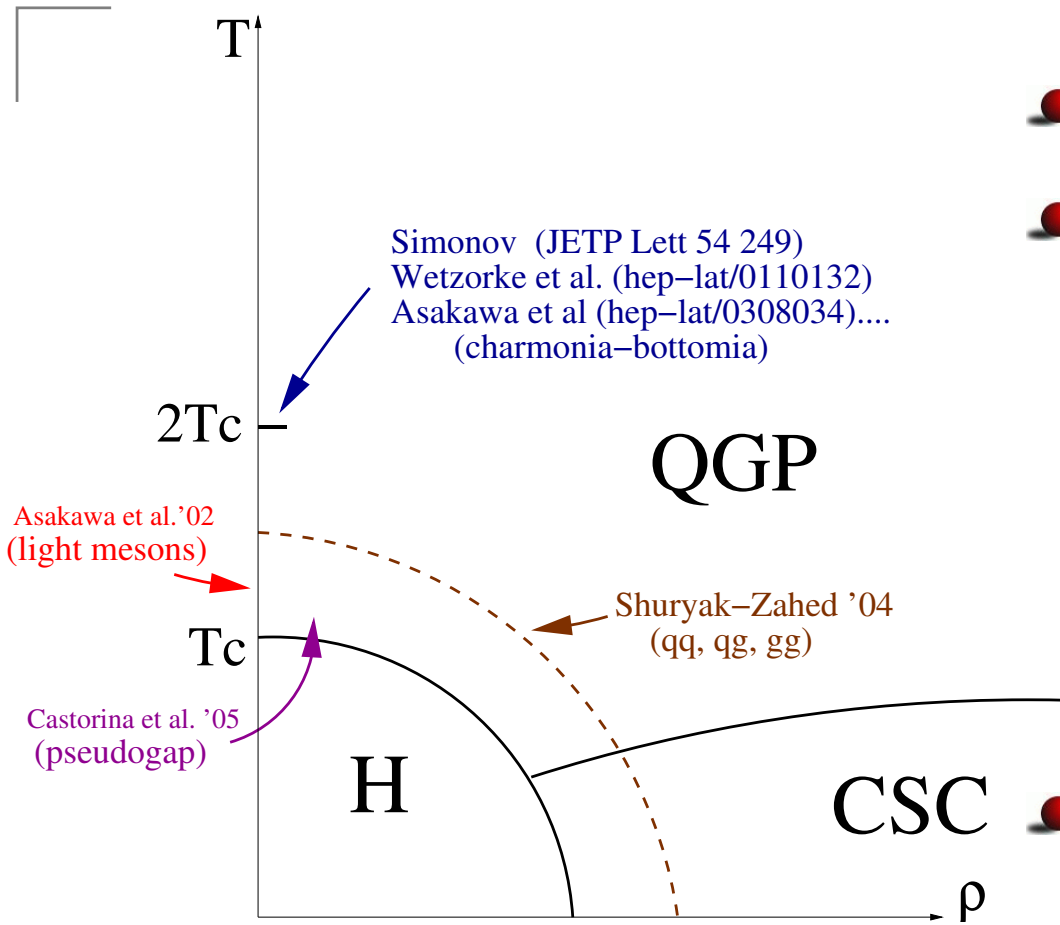
# QCD phase diagram



- No isospin axis
- Three possible phases
  - H Hadronic
  - QGP Quark-gluon plasma
  - CSC Color Superconductor
- $T_c \sim 170$  MeV from lattice simulations

**Plasma:** System of charged and neutral particles which exhibit collective behavior.

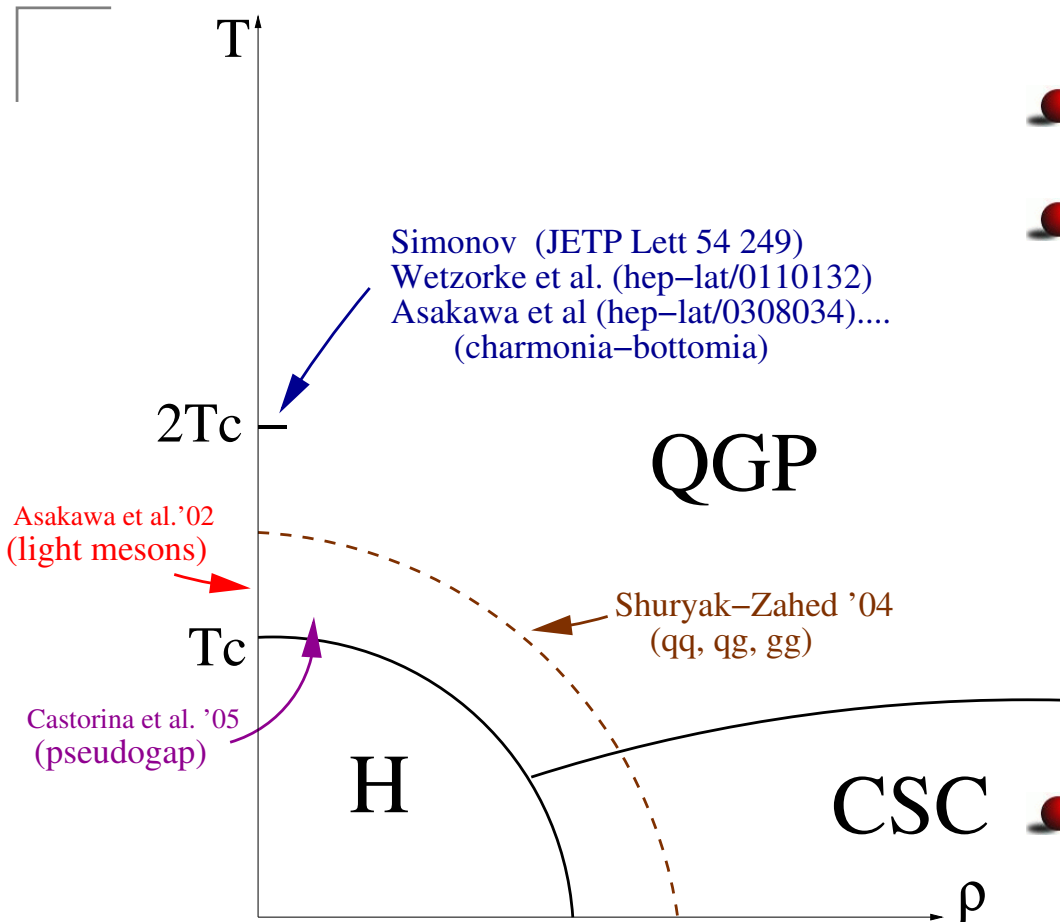
# QCD phase diagram



- No isospin axis
- Three possible phases
- H Hadronic
- QGP Quark-gluon plasma
- CSC Color Superconductor
- $T_c \sim 170$  MeV from lattice simulations

**Plasma:** System of charged and neutral particles which exhibit collective behavior.

# QCD phase diagram



- No isospin axis
- Three possible phases
- H Hadronic
- QGP Quark-gluon plasma
- CSC Color Superconductor
- $T_c \sim 170$  MeV from lattice simulations

**Plasma:** System of charged and neutral particles which exhibit collective behavior. **A definition a bit too generic**

# Our model

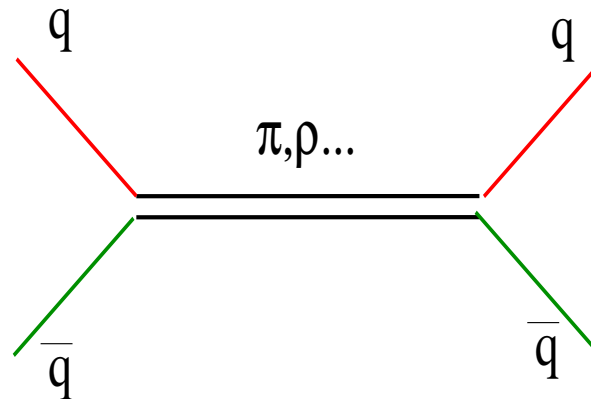
We assume that

- The QGP consists of quarks, gluons. Quarks can form bound (resonant) **mesonic states**.
- **Potential model** with the quark interactions in the singlet and octet channels can be extracted from lattice QCD simulations.
- The Dirac-Brueckner approach is used to take into account **medium effects** on quarks (quasiparticle description)

# Our model

We assume that

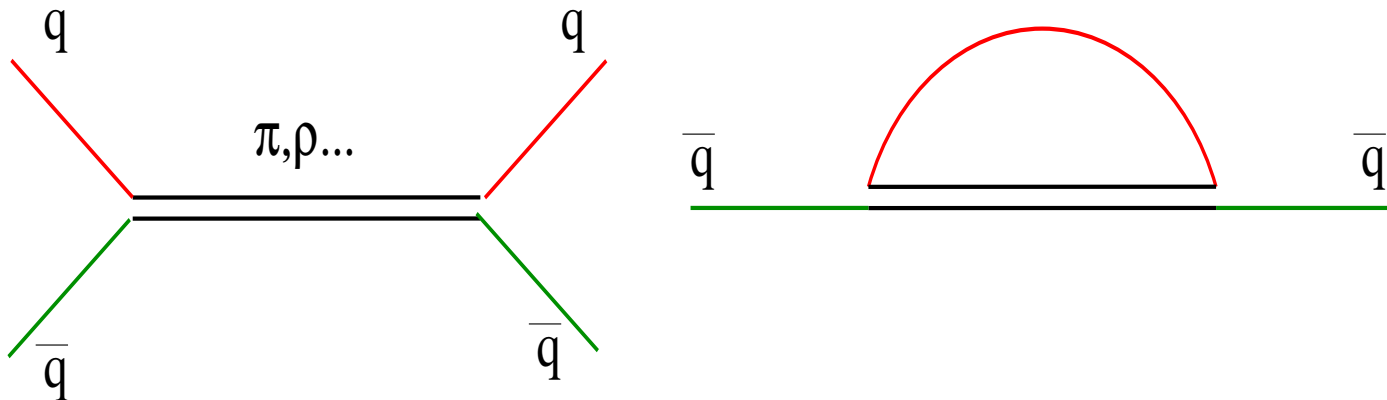
- The QGP consists of quarks, gluons. Quarks can form bound (resonant) **mesonic states**.
- **Potential model** with the quark interactions in the singlet and octet channels can be extracted from lattice QCD simulations.
- The Dirac-Brueckner approach is used to take into account **medium effects** on quarks (quasiparticle description)



# Our model

We assume that

- The QGP consists of quarks, gluons. Quarks can form bound (resonant) **mesonic states**.
- **Potential model** with the quark interactions in the singlet and octet channels can be extracted from lattice QCD simulations.
- The Dirac-Brueckner approach is used to take into account **medium effects** on quarks (quasiparticle description)

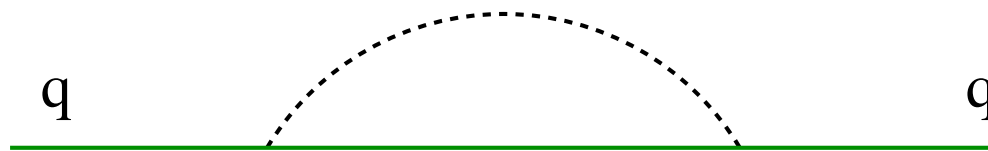


# Our model

We assume that

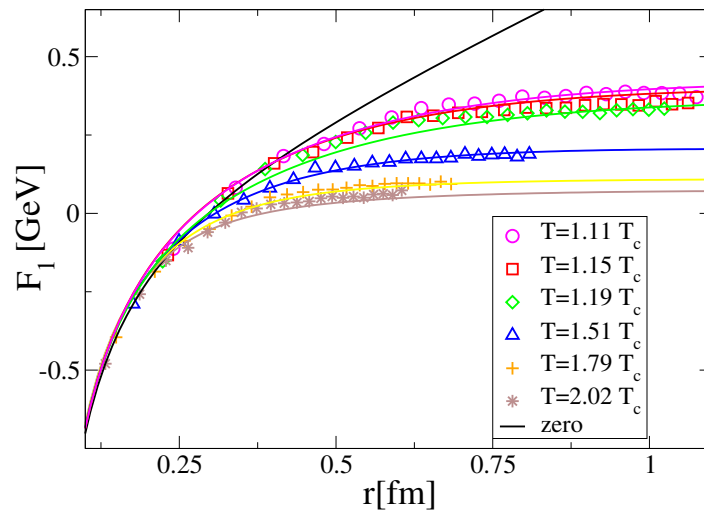
- The QGP consists of quarks, gluons. Quarks can form bound (resonant) **mesonic states**.
- **Potential model** with the quark interactions in the singlet and octet channels can be extracted from lattice QCD simulations.
- The Dirac-Brueckner approach is used to take into account **medium effects** on quarks (quasiparticle description)

Note that the typical HTL quark self-energy diagram is not **directly** included in our procedure



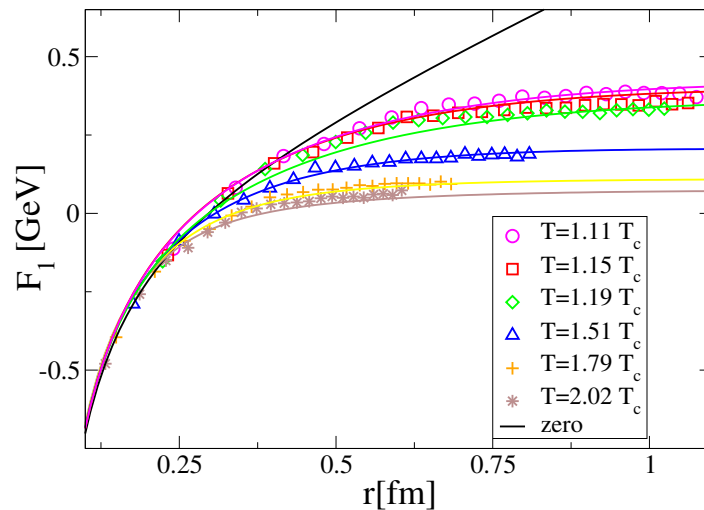
# Parametrization of the Free energy

We extract the potential from the variation of the free-energy due to a  $q\bar{q}$  pair (*Petreczky et al. '04*):



# Parametrization of the Free energy

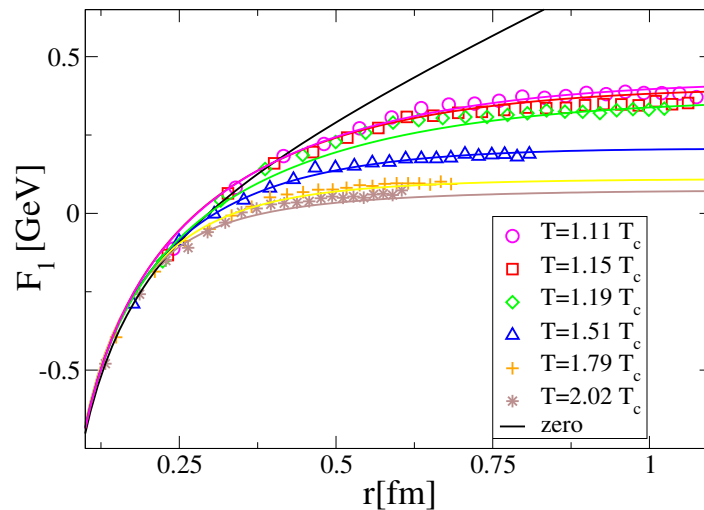
We extract the potential from the variation of the free-energy due to a  $q\bar{q}$  pair (*Petreczky et al. '04*):



● At small distance same behavior

# Parametrization of the Free energy

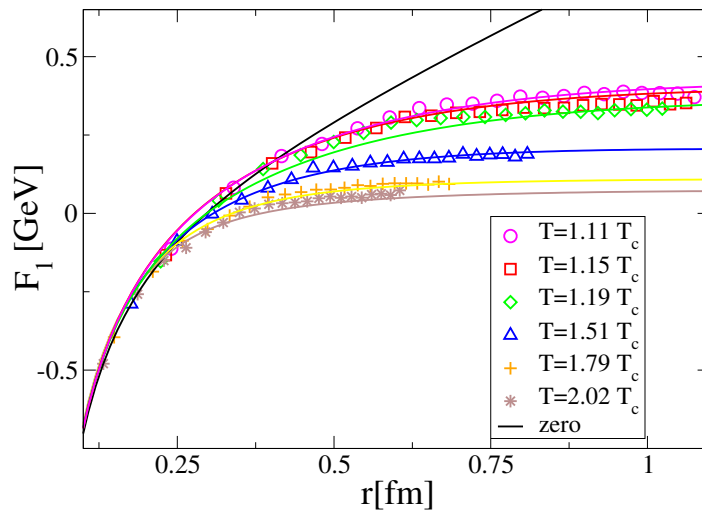
We extract the potential from the variation of the free-energy due to a  $q\bar{q}$  pair (*Petreczky et al. '04*):



- At small distance same behavior
- At large distance screening effects

# Parametrization of the Free energy

We extract the potential from the variation of the free-energy due to a  $q\bar{q}$  pair (*Petreczky et al. '04*):



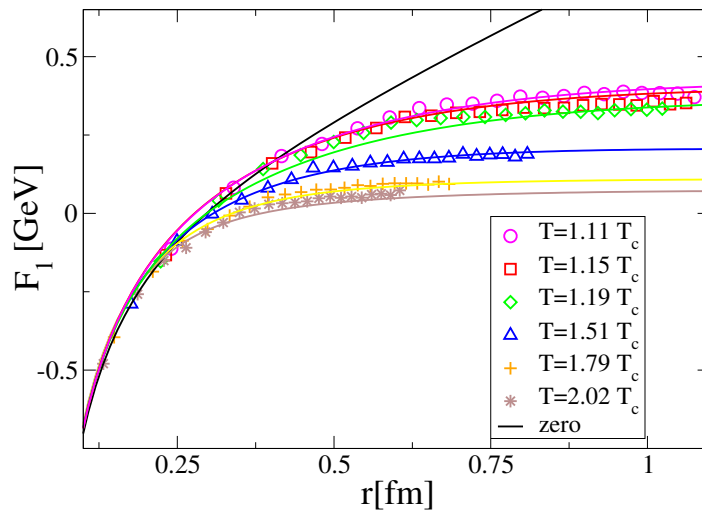
- At small distance same behavior
- At large distance screening effects

Fitting function (similar to *Karsch et al. '88*) :

$$F_1(r, T) = -\frac{\alpha}{r} e^{-A\mu r} + \frac{\sigma}{\mu} (1 - e^{-\mu r})$$

# Parametrization of the Free energy

We extract the potential from the variation of the free-energy due to a  $q\bar{q}$  pair (*Petreczky et al. '04*):



- At small distance same behavior
- At large distance screening effects

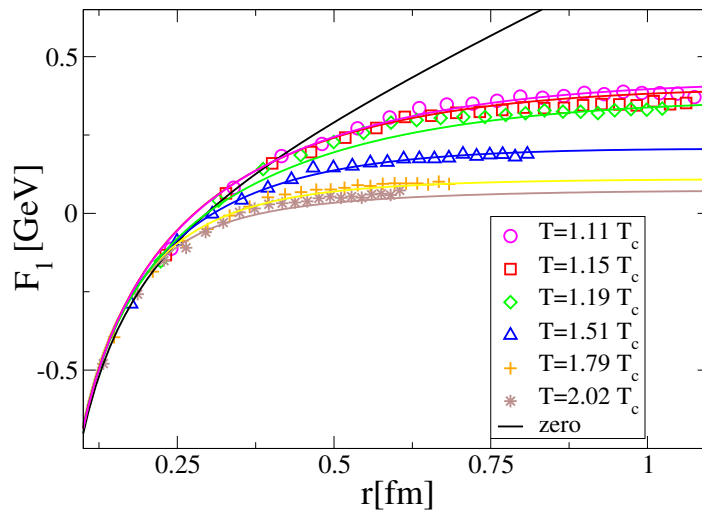
Fitting function (similar to *Karsch et al. '88*) :

$$F_1(r, T) = -\frac{\alpha}{r} e^{-A\mu r} + \frac{\sigma}{\mu} (1 - e^{-\mu r})$$

$A$  fitting function;

# Parametrization of the Free energy

We extract the potential from the variation of the free-energy due to a  $q\bar{q}$  pair (*Petreczky et al. '04*):



- At small distance same behavior
- At large distance screening effects

Fitting function (similar to *Karsch et al. '88*) :

$$F_1(r, T) = -\frac{\alpha}{r} e^{-A\mu r} + \frac{\sigma}{\mu} (1 - e^{-\mu r})$$

$A$  fitting function;  $\mu$  "screening mass"

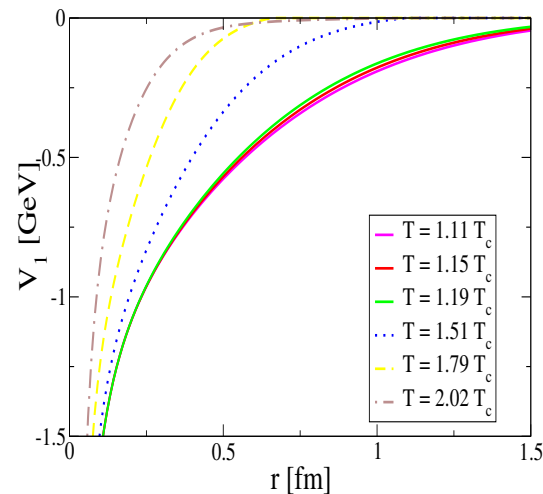
# Potential

Internal energy

$$E_1 = F_1 - T \frac{dF_1}{dT}.$$

We assume that the asymptotic value of the internal is the in-medium quark mass. Therefore the potential in the color-singlet channel is

$$V_1(r, T) = E_1(r, T) - E_1(\infty, T) \quad V_8 = -\frac{1}{8} V_1$$



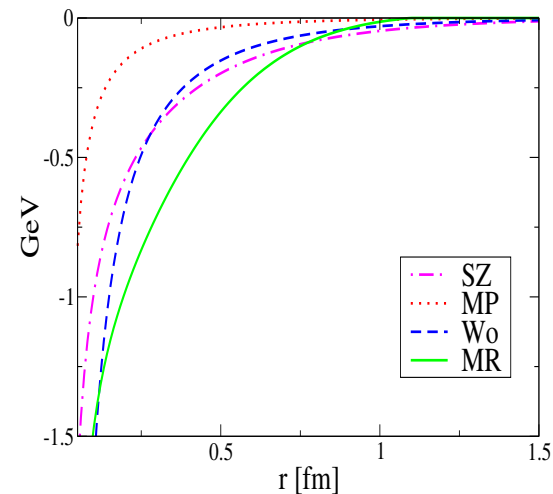
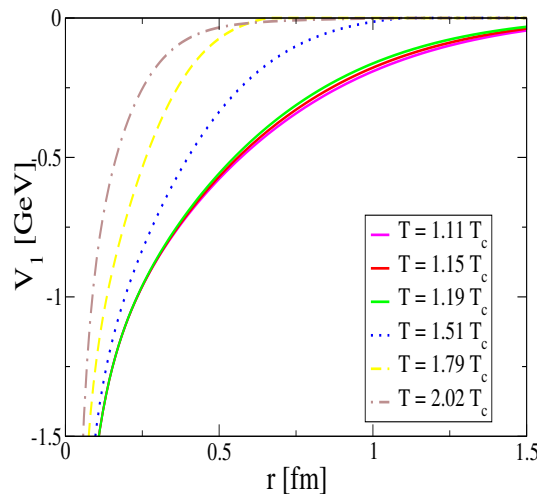
# Potential

Internal energy

$$E_1 = F_1 - T \frac{dF_1}{dT}.$$

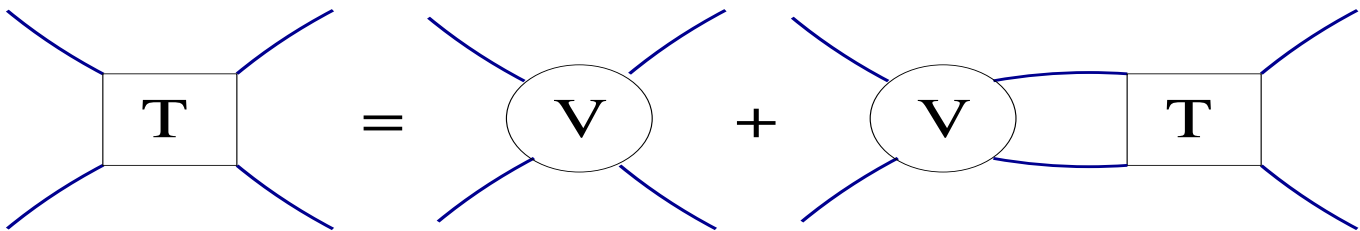
We assume that the asymptotic value of the internal is the in-medium quark mass. Therefore the potential in the color-singlet channel is

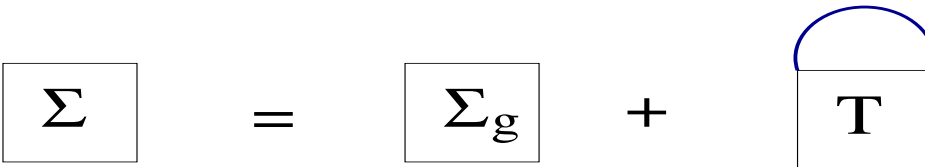
$$V_1(r, T) = E_1(r, T) - E_1(\infty, T) \quad V_8 = -\frac{1}{8} V_1$$



# Self-consistency

The full set of Dirac-Brueckner equations

**B.S.** 

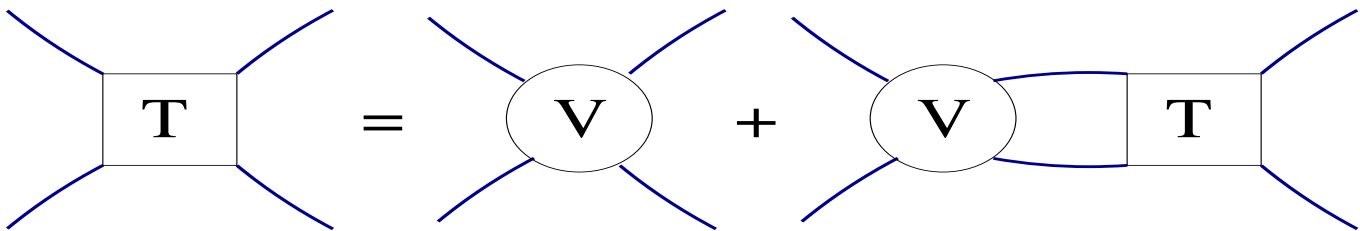
**"Link"** 

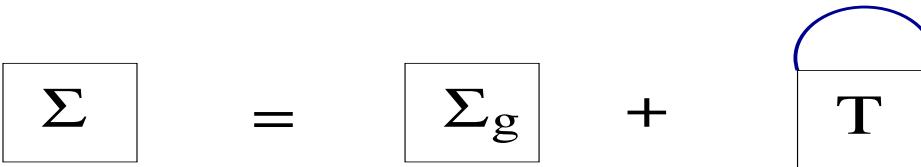
**S.D.** 

has to be solved self-consistently (by [iteration](#)).

# Self-consistency

The full set of Dirac-Brueckner equations

**B.S.** 

**"Link"** 

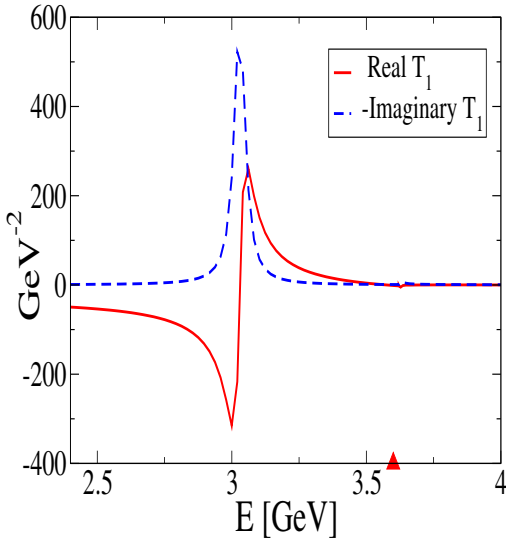
**S.D.** 

has to be solved self-consistently (by [iteration](#)).

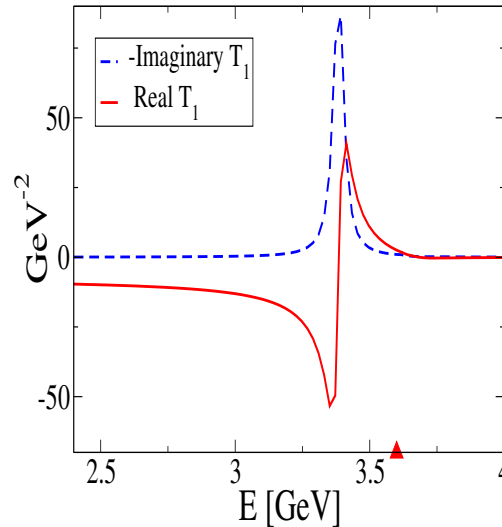
$\Sigma_g$  is a "gluon-induced" self-energy, parameterized as a mass  $m$  in the dispersion law:

$$\omega_k = \sqrt{k^2 + m^2 + \Sigma_R(\omega_k, k)}$$

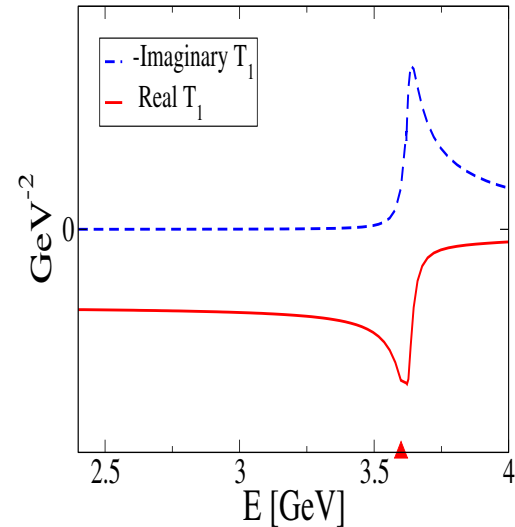
# A check: " $J/\Psi$ "



$$T = 1.2 T_c$$



$$T = 1.5 T_c$$



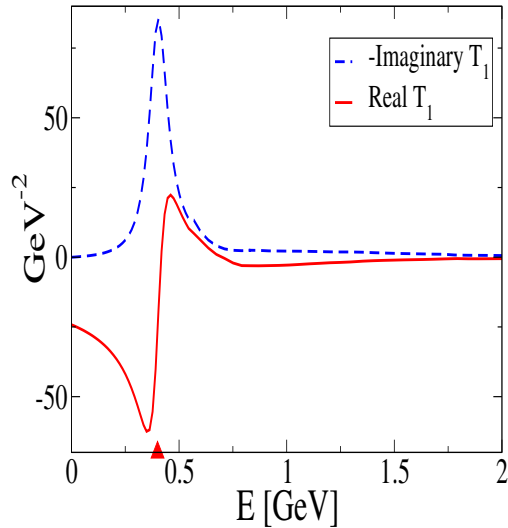
$$T = 2.0 T_c$$

**Real** and **imaginary** part of  $T$ -matrix in the singlet channel for the in medium " $J/\Psi$ " meson as a function of  $CM$  energy  $E$  charm effective mass  $m = 1.8$  GeV

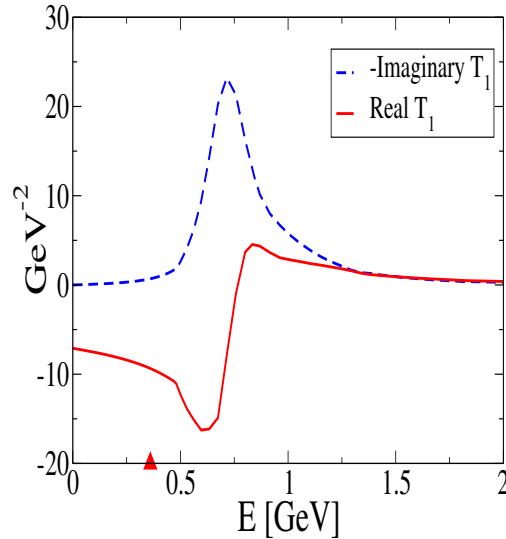
**No self-consistency:**  $\Sigma_R = 0$  and  $\Sigma_I = -10$  MeV.

Dissociation at  $T_D \sim 2T_c$ .

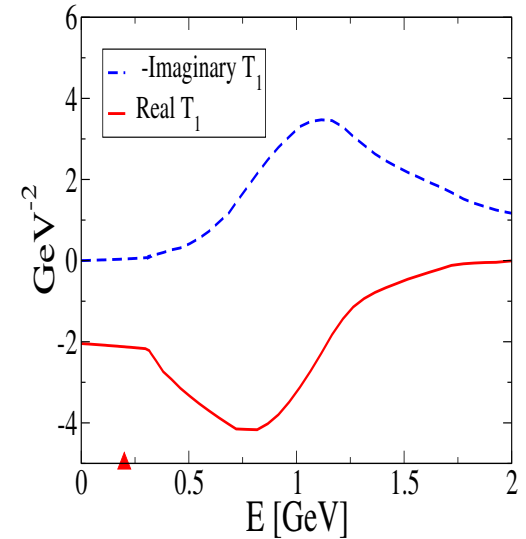
# Light quarks T-matrix (singlet)



$$T = 1.2 T_c$$



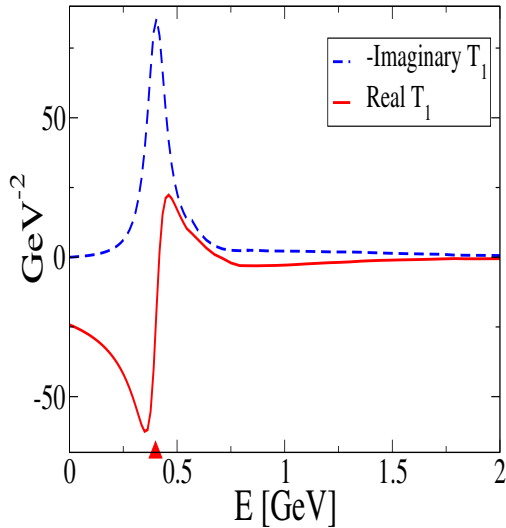
$$T = 1.5 T_c$$



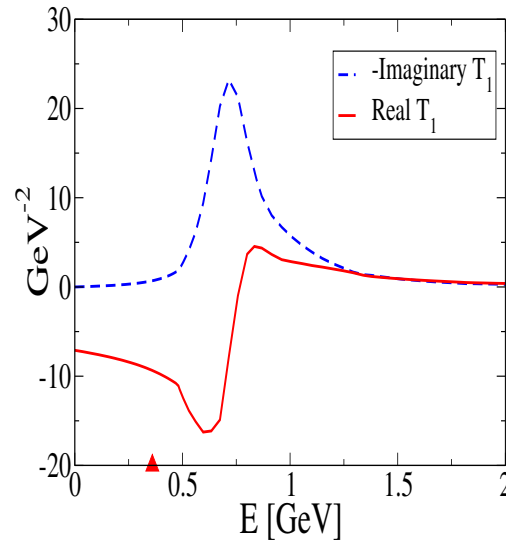
$$T = 1.75 T_c$$

**Real** and **imaginary** part of the light-quark (on-shell)  $T$ -matrix in the singlet color channel with a "gluon-induced" mass  $m = 0.1$  GeV.

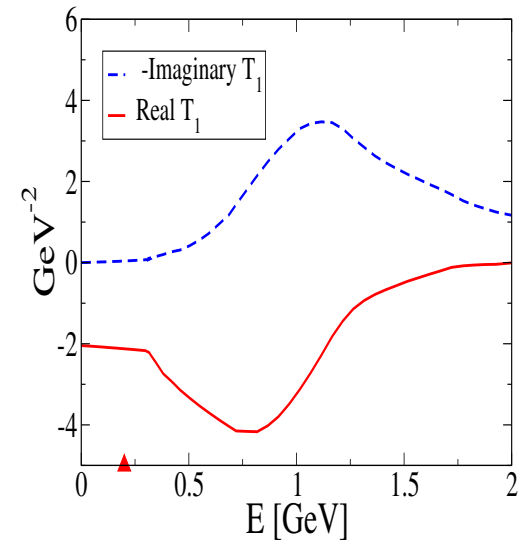
# Light quarks T-matrix (singlet)



$$T = 1.2 T_c$$



$$T = 1.5 T_c$$

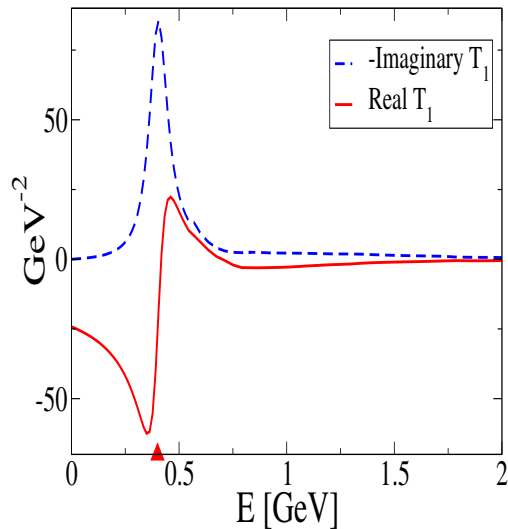


$$T = 1.75 T_c$$

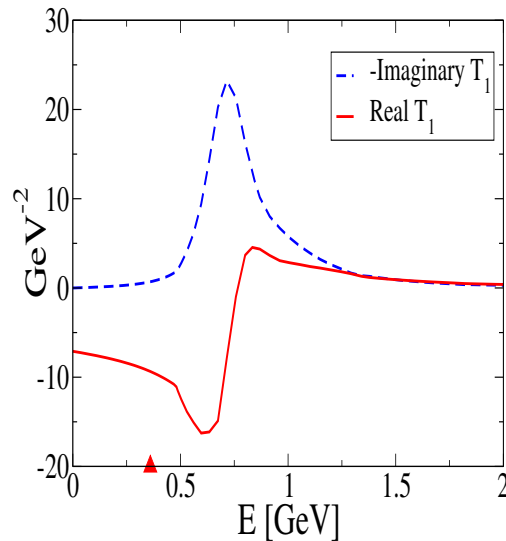
**Real** and **imaginary** part of the light-quark (on-shell)  $T$ -matrix in the singlet color channel with a "gluon-induced" mass  $m = 0.1$  GeV.

- Dissociation temperature  $T \sim 1.2 T_c$ .

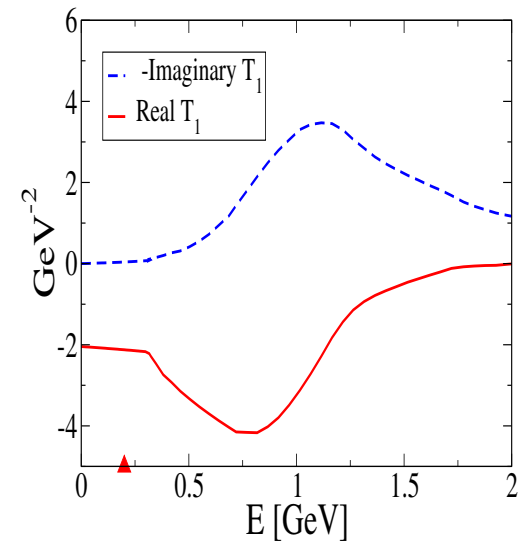
# Light quarks T-matrix (singlet)



$$T = 1.2 T_c$$



$$T = 1.5 T_c$$

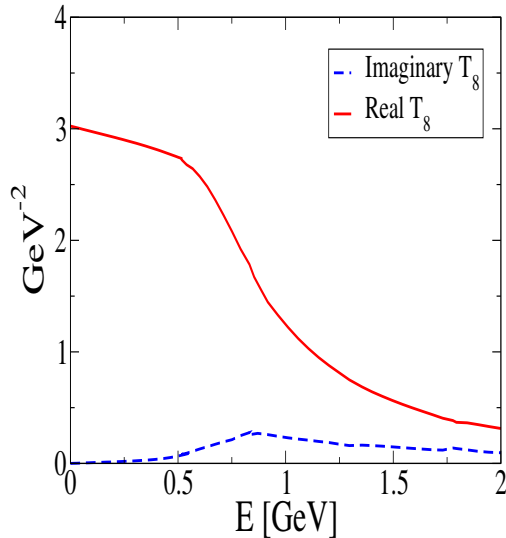


$$T = 1.75 T_c$$

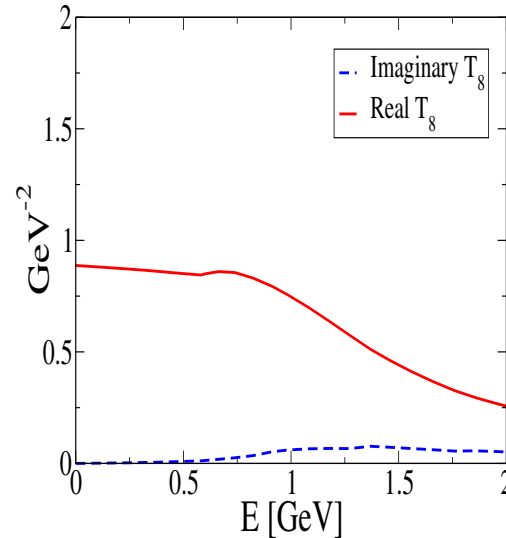
**Real** and **imaginary** part of the light-quark (on-shell)  $T$ -matrix in the singlet color channel with a "gluon-induced" mass  $m = 0.1$  GeV.

- Dissociation temperature  $T \sim 1.2 T_c$ .
- For  $T > 1.2 T_c$  mesonic states survive as resonant states.

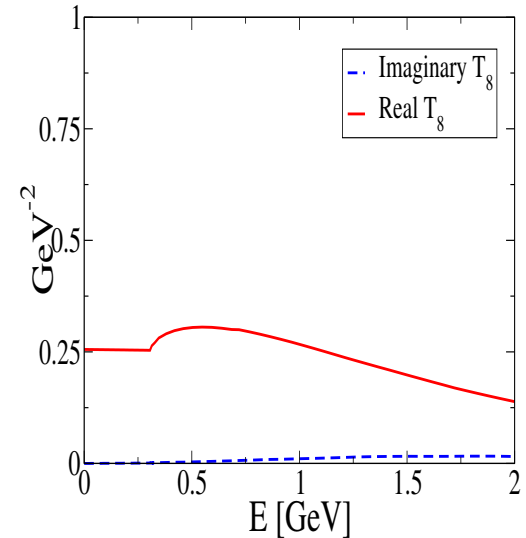
# Light quarks T-matrix (octet)



$$T = 1.2 T_c$$



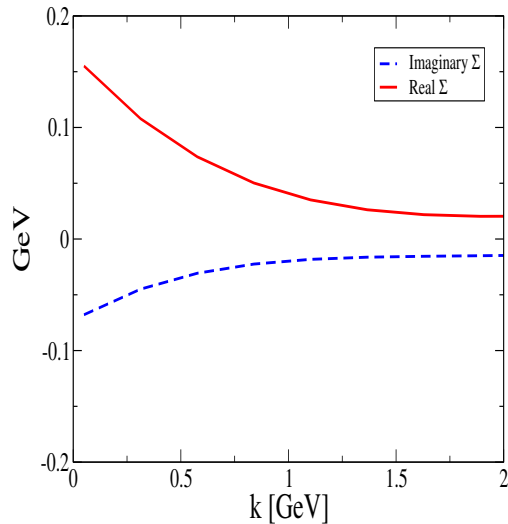
$$T = 1.5 T_c$$



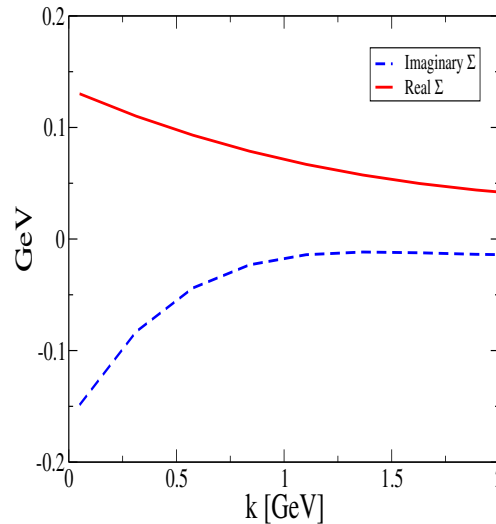
$$T = 1.75 T_c$$

**Real** and **imaginary** part of the light-quark (on-shell)  $T$ -matrix in the color octet channel;  $m = 0.1$  GeV.

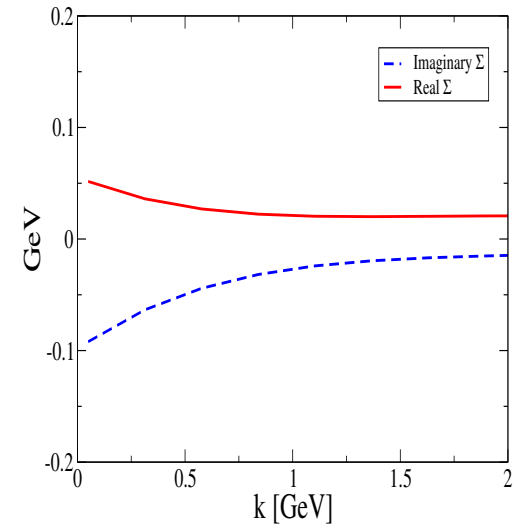
# Light quarks Self-energy



$$T = 1.2 T_c$$



$$T = 1.5 T_c$$

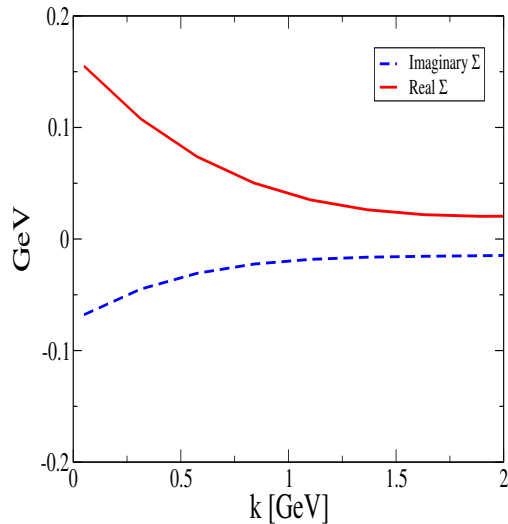


$$T = 1.75 T_c$$

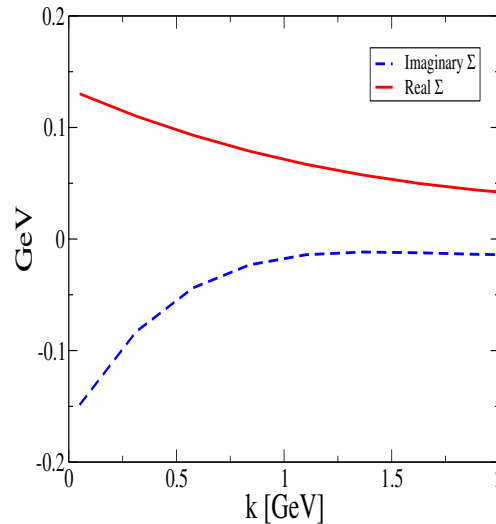
**Real** and **imaginary** part of the light-quark (on-shell) self-energy (singlet+octet);  $m = 0.1$  GeV.

- Real part dominated by the octet contribution.

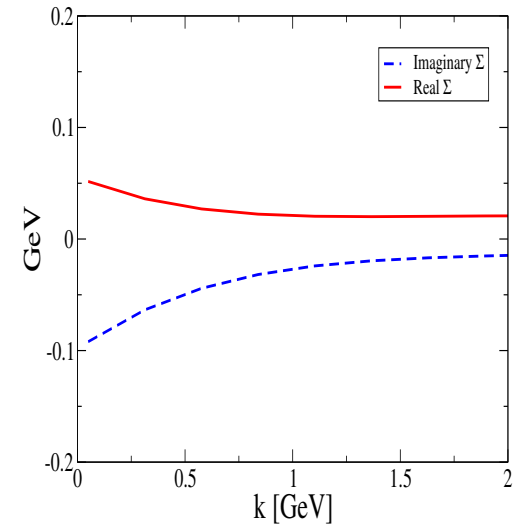
# Light quarks Self-energy



$$T = 1.2 T_c$$



$$T = 1.5 T_c$$



$$T = 1.75 T_c$$

**Real** and **imaginary** part of the light-quark (on-shell) self-energy (singlet+octet);  $m = 0.1$  GeV.

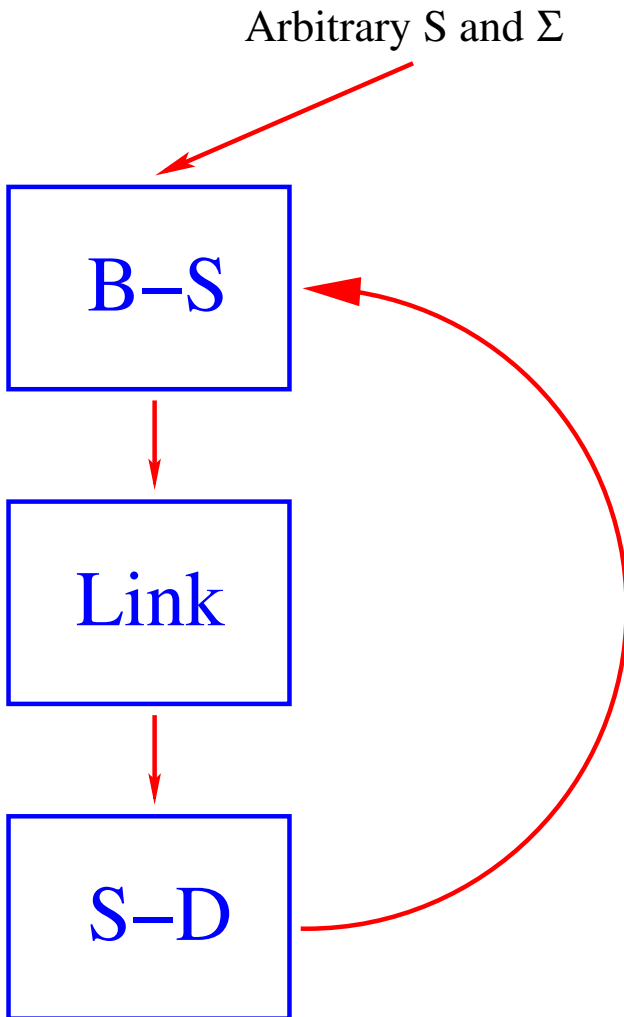
- Real part dominated by the octet contribution.
- Imaginary part dominated by the singlet contribution.

# Conclusions and outlook

- We have analyzed the mechanism of scattering of quarks in bound or resonant states in the QGP  $T \sim 1 - 2T_c$
- "Quasiparticles" acquire a large mass  $\sim 150$  MeV from the octet channel and a large width  $\sim 200$  MeV from the singlet
- The mechanism which produces the large width is the scattering into resonant states
- This could be a mechanism to quickly thermalize the QGP and to explain the collective behavior observed.
- Next step is to study heavy quarks and to find a way to consider in a self-consistent way the gluons.

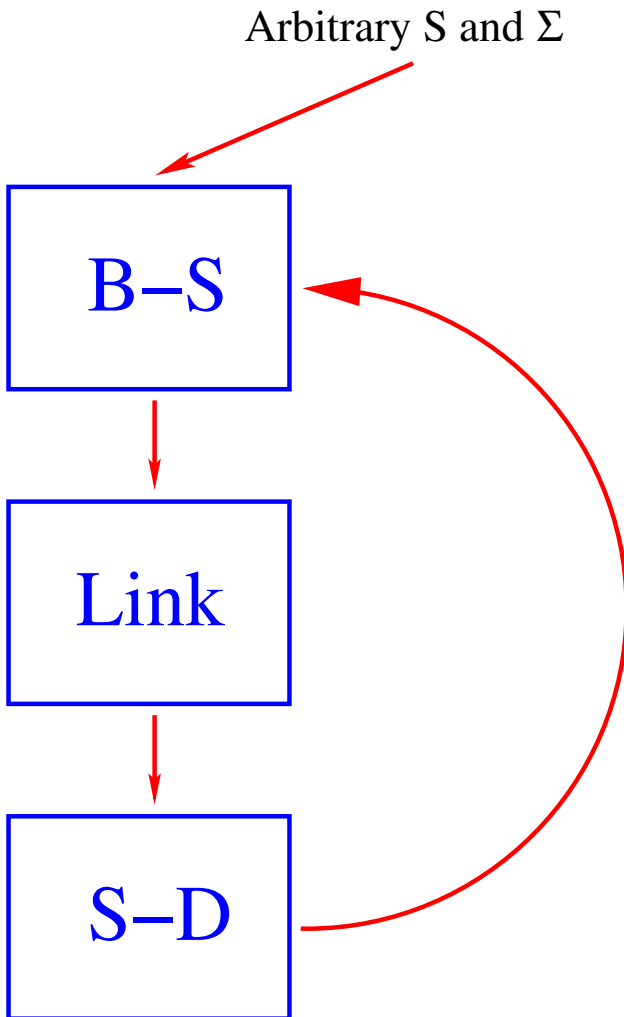
# Numerical Solution

We "solve" the Dirac-Brueckner equations by iteration



# Numerical Solution

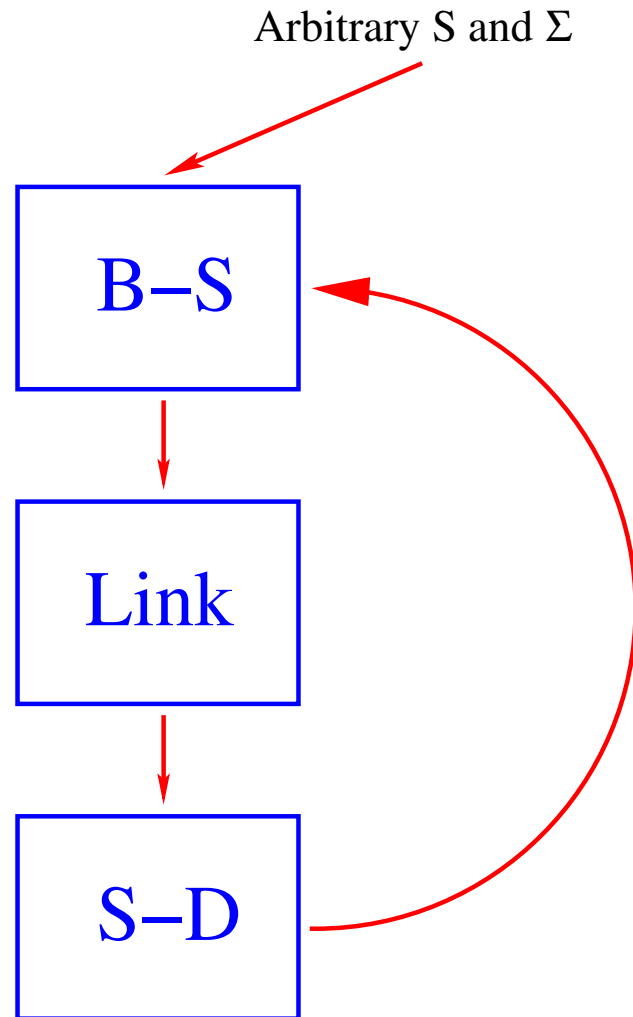
We "solve" the Dirac-Brueckner equations by iteration



- Begin with an arbitrary self-energy
- Solve the Bethe-Salpeter
- Determine the Self-energy
- Solve the Schwinger-Dyson
- Iterate

# Numerical Solution

We "solve" the Dirac-Brueckner equations by iteration



- Begin with an arbitrary self-energy
- Solve the Bethe-Salpeter
- Determine the Self-energy
- Solve the Schwinger-Dyson
- Iterate

Check of stability and independence on step 1

# Potential models ( $\mu = 0$ )

Charm and bottom quarks are heavy. In bound states and for  $m \ll E_k$  one can use non-relativistic potential model

$$\left( 2m_a + \frac{\nabla^2}{m_a} + V_1(r, T) \right) \psi_a = M(T)\psi_a ,$$

# Potential models ( $\mu = 0$ )

Charm and bottom quarks are heavy. In bound states and for  $m \ll E_k$  one can use non-relativistic potential model

$$\left( 2m_a + \frac{\nabla^2}{m_a} + V_1(r, T) \right) \psi_a = M(T) \psi_a ,$$

with  $a = b, c$ . Without a medium ( $T \simeq 0$ ) the Cornell potential

$$V_1(r, T) = -\frac{\alpha}{r} + \sigma r$$

works for  $c\bar{c}$  and  $b\bar{b}$ : *Eichten et al.* Phys. Rev. D 17, 3090 (1978).

# Potential models ( $\mu = 0$ )

Charm and bottom quarks are heavy. In bound states and for  $m \ll E_k$  one can use non-relativistic potential model

$$\left( 2m_a + \frac{\nabla^2}{m_a} + V_1(r, T) \right) \psi_a = M(T) \psi_a ,$$

with  $a = b, c$ . Without a medium ( $T \simeq 0$ ) the Cornell potential

$$V_1(r, T) = -\frac{\alpha}{r} + \sigma r$$

works for  $c\bar{c}$  and  $b\bar{b}$ : *Eichten et al.* Phys. Rev. D 17, 3090 (1978).

To fix the parameters

- Phenomenology
- Lattice
- pQCD

# Potential models ( $\mu = 0$ )

Charm and bottom quarks are heavy. In bound states and for  $m \ll E_k$  one can use non-relativistic potential model

$$\left( 2m_a + \frac{\nabla^2}{m_a} + V_1(r, T) \right) \psi_a = M(T) \psi_a ,$$

with  $a = b, c$ . Without a medium ( $T \simeq 0$ ) the Cornell potential

$$V_1(r, T) = -\frac{\alpha}{r} + \sigma r$$

works for  $c\bar{c}$  and  $b\bar{b}$ : *Eichten et al.* Phys. Rev. D 17, 3090 (1978).

To fix the parameters

- Phenomenology
- **Lattice** *Wong '04* and *Mocsy et al. '04*  $T_{J/\Psi}^D \sim 2T_c$
- pQCD

# Dirac-Brueckner approach

Quark-antiquark scattering can be described covariantly by the Bethe-Salpeter (B.S.) equation for the T-matrix

$$T = K + \int KSST,$$

where  $K$  is the interaction kernel and

$$S = S_0 + S_0\Sigma S$$

is the single-particle propagator.

# Dirac-Brueckner approach

Quark-antiquark **scattering** can be described covariantly by the Bethe-Salpeter (B.S.) equation for the T-matrix

$$T = K + \int KSST,$$

where  $K$  is the interaction kernel and

$$S = S_0 + S_0\Sigma S$$

is the single-particle propagator.

**Approximations**

# Dirac-Brueckner approach

Quark-antiquark scattering can be described covariantly by the Bethe-Salpeter (B.S.) equation for the T-matrix

$$T = K + \int KSST,$$

where  $K$  is the interaction kernel and

$$S = S_0 + S_0\Sigma S$$

is the single-particle propagator.

## Approximations

- Only S-wave channel

# Dirac-Brueckner approach

Quark-antiquark **scattering** can be described covariantly by the Bethe-Salpeter (B.S.) equation for the T-matrix

$$T = K + \int KSST,$$

where  $K$  is the interaction kernel and

$$S = S_0 + S_0\Sigma S$$

is the single-particle propagator.

## Approximations

- Only S-wave channel
- For B.S. we use the ladder approximation with  $K \rightarrow V$

# Dirac-Brueckner approach

Quark-antiquark scattering can be described covariantly by the Bethe-Salpeter (B.S.) equation for the T-matrix

$$T = K + \int KSST,$$

where  $K$  is the interaction kernel and

$$S = S_0 + S_0\Sigma S$$

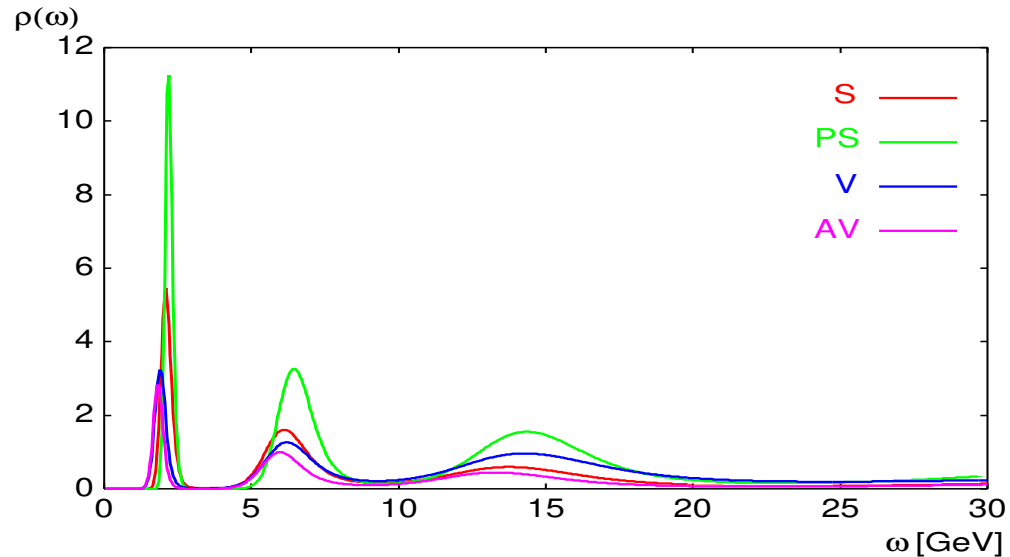
is the single-particle propagator.

## Approximations

- Only S-wave channel
- For B.S. we use the ladder approximation with  $K \rightarrow V$
- Interaction in the octet channel with  $V_8 = -\frac{1}{8}V_1$

# Mesonic states above $T_c$

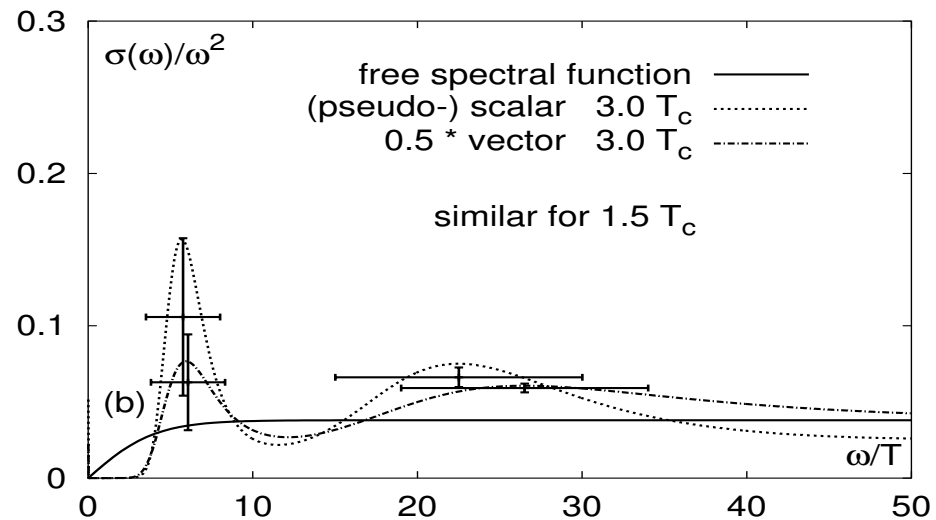
Maximum Entropy Method (MEM) analysis of correlators allows the evaluation of mesonic spectral functions:



Quenched  $32^3 \times 54$  ( $T = 1.4 T_c$ ) with  $m_\pi/m_\rho \simeq 0.7$   
Asakawa et al. Nucl. Phys. A 715 (2003) 701

# Mesonic states above $T_c$

Maximum Entropy Method (MEM) analysis of correlators allows the evaluation of mesonic spectral functions:

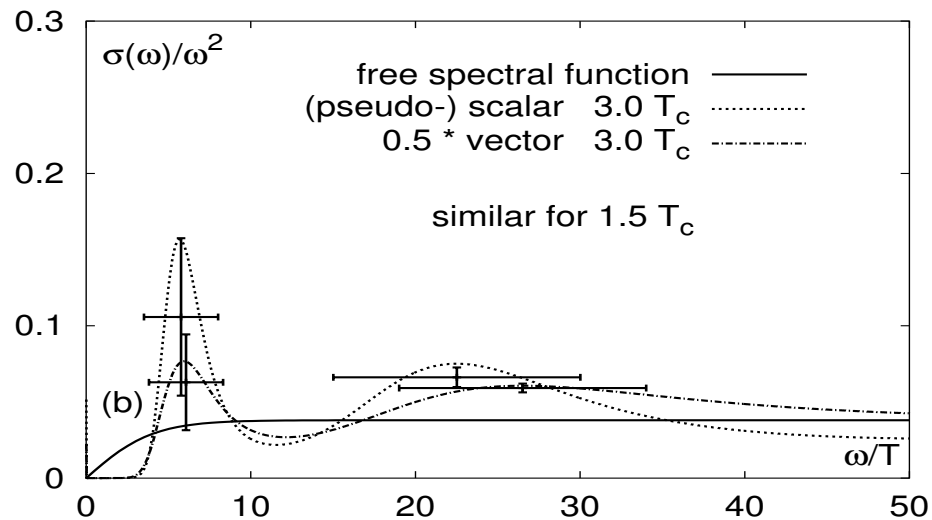


Quenched  $24^3 \times 32$ ; vanishing quark mass

Wetzorke et al. hep-lat/0110132

# Mesonic states above $T_c$

Maximum Entropy Method (MEM) analysis of correlators allows the evaluation of mesonic spectral functions:



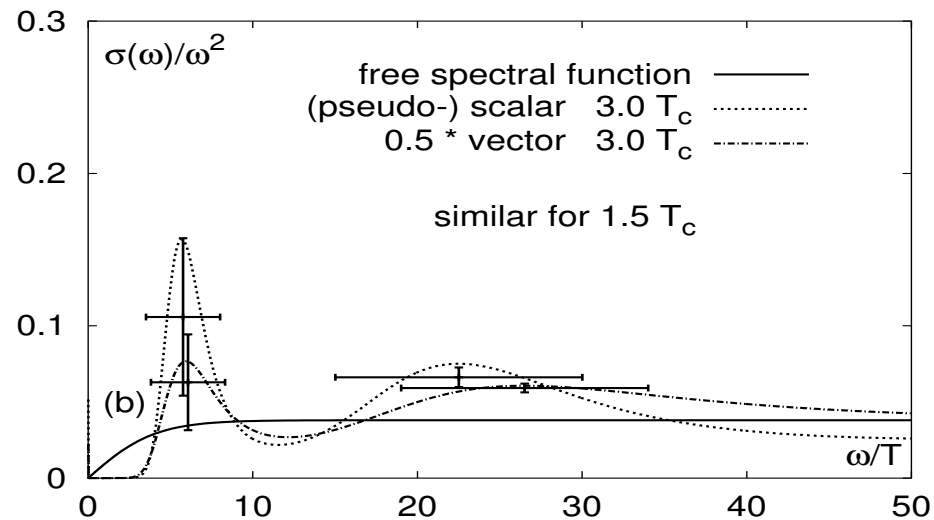
Quenched  $24^3 \times 32$ ; vanishing quark mass

Wetzorke et al. hep-lat/0110132

If no structure is present the spectral function is flat.

# Mesonic states above $T_c$

Maximum Entropy Method (MEM) analysis of correlators allows the evaluation of mesonic spectral functions:



Quenched  $24^3 \times 32$ ; vanishing quark mass

Wetzorke et al. hep-lat/0110132

If no structure is present the spectral function is flat. Powerful method but one needs a big calculation power.